

The Most Likely Evolution of Diffusing and Vanishing Particles in the Spirit of Erwin Schrödinger:

Constructing bridges with unbalanced marginals

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joint work with

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Dynamics and Discretization: PDEs, Sampling, and Optimization
Simons Institute, Berkeley

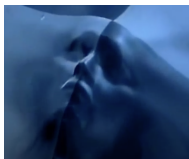
Oct 25 – Oct 29, 2021

Outline of the talk

- Some context/motivation: interpolation of distributions (bridges)
- Optimal mass transport and Schrödinger's bridge problem
- Diffusing and vanishing Particles in the spirit of Schrödinger
 - Bridges between unbalanced marginals

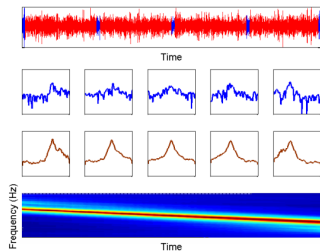
Interpolation of distributions

aka Morphing

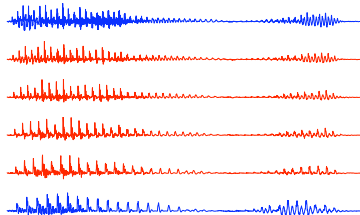


Interpolation of distributions

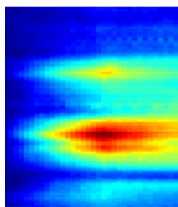
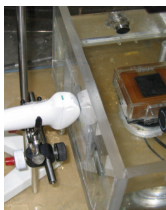
Time-series analysis



Doppler frequency tracking

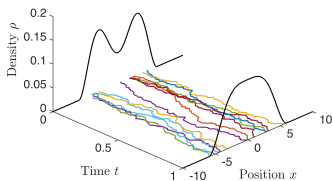
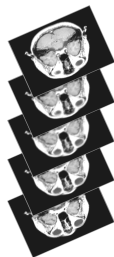
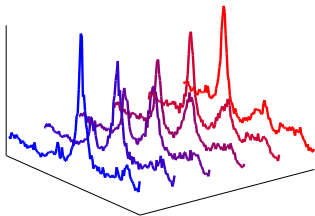


Voice morphing



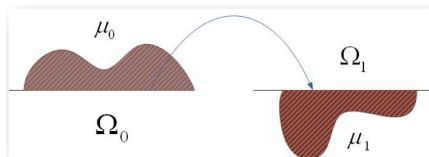
Noninvasive temperature sensing - temperature field

Interpolation of distributions – unbalanced marginals





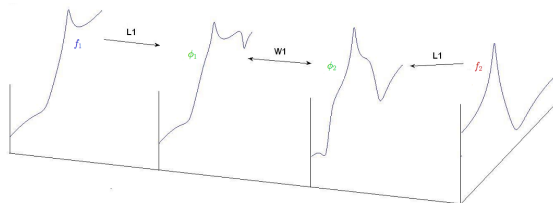
G. Monge (1871)



L. Kantorovich (1942)

and then McCann, Gangbo, Brenier,
Benamou, Ambrosio, ... (1990's on)
Rachev-Ruschendorf, Villani, ...

Unbalanced marginals - before 2010



$$\mu_0(\Omega_0) \neq \mu_1(\Omega_1)$$

$$\begin{aligned} d_{\text{mixed}, \kappa}(\mu_0, \mu_1) &= \inf_{\hat{\mu}_0, \hat{\mu}_1} d_W(\hat{\mu}_0, \hat{\mu}_1) + \kappa \sum_{i=0}^1 \|\hat{\mu}_i - \mu_i\|_{\text{TV}} \\ &= \sup_f \left\{ \int f d(\mu_0 - \mu_1) \mid \|f\|_{\text{Lip}} \leq 1, \|f\|_{\infty} \leq \kappa \right\} \end{aligned}$$

