

# *EPR ALIGNMENT OF REFERENCE FRAMES*

*Giulio Chiribella, Rui Chao, and Yuxiang Yang*  
*Tsinghua University, Beijing*

May 6th 2015, Quantum Hamiltonian Complexity Reunion,  
Simons Institute for the Theory of Computing.



RECRUITMENT  
PROGRAM OF GLOBAL EXPERTS



National Natural Science  
Foundation of China



# Application of non-locality to precision measurements?

Non-locality has applications to a number of pure information-theoretic tasks.

Can we find applications to **physical** tasks, such as **clock synchronization** and **direction alignment**?

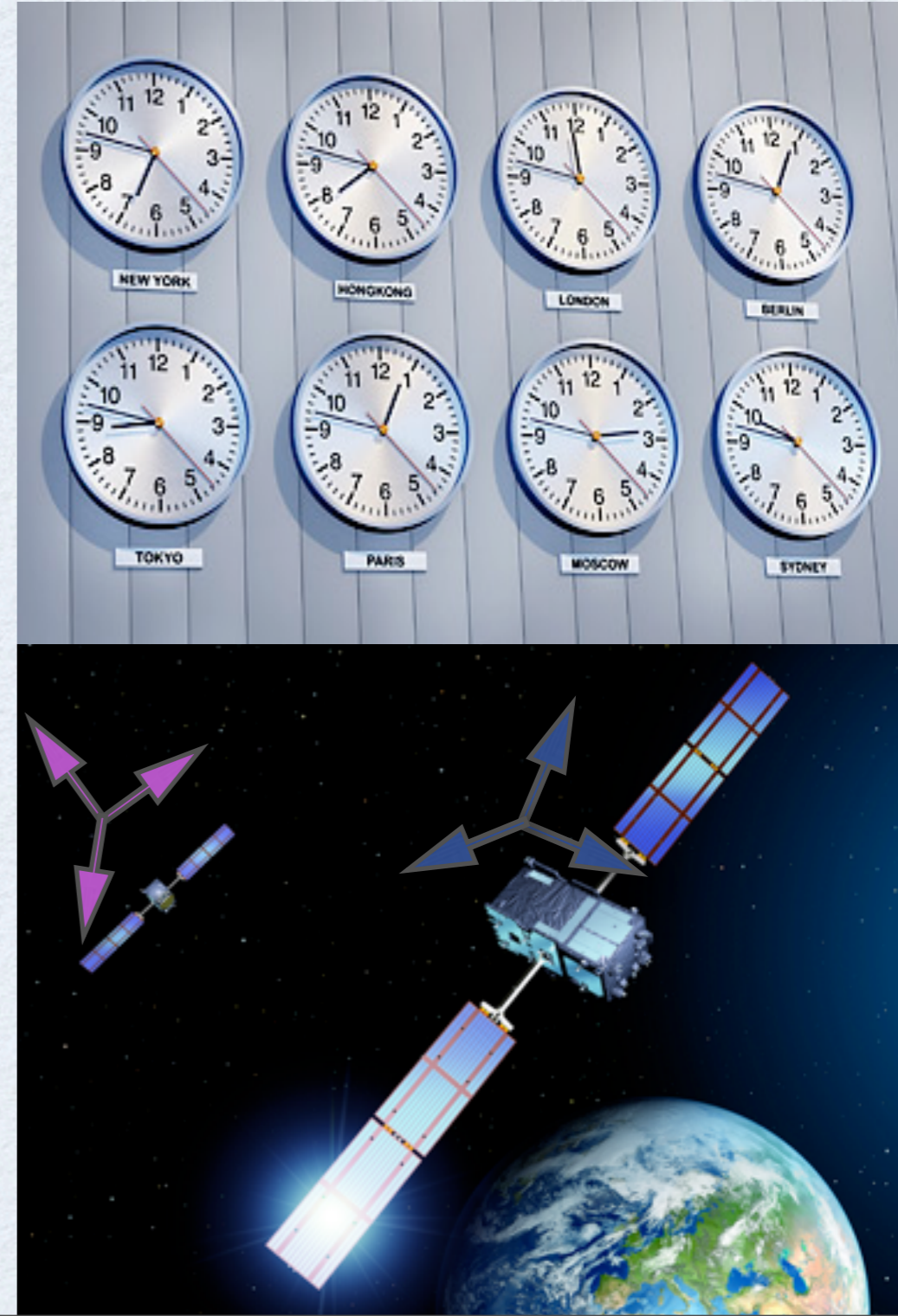
**quantum reference frames:**

Aharonov-Kaufherr PRD 1984,

Gisin-Popescu PRL 1998,

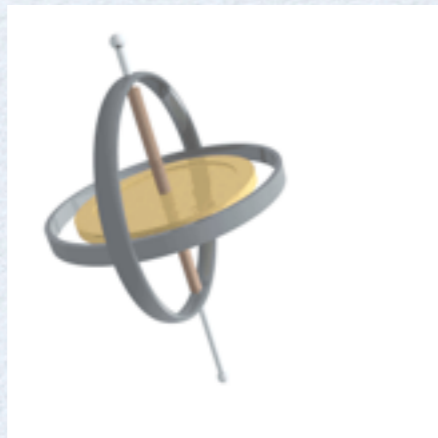
Peres-Scudo, 2001, ...