$\hat{\mu}_0, \hat{\mu}_1$: noise-free measures

$\mu_i - \hat{\mu}_i$: noise components

e.g., see G-Karlssohn-Takyar 2009

Unbalanced marginals - post 2010

$$\inf_{\rho, \nu, \tilde{\rho}_1} \int_0^1 \int_{\mathbb{R}^m} \rho(t, x) \|\nu\|^2 dx dt + \alpha \int_{\mathbb{R}^m} (\rho_1(x) - \tilde{\rho}_1(x))^2 dx,$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \nu) = 0, \quad \rho(0, \cdot) = \rho_0(\cdot), \quad \rho(1, \cdot) = \tilde{\rho}_1(\cdot) \text{ (not necessarily } = \rho_1(\cdot)).$$

$$\inf_{\rho, \nu, s} \int_0^1 \int_{\mathbb{R}^m} \left\{ \rho(t, x) \|\nu\|^2 + \alpha s(t, x)^2 \right\} dx dt,$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \nu) = s, \quad \rho(0, \cdot) = \rho_0(\cdot), \quad \rho(1, \cdot) = \rho_1(\cdot).$$

$$\inf_{\rho, \nu, r} \int_0^1 \int_{\mathbb{R}^m} \left\{ \rho(t, x) \|\nu\|^2 + \alpha \frac{s^2}{\rho(t, x)} \right\} dx dt$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \nu) = s, \quad \rho(0, \cdot) = \rho_0(\cdot), \quad \rho(1, \cdot) = \rho_1(\cdot).$$

see Liero-Mielke-Savaré (arxiv.org/pdf/1508.07941),
Peyré-Cuturi 2020, also Chen-G-Tannenbaum 2018

Optimal Mass Transport regularization: Schrödinger's Bridge Problem (SBP)

Balanced marginals for now

A problem in large-deviations that leads to:

$$\inf_{(\rho, \nu)} \int_{\mathbb{R}^n} \int_0^1 \rho(t, x) \|v(t, x)\|^2 dt dx,$$
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\nu \rho) = \frac{1}{2} \Delta \rho$$
$$\rho(0, x) = \rho_0(x), \quad \rho(1, y) = \rho_1(y)$$

And a fluid-dynamic, time-symmetric, formulation:

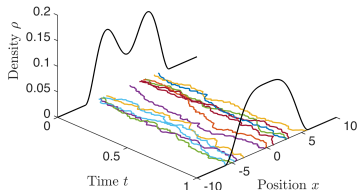
$$\inf_{(\rho, \nu)} \int_{\mathbb{R}^n} \int_0^1 \left[\|v(t, x)\|^2 + \left\| \frac{1}{2} \nabla \log \rho(t, x) \right\|^2 \right] \rho(t, x) dt dx,$$
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \nu) = 0,$$
$$\rho(0, x) = \rho_0(x), \quad \rho(1, y) = \rho_1(y).$$

Blaquière, Dai Pra, Pavon-Wakolbinger, Filliger-Hongler-Streit, Mikami,
Thieulien, Leonard, Chen-G-Pavon

Schrödinger's Bridge Problem (SBP)



Erwin Schrödinger
Schrödinger bridges 1931/32

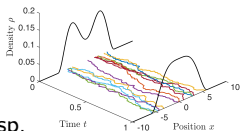


~ Nelson's stochastic mechanics

Schrödinger's Bridge Problem (SBP)

Consider:

- Cloud of N independent Brownian particles (N large)
- empirical distr. $\rho_0(x)dx$ and $\rho_1(y)dy$ at $t = 0$ and $t = 1$, resp.
- ρ_0 and ρ_1 not compatible with transition mechanism



$$\rho_1(y) \neq \int_0^1 p(0, x, 1, y) \rho_0(x) dx,$$

where

$$p(s, y, t, x) = [2\pi(t-s)]^{-\frac{n}{2}} \exp\left[-\frac{|x-y|^2}{2(t-s)}\right], \quad s < t$$

Particles have been transported in an unlikely way

Schrödinger (1931): Of the many unlikely ways in which this could have happened, which one is the most likely?

Schrödinger's Bridge Problem (SBP)

Large deviations formulation

$$\min_P H(P|R) = \min_P E_P \left[\log \frac{dP}{dR} \right]$$

over $P \in \{\text{distributions on paths with marginals } \rho_0, \rho_1\}$;

$H(\cdot|\cdot)$ is the relative entropy

R reference Wiener measure

Föllmer 1988: SBP is a large deviations problem of the empirical distribution on paths \equiv maximum entropy problem via Sanov's thm

Connection to stochastic control & OMT:

For prior the law of a diffusion: $dX = vdt + dB$, Girsanov's thm:

$$E_Q \left[\log \frac{dQ}{dR} \right] = E_Q \left[\frac{1}{2} \int_0^1 \|v\|^2 ds \right]$$

Schrödinger's Bridge Problem (SBP)