Bartlett-Rudolph-Spekkens RMP 2007



# Gyroscopes

Classical gyroscope = physical system whose angular momentum indicates a direction in space

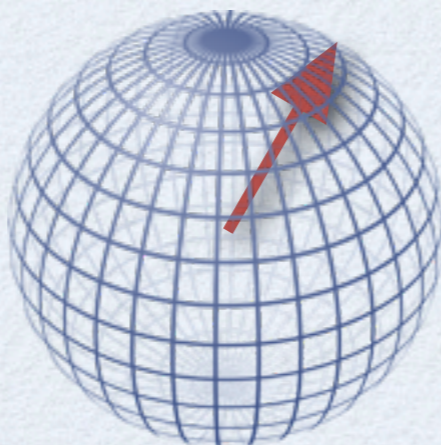


large angular momentum

→ more stable gyroscope

Quantum gyroscope = quantum system whose angular momentum indicates a direction in space

e.g. a spin- $j$  particle

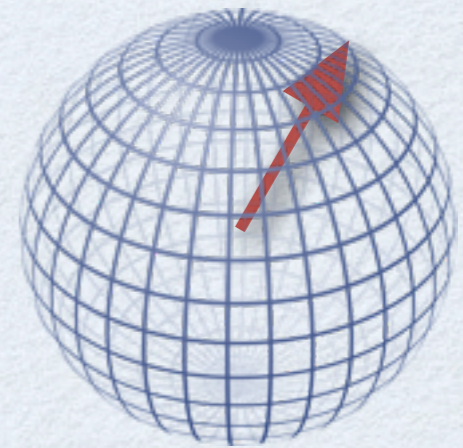


large  $j$

→ more precise gyroscope

# Spin $j$ degrees of freedom

$2j+1$  Hilbert space



Rotation of an angle  $\varphi$  around the axis  $\mathbf{n} = \begin{pmatrix} n_x \\ n_y \\ n_z \end{pmatrix}$ :

$$e^{-i\varphi \mathbf{n} \cdot \mathbf{J}}$$

where  $J_x, J_y, J_z$  are the **angular momentum operators**

$$[j_x, j_y] = i j_z \quad [j_y, j_z] = i j_x \quad [j_z, j_x] = i j_y$$

$$j_x^2 + j_y^2 + j_z^2 = j(j+1) I$$

# Quantifying the error

To find out the direction, one has to perform a **measurement**,

mathematically described by a POVM  $P(d\hat{\mathbf{n}})$

whose outcome gives an estimate  $\hat{\mathbf{n}}$  of the unknown direction.

The error is quantified by the **worst-case square distance**

$$\langle d^2 \rangle = \sup_{\mathbf{n}} \int p(d\hat{\mathbf{n}} | \mathbf{n}) \|\hat{\mathbf{n}} - \mathbf{n}\|^2$$

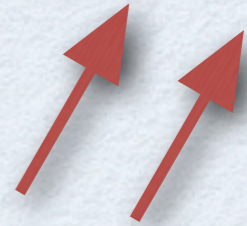
$$p(d\hat{\mathbf{n}} | \mathbf{n}) = \text{Tr} [P(d\hat{\mathbf{n}}) \rho_{\mathbf{n}}]$$

# What was known: $j = 1/2$

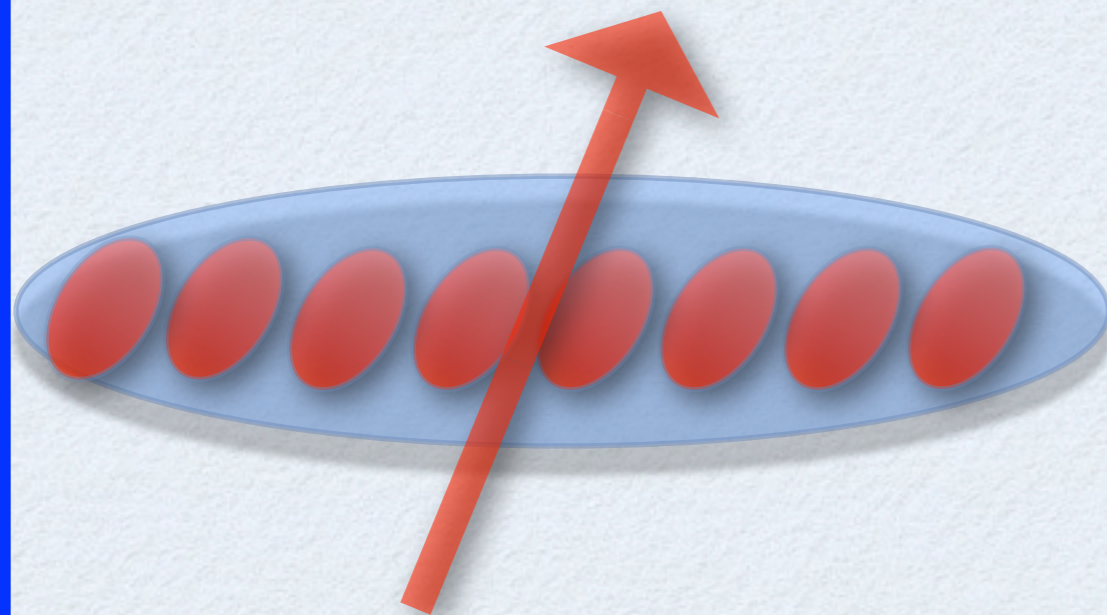
Gisin-Popescu PRL 1998:



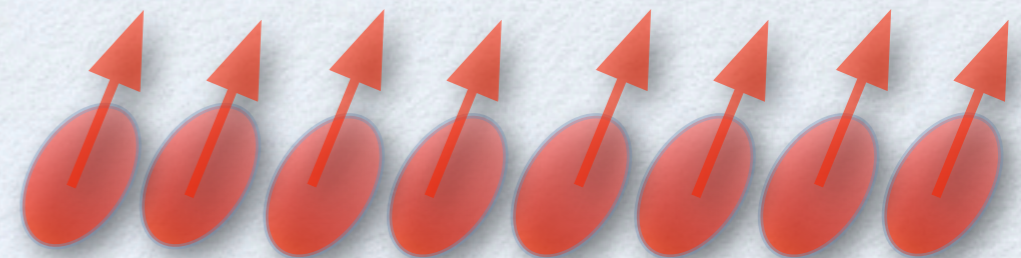
is better than



GC et al PRL 2004, Bagan et al PRA 2004, Hayashi PLA 2006



is better than

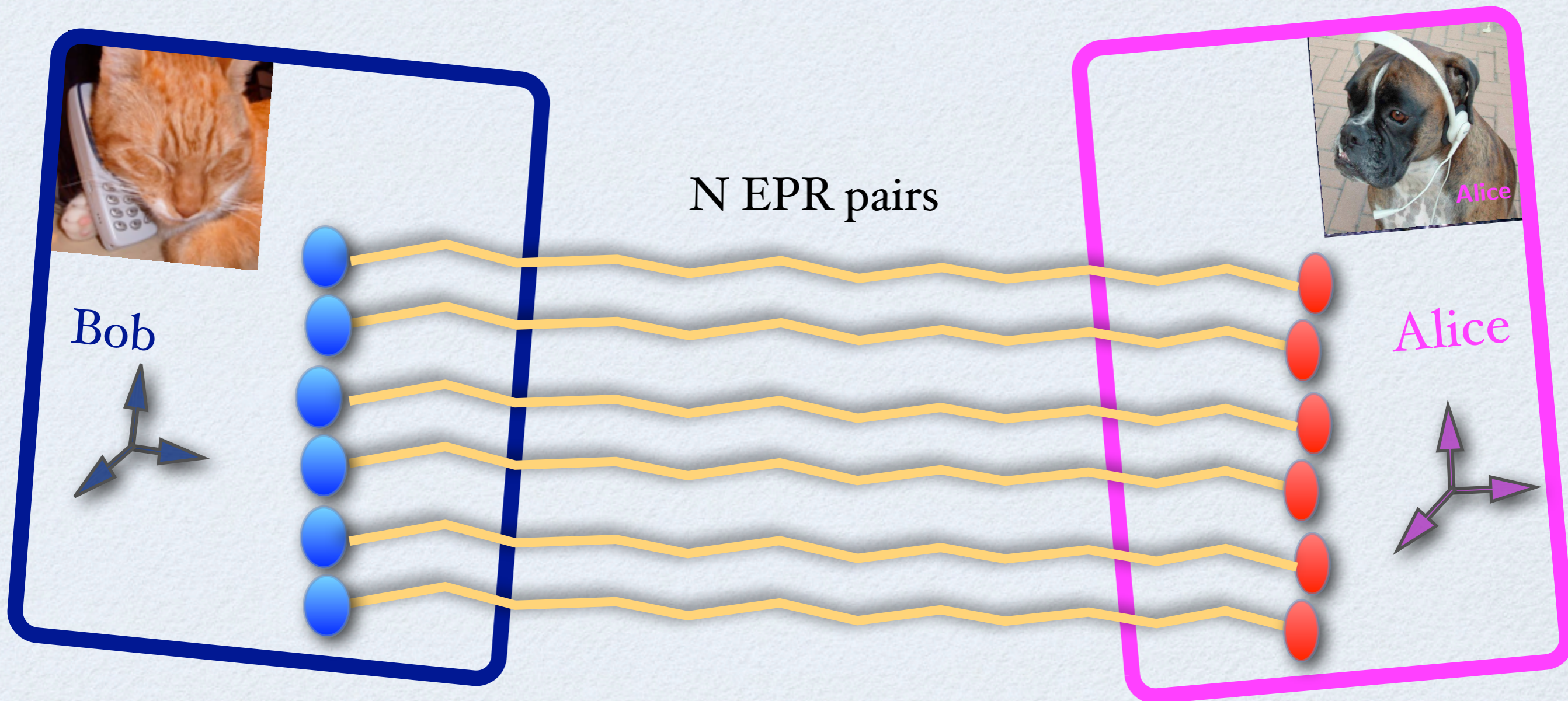


entangling  $N$  particles reduces the error by  $N^2$  (instead of  $N$ )

# The two-party scenario

Rudolph 1999 arXiv

Alice measures her spins along the z-axis, Bob tries to find out the direction.



Scaling up Gisin-Popescu result?