Stochastic control formulation & structure of solutions

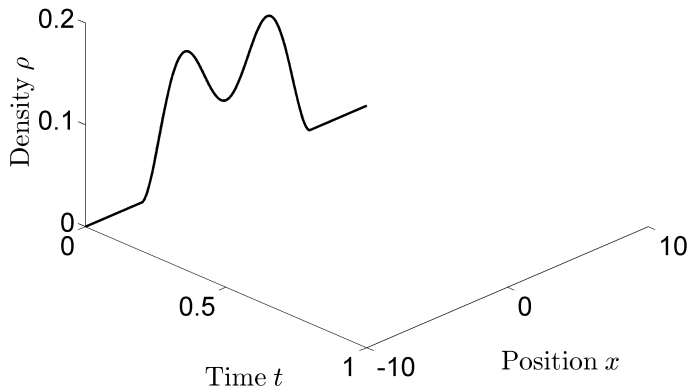
- Girsanov's thm gives:

$$\inf_{(\rho, \nu)} \int_{\mathbb{R}^n} \int_0^1 \|v(t, x)\|^2 \rho(t, x) dt dx,$$
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\nu \rho) = \frac{1}{2} \Delta \rho$$
$$\rho(0, x) = \rho_0(x), \quad \rho(1, y) = \rho_1(y)$$

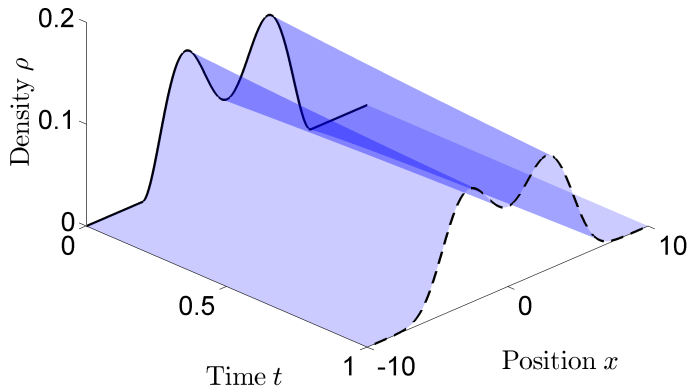
- $\min H(P|R) \Rightarrow \rho(t, x) = \varphi(t, x) \hat{\varphi}(t, x)$ (t -time marginal of P)
where φ and $\hat{\varphi}$ solve the **Schrödinger's system**:

$$\varphi(t, x) = \int \rho(t, x, 1, y) \varphi(1, y) dy, \quad \varphi(0, x) \hat{\varphi}(x, 0) = \rho_0(x)$$
$$\hat{\varphi}(t, x) = \int \rho(0, y, t, x) \hat{\varphi}(0, y) dy, \quad \varphi(1, x) \hat{\varphi}(1, x) = \rho_1(x).$$

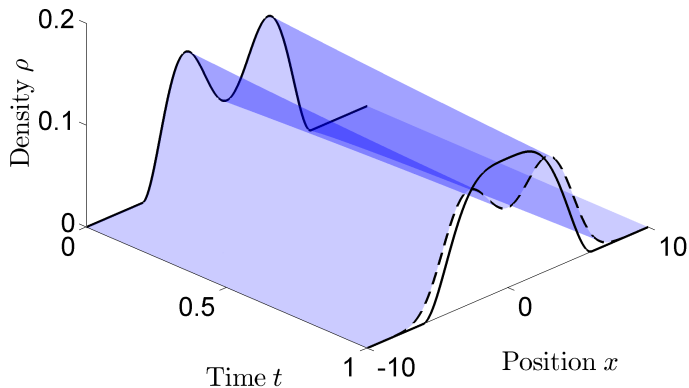
SBP schematic - marginal



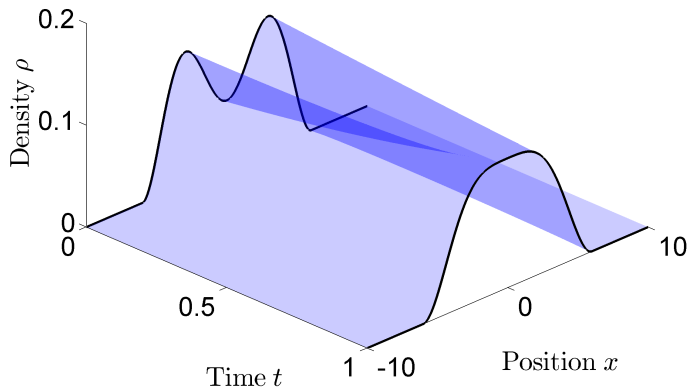
SBP schematic - prior



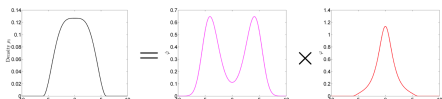
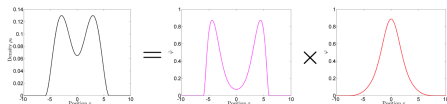
SBP schematic - prior vs. mismatched end-point marginal