# ...not much

## Deterministic strategies:

- $O(1/N)$  error with Rudolph's protocol
- $O(1/N)$  error with the optimal protocol

CLASSICAL  
SCALING



## Probabilistic strategies:

- $O(1/N^2)$  error with Rudolph's protocol + postselection
- $O(1/N^2)$  error with the optimal protocol + postselection

probability of success:  $O(2^{-N})$

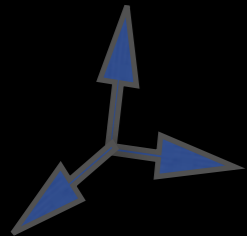
LOW PROBABILITY  
CURSE





USING  
ENTANGLED GYROSCOPES  
OF  
LARGER ANGULAR  
MOMENTUM

Bob



one EPR pair  
of spin- $j$  particles

Alice



Alice

$$|S_j\rangle = \frac{1}{\sqrt{2j+1}} \sum_{m=-j}^j (-1)^m |j, m\rangle \otimes |j, -m\rangle$$
$$J_z |j, m\rangle = m |j, m\rangle$$

optimal state for alignment

# The error

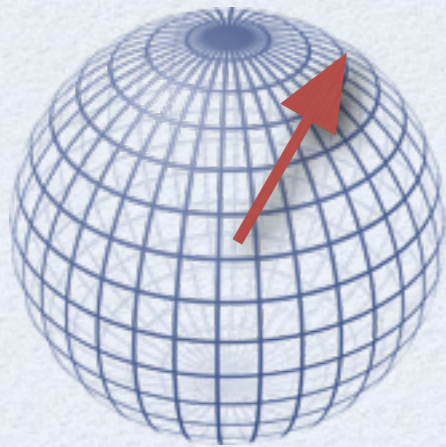
Error

$$\langle d^2 \rangle$$

**CLASSICAL  
SCALING**



No heuristic strategies.



spin  $j$   
gyroscope



length  $j$



solid angle  
 $O(1/j)$

classical gyroscope  
with angular momentum  $j$ ,  
disturbed by a random force  
of fixed intensity,  
error due to precession.

# Full Cartesian frames?

Irreducible error  $\langle d^2 \rangle \geq \frac{4}{3} \quad \forall j$

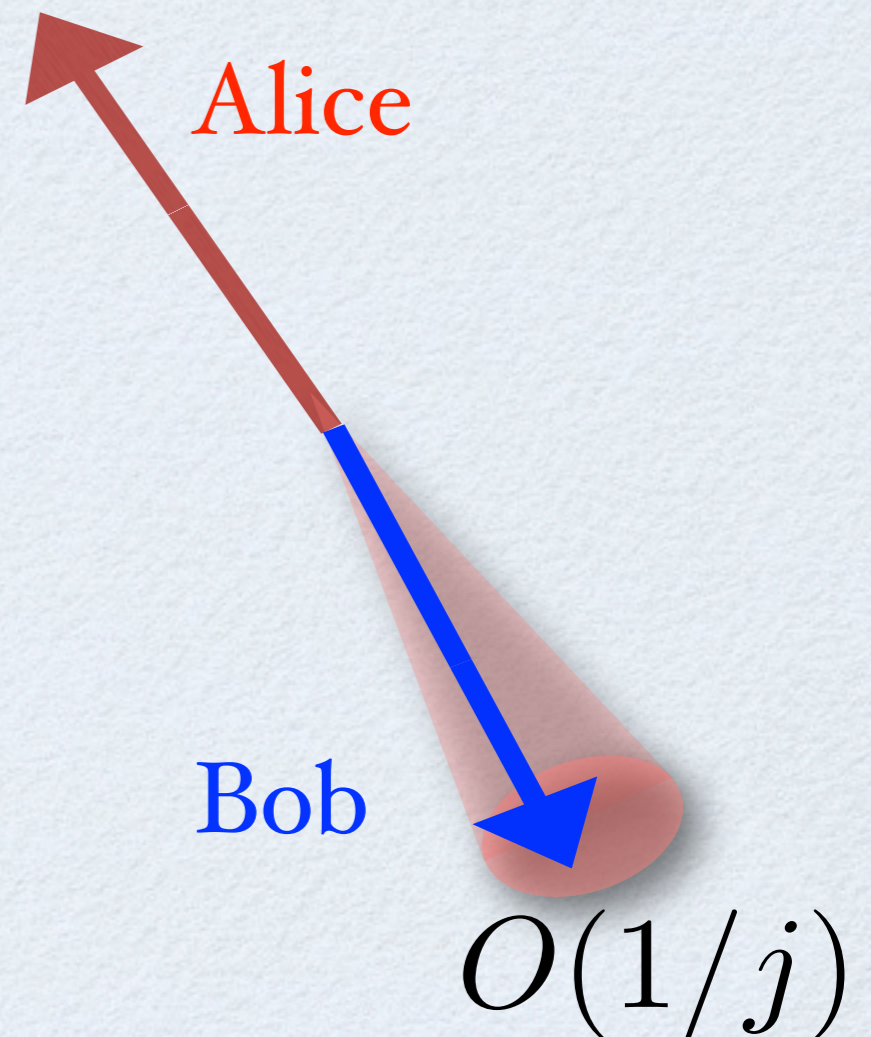
No help from probabilistic



CLASSICAL  
NON-SCALING

Classical

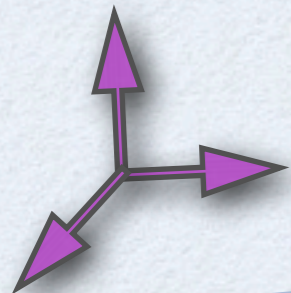
Alice and Bob have a pair of classical gyroscopes pointing in **random, anti-correlated** directions, up to an error  $O(1/j)$