SBP schematic - Schrödinger bridge



Schrödinger system

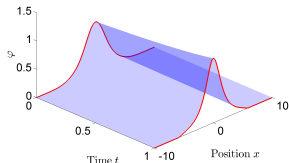
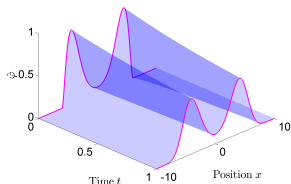
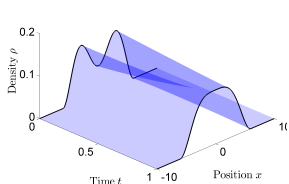


$$-\frac{\partial \varphi}{\partial t}(t, x) = \frac{1}{2} \Delta \varphi(t, x)$$

$$\frac{\partial \hat{\varphi}}{\partial t}(t, x) = \frac{1}{2} \Delta \hat{\varphi}(t, x)$$

$$\varphi(0, x) \hat{\varphi}(0, x) = \rho_0(x)$$

$$\varphi(1, x) \hat{\varphi}(1, x) = \rho_1(x)$$



For $dX_t = b(t, X_t)dt + \sigma(t, X_t)dW_t$ (\star)

SBP theory outline

$a(t, X) := \sigma(t, X)\sigma(t, X)' > 0$

Notation:

R : “prior” law of (\star) on paths

R_t, R_{st} : marginals at times t , and jointly t, s

$R^{xy}(\cdot)$ law conditioned on $X_0 = x, X_1 = y$

disintegration of measure $R(\cdot) = \int_x \int_y R^{xy}(\cdot)R_{01}(dxdy)$

SBP: Find

$$P^* = \arg \min_P \{H(P|R) \mid P_0 = \rho_0, P_1 = \rho_1\}$$

$$H(P|R) = H(P_{01}|R_{01}) + \int H(P^{xy}|R^{xy})P_{01}(dxdy)$$

Static SBP: Find

$$P_{01}^* = \arg \min_{P_{01}} \{H(P_{01}|R_{01}) \mid P_0 = \rho_0, P_1 = \rho_1\}$$

Relation static-dynamic SBP:

$$P^*(\cdot) = \int_x \int_y R^{xy}(\cdot) P_{01}^*(dxdy)$$

Solution:

Under mild/natural assumptions, $\exists f, g$ so that:

$P_{01}^* = f(X_0)g(X_1)R_{01}$. These are solutions of the **Schrödinger system**

$$\frac{d\rho_0}{dR_0}(x) = f(x)R(g(X_1) | X_0 = x),$$

$$\frac{d\rho_1}{dR_1}(y) = g(y)R(f(X_0) | X_1 = y).$$

$$P_{01}^* = f(X_0)g(X_1)R_{01} \Leftrightarrow P^* = f(X_0)g(X_1)R.$$

Solution:for $\hat{\varphi}(0, x) := f(x)R_0(x)$, $\varphi(1, y) := g(y)$

$$\partial_t \hat{\varphi} = -\nabla \cdot (b\hat{\varphi}) + \frac{1}{2} \sum_{i,j=1}^n \frac{\partial^2 (a_{ij}\hat{\varphi})}{\partial x_i \partial x_j}$$

$$\partial_t \varphi = -b \cdot \nabla \varphi - \frac{1}{2} \sum_{i,j=1}^n a_{ij} \frac{\partial^2 \varphi}{\partial x_i \partial x_j}$$

$$\rho_0 = \varphi(0, \cdot) \hat{\varphi}(0, \cdot)$$

$$\rho_1 = \varphi(1, \cdot) \hat{\varphi}(1, \cdot).$$

Then, $P_t^* = \rho(t, \cdot) = \phi(t, \cdot) \hat{\phi}(t, \cdot)$ (t -time marginal) of the law of

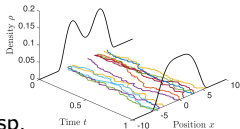
$$dX_t = (b(t, X_t) + a(t, X_t) \nabla \log \varphi(t, X_t))dt + \sigma(t, X_t)dW_t$$

Schrödinger's Bridge with losses

most likely evolution of diffusing and vanishing particles

Consider:

- Cloud of N “tracer” particles (N large)
- empirical distr. $\rho_0(x)dx$ and $\rho_1(y)dy$ at $t = 0$ and $t = 1$, resp.
- ρ_0 and ρ_1 not compatible with transition mechanism



$$\rho_1(y) \neq \int_0^1 p(t_0, x, t_1, y) \rho_0(x) dx,$$

Besides having been transported in an unlikely way,

the particles remain in suspension for a duration of time, and thus,

at $t = 1$ a random portion of the particles have been lost (sunk), and $\int \rho_1 < \int \rho_0$

Question - in the spirit of Schrödinger:

What is the most likely evolution that accounts for losses?

Stochastic transport with losses – Prior:

$$dX_t = b(t, X_t)dt + \sigma(t, X_t)dW_t \quad (**)$$

with killing rate $V(t, x)$

State space: $\mathcal{X} = \mathbb{R}^n \cup \{c\}$ with c a “coffin state”

Paths $\Omega = D([0, 1], \mathcal{X})$ càdlàg

$(X_t$ on \mathbb{R}^n with killing) \equiv $(\mathbf{X}_t$ on \mathcal{X} with a law on $\mathcal{P}(\Omega)$)

p_0, p_1 natural augmentation of ρ_0, ρ_1 so that $p_0, p_1 \in \mathcal{P}(\mathcal{X})$

i.e., assuming $\int \rho_1 = 1$, **set** $p_0 = (\rho_0(\cdot), 0)$, and $p_1 = (\rho_1(\cdot), 1 - \int \rho_1)$

$$\mathbf{P}^* := \arg \min_{\mathbf{P} \in \mathcal{P}(\Omega)} \{H(\mathbf{P} \mid \mathbf{R}) \mid \mathbf{P}_0 = p_0, \mathbf{P}_1 = p_1\}.$$

Schrödinger Bridge with losses

unbalanced SBP – $\int \rho_0 > \int \rho_1$

Prior: Fokker-Planck equation for a diffusion with killing rate $V(t, x)$

$$\partial_t R_t + \nabla \cdot (bR_t) + VR_t = \frac{1}{2} \sum_{i,j=1}^n \frac{\partial^2 (a_{ij} R_t)}{\partial x_i \partial x_j}.$$

SB with losses: “new ϕ ” = $(\phi(t, \cdot), \psi(t))$ on \mathcal{X} , same for “new ψ ,” via the **Schrödinger system:**

$$\partial_t \hat{\varphi} = -\nabla \cdot (b\hat{\varphi}) - V\hat{\varphi} + \frac{1}{2} \sum_{i,j=1}^n \frac{\partial^2 (a_{ij} \hat{\varphi})}{\partial x_i \partial x_j}$$

$$\frac{d\hat{\psi}}{dt} = \int V\hat{\varphi}(t, x) dx$$

$$\partial_t \varphi = -b \cdot \nabla \varphi + V\varphi - \frac{1}{2} \sum_{i,j=1}^n a_{ij} \frac{\partial^2 \varphi}{\partial x_i \partial x_j} - V\psi$$

$$\frac{d\psi}{dt} = 0$$

with b.c.

$$\rho_0 = \varphi(0, \cdot) \hat{\varphi}(0, \cdot)$$

$$\rho_1 = \varphi(1, \cdot) \hat{\varphi}(1, \cdot)$$

$$\hat{\psi}(0) = 0$$

$$\psi(1) \hat{\psi}(1) = 1 - \int \rho_1.$$

$$\Rightarrow \mathbf{P}^* = f(\mathbf{X}_0)g(\mathbf{X}_1)\mathbf{R}$$

$$\mathbf{P}^* = (P_t^*, q_t^*), \mathbf{R} = (R_t, s_t)$$

Schrödinger Bridge with losses – dynamic formulation

\mathbf{P}^* is the law of a diffusion

$$dX_t = (b(t, X_t) + a(t, X_t)\nabla \log \varphi(t, X_t))dt + \sigma(t, X_t)dW_t$$

with killing rate $\psi V/\varphi$, and Fokker-Planck equation

$$\partial_t P_t + \nabla \cdot ((b + a\nabla \log \varphi)P_t) = \frac{1}{2} \sum_{i,j=1}^n \frac{\partial^2 (a_{ij}P_t)}{\partial x_i \partial x_j} - \frac{\psi}{\varphi} VP_t.$$

mass $q(t)$ on c :

$$\frac{dq_t}{dt} = \psi(t) \int V \hat{\varphi}(t, x) dx = \int \frac{\psi}{\varphi} VP_t dx$$