TWO EPR PAIRS  
WITH  
THE ASSISTANCE  
OF  
LOGICAL ENTANGLEMENT

# The teleportation trick

Alice



Bob



EPR pair of spin- $j$  gyroscopes

EPR pair of spin- $j$  gyroscopes

$\log(2j+1)$  logical EPR pairs

$\log(2j+1)$  logical EPR pairs

Bell measurement

$$|S_{j,g}\rangle = (e^{i\varphi} \mathbf{n} \cdot \mathbf{j} \otimes I) |S_j\rangle$$

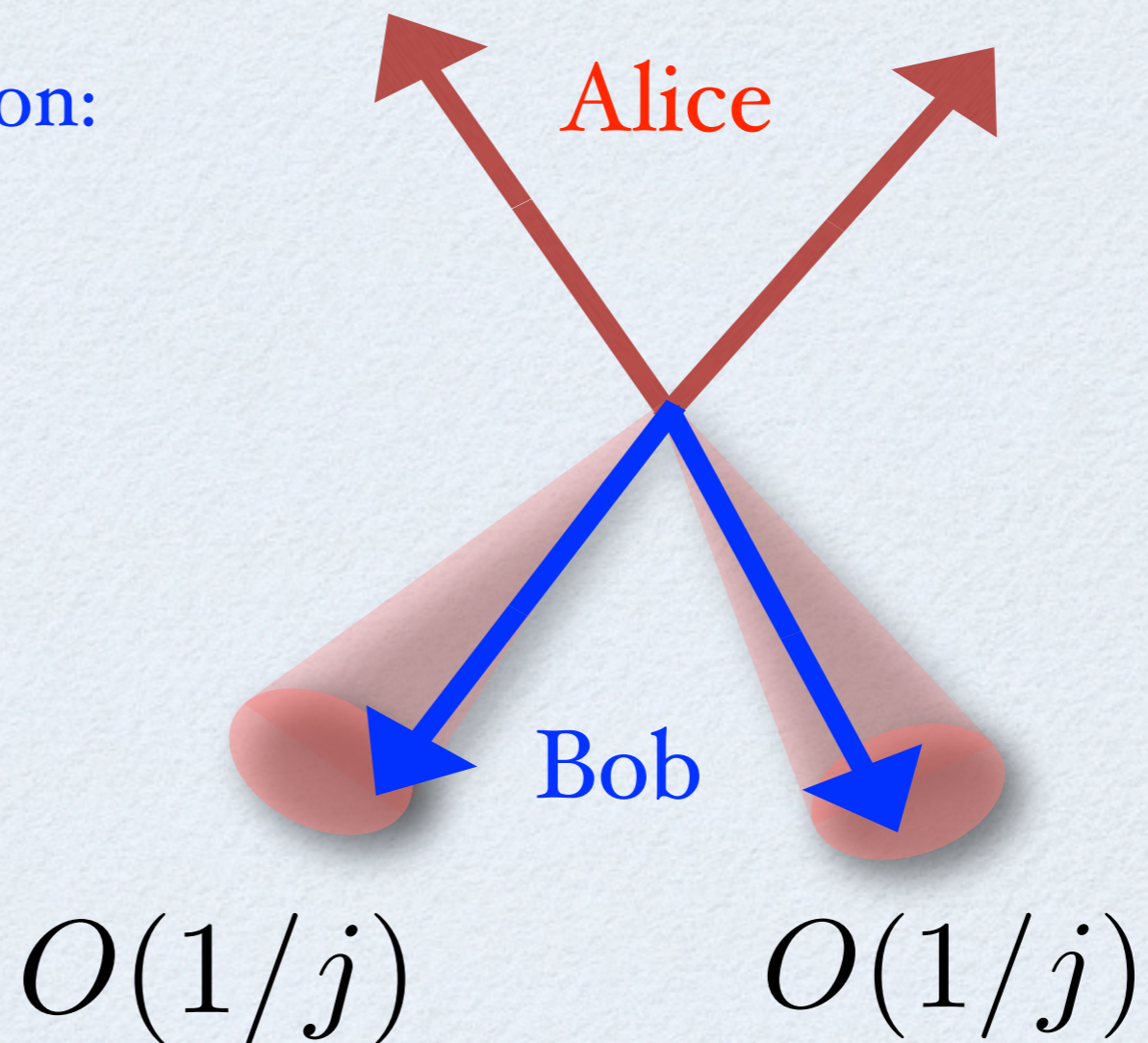
Correction operation

**MORE EPR PAIRS**

# Error for the full Cartesian frame

$$\langle d^2 \rangle = \frac{2}{3j} + O\left(\frac{\log j}{j^2}\right)$$

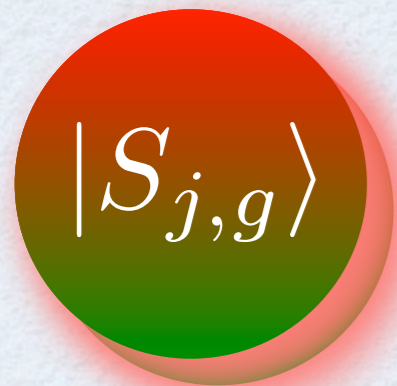
Again, classical explanation:





# Probabilistic strategies

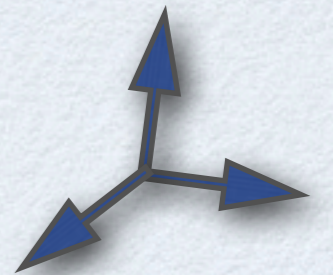
Bob's lab



Filter



Unfavorable  
case

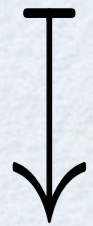


Suppose that Bob uses a probabilistic filter,  
with two outcomes “yes” and “no”

# The yes case

The filter implements the transformation

$$|S_{j,g}\rangle^{\otimes 2} = \bigoplus_{k=0}^{2j} \sqrt{\frac{2k+1}{2j+1}} |S_{k,g}\rangle$$



$$|\Phi_{j,g}^{\text{yes}}\rangle = \bigoplus_{k=0}^{2j} \frac{\sin \frac{\pi(k+1)}{2(j+1)}}{\sqrt{j+1}} |S_{k,g}\rangle$$

= optimal state for transmitting  
a Cartesian frame (GC, D'Ariano,  
Perinotti, Sacchi PRL 2004,  
Bagan et al PRA 2004,  
Hayashi PLA 2006)

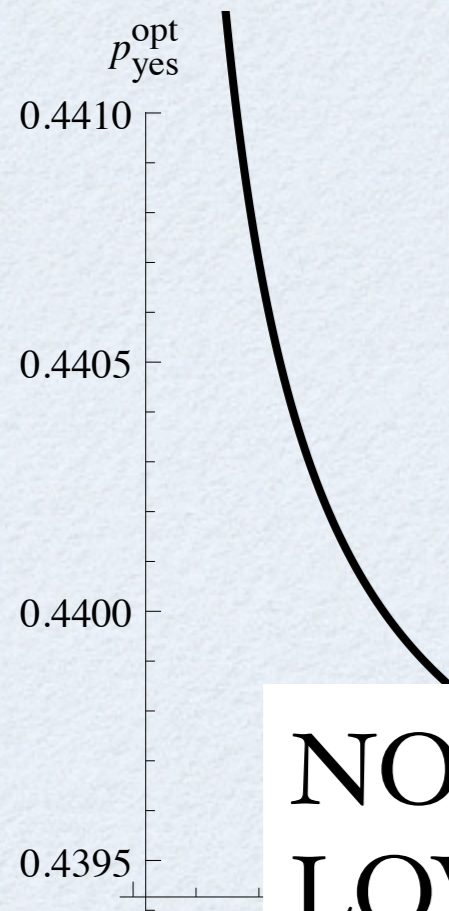
$$\Rightarrow \langle d^2 \rangle = \frac{\pi^2}{6j^2} + O\left(\frac{1}{j^3}\right)$$



# Probabilistic super-activation

In the classical model, Alice and Bob cannot align their axes with an error smaller than  $O(1/j)$

“Probabilistic super-activation” ...what is the probability of seeing it?



$$p_{\text{yes}}^{\text{opt}} = \min_k \frac{(2k + 1)(j + 1)}{(2j + 1)^2 \sin \left[ \frac{\pi(k+1)}{2(j+1)} \right]^2}$$

→ 43.9% for  $j \rightarrow \infty$ !

NO  
LOW PROBABILITY CURSE



# The no case

In the unfavorable instance, the error is  $\langle d^2 \rangle \approx \frac{1.189}{j}$

Of course, no gain on average

$$\langle d^2 \rangle_{\text{average}} \approx p_{\text{yes}}^{\text{opt}} \frac{\pi^2}{6j^2} + (1 - p_{\text{yes}}^{\text{opt}}) \frac{1.189}{j}$$

$$\approx (1 - p_{\text{yes}}^{\text{opt}}) \frac{1.189}{j}$$

$$\approx \frac{0.668}{j} \geq \frac{0.666}{j} = \frac{2}{3j}$$

- BUT
- 1) almost same average performance
  - 2) the quadratic improvement is heralded