Schrödinger Bridge with losses – fluid dynamic formulation

Contrast with original SB the **added terms**:

$$\min_{P_t(\cdot), u(t, \cdot)} \int_0^1 \int_{\mathbb{R}^n} \left[\frac{1}{2} \|u(t, x)\|^2 P_t + (\alpha \log \alpha - \alpha + 1) VP_t \right] dx dt$$
$$\partial_t P_t + \nabla \cdot ((b + \sigma u) P_t) + \alpha VP_t - \frac{1}{2} \sum_{i,j=1}^n \frac{\partial^2 (a_{ij} P_t)}{\partial x_i \partial x_j} = 0$$
$$P_0 = \rho_0, \quad P_1 = \rho_1.$$

$$u^*(t, x) = \sigma(t, x)' \nabla \log \varphi(t, x)$$

$$\alpha^*(t, x) = \frac{\psi(t)}{\varphi(t, x)}$$

with marginals:

$$P_t^*(x) = \varphi(t, x) \hat{\varphi}(t, x) \text{ on } \mathbb{R}^n$$

$$q_t = \psi(t) \hat{\psi}(t) \text{ on } \mathfrak{c}$$

SBP on Feynman-Kac reweighed processes

Earlier attempts to “model” losses – Nagasawa, Wakolbinger, Leonard, Chen-G-Pavon, ...

Feynman-Kac reweighing of the prior

$$\hat{R} := \exp\left(-\int_0^1 V(t, X_t) dt\right) R \quad \mapsto \quad \hat{P}^* = f(X_0) \exp\left(-\int_0^1 V(t, X_t) dt\right) g(X_1) R$$

via

$$\hat{P}^* := \min_{P \in \mathcal{P}(\Omega)} \left\{ H(P | \hat{R}) \mid P_0 = \rho_0, P_1 = \hat{\rho}_1 \right\},$$

with $\hat{\rho}_1$ normalized distribution of **survived particles**

- **upside:** simpler Schrödinger system
- **downside:** not physical & inconsistent with Schrödinger's dictum

ρ_0 distribution of all starting particles, ρ_1 surviving particles
starting distribution of survived particles in **not knowable**

no mechanism to update V

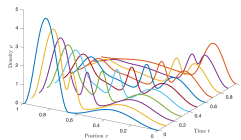
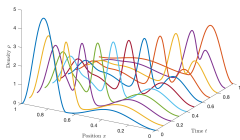
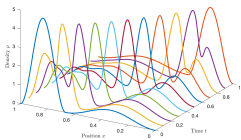
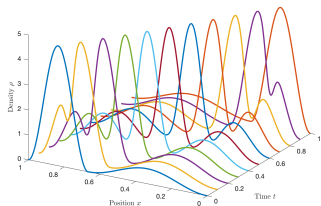
when ρ_1 consistent with prior model and losses in V , $\hat{P}^* \neq \hat{R}$

marginals of \hat{R} and \hat{P}^* have constant mass

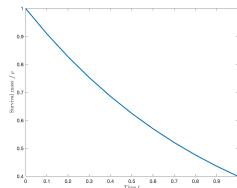
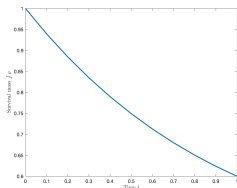
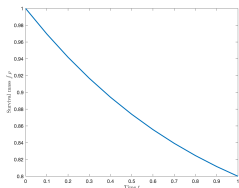
Numerical example

Prior:

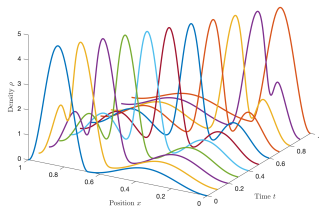
$$dX_t = \sigma dW_t \text{ with killing rate } V(t, x) = 1.$$



Numerical example



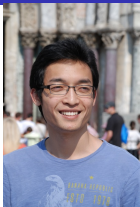
survived mass



reweighted process, regardless of end-point mass

Diffusing and Vanishing Particles in the Spirit of Schrödinger

Bridges with unbalanced marginals



Yongxin Chen



Michele Pavon

Thank you for your attention

CGP arxiv.org/abs/2108.02879