MORE  
EPR PAIRS

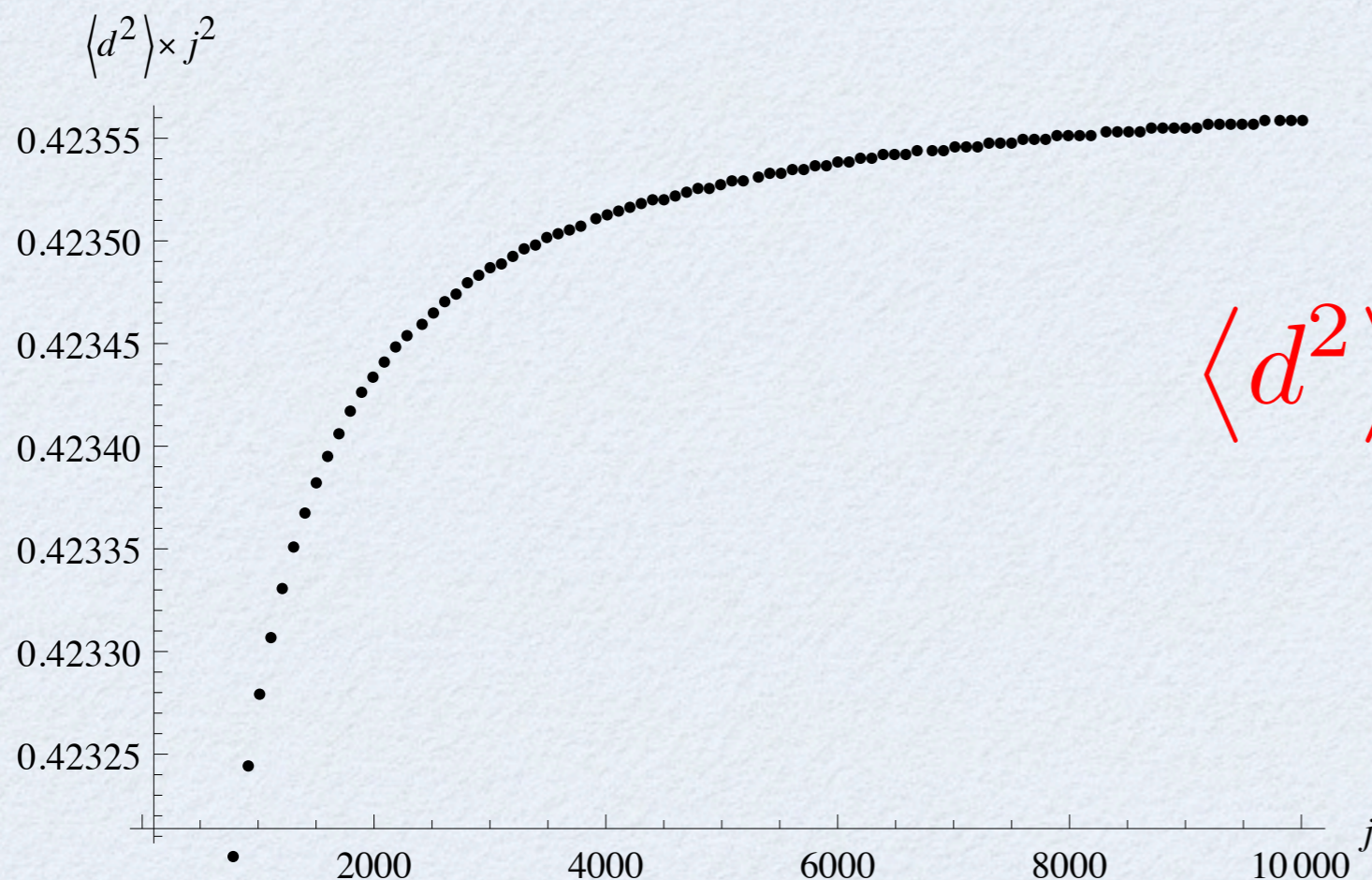
# Deterministic super-activation

$2N$  EPR pairs  $\longrightarrow$  Heisenberg scaling with probability

$$p_{\text{yes}} \geq 1 - (0,561)^N \quad (\text{brute force repetition})$$

In fact, the optimal strategy can do much better:

for **4 pairs** the Heisenberg scaling is achieved **deterministically**



$$\langle d^2 \rangle = \frac{11 \ln 2}{18j^2} + O\left(\frac{1}{j^3}\right)$$

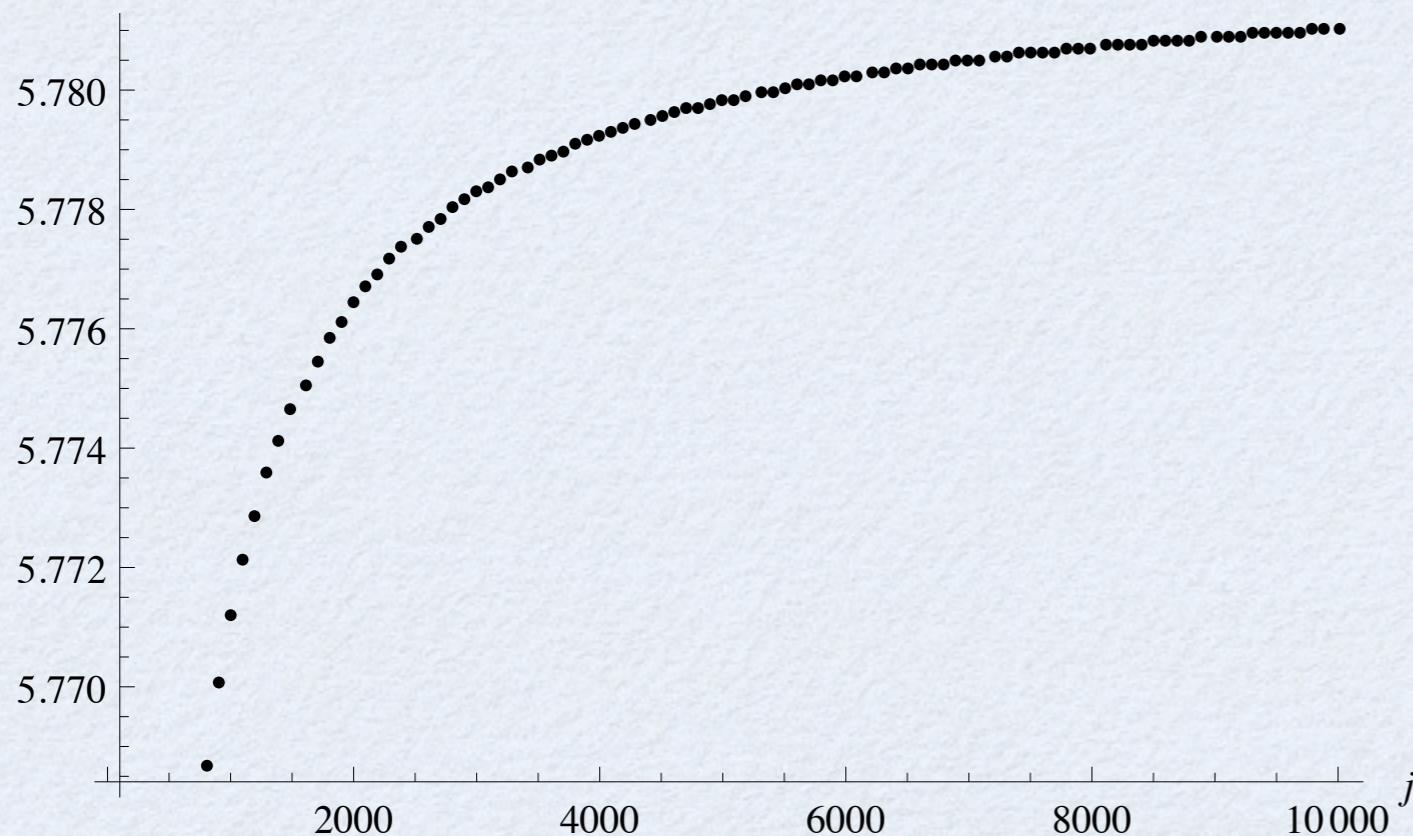
# Quasi-Heisenberg scaling for 3 EPR pairs

4 EPR pairs are the minimum to achieve the Heisenberg scaling deterministically.

Still, if Alice and Bob have only 3 pairs, they can still beat the classical scaling.

The optimal strategy yields:

$\langle d^2 \rangle \times 8j^2 - \ln j$



$$\langle d^2 \rangle = \frac{\ln j}{8j^2} + O\left(\frac{1}{j^2}\right)$$

(quasi-Heisenberg scaling)

A CAUTIONARY TALE  
ABOUT  
THE  
QUANTUM CRAMÈR-RAO  
BOUND  
IN  
NON-ASYMPTOTIC  
SCENARIOS



# Cramèr-Rao bound

State parametrization  $|S_{j,\boldsymbol{\theta}}\rangle = (e^{-i\boldsymbol{\theta}\cdot\mathbf{j}} \otimes I) |S_j\rangle$      $\boldsymbol{\theta} = \begin{pmatrix} \theta_x \\ \theta_y \\ \theta_z \end{pmatrix}$

Quantum Fisher Information matrix  $F_Q = \frac{4j(j+1)}{3} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$

Quantum Cramèr-Rao bound  $V_{\boldsymbol{\theta}} \geq F_Q^{-1} = \frac{3}{4j(j+1)} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$

achievable in the asymptotic regime of large number of copies

# Caveat

A naive application of the CRB would promise Heisenberg scaling of the error:

$$V_{\theta} \approx O\left(\frac{1}{j^2}\right) \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \implies \langle d^2 \rangle \approx O\left(\frac{1}{j^2}\right)$$

in contradiction with the analytical results for 1,2 and 3 pairs

For spin- $j$  singlets,  
the CRB is not achievable in the finite-copy regime

# Asymptotic achievability of CRB

To achieve the CRB, the number of copies must be large.  
But how large?

Hopefully not large compared to  $j^2$  ...

Covariance matrix for the optimal measurement:

$$V_{\boldsymbol{\theta},n}^{\text{opt}} = \frac{3}{4nj(j+1)} I + O\left(\max\left\{n^{-3/2}j^{-3}, n^{-2}j^{-2}\right\}\right)$$

CRB achieved whenever  $n \gg 1$ , uniformly in  $j$

# CONCLUSIONS

# Conclusions

Super-activation of quantum gyroscopes:

- 1 EPR pair  $\longrightarrow$  no scaling with the size
- 2 EPR pairs  $\longrightarrow$  Heisenberg scaling with non-vanishing prob.
- 3 EPR pairs  $\longrightarrow$  quasi-Heisenberg
- $\geq 4$  EPR pairs  $\longrightarrow$  Heisenberg with certainty

The moral:

- not having a pre-defined direction helps
- two spin- $j$  particles are more useful than a single spin- $(2j)$  particle
- logical qubits help
- ...be careful using the quantum CRB in non-asymptotic scenarios!

THANK YOU

FOR YOUR

ATTENTION!