

# Provably efficient machine learning for quantum many-body problems

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Joint work with Richard Kueng, Giacomo Torlai, Victor V. Albert, John Preskill

arXiv:2106.12627, 2021

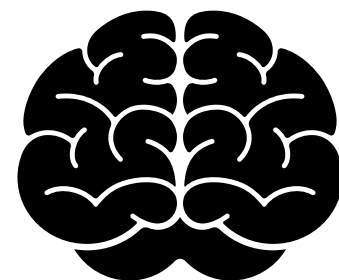


# Motivation

- Machine learning (ML) has received great attention in the quantum community these days.
- Yet, many fundamental questions remain to be answered.

## Classical ML for quantum physics/chemistry


The goal 🎯:  
Solve challenging  
quantum many-body problems  
**better** than  
traditional classical algorithms

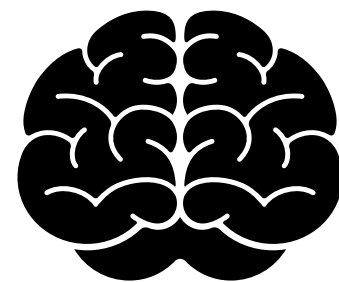


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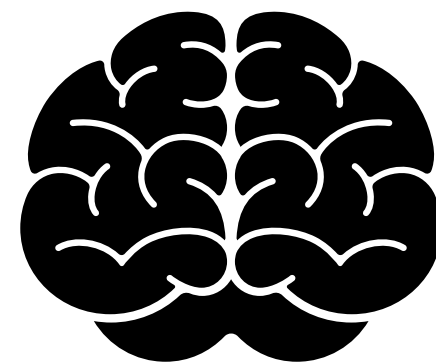
## Classical ML for quantum physics/chemistry

The question :  
How can ML algorithms be more useful  
than non-ML algorithms?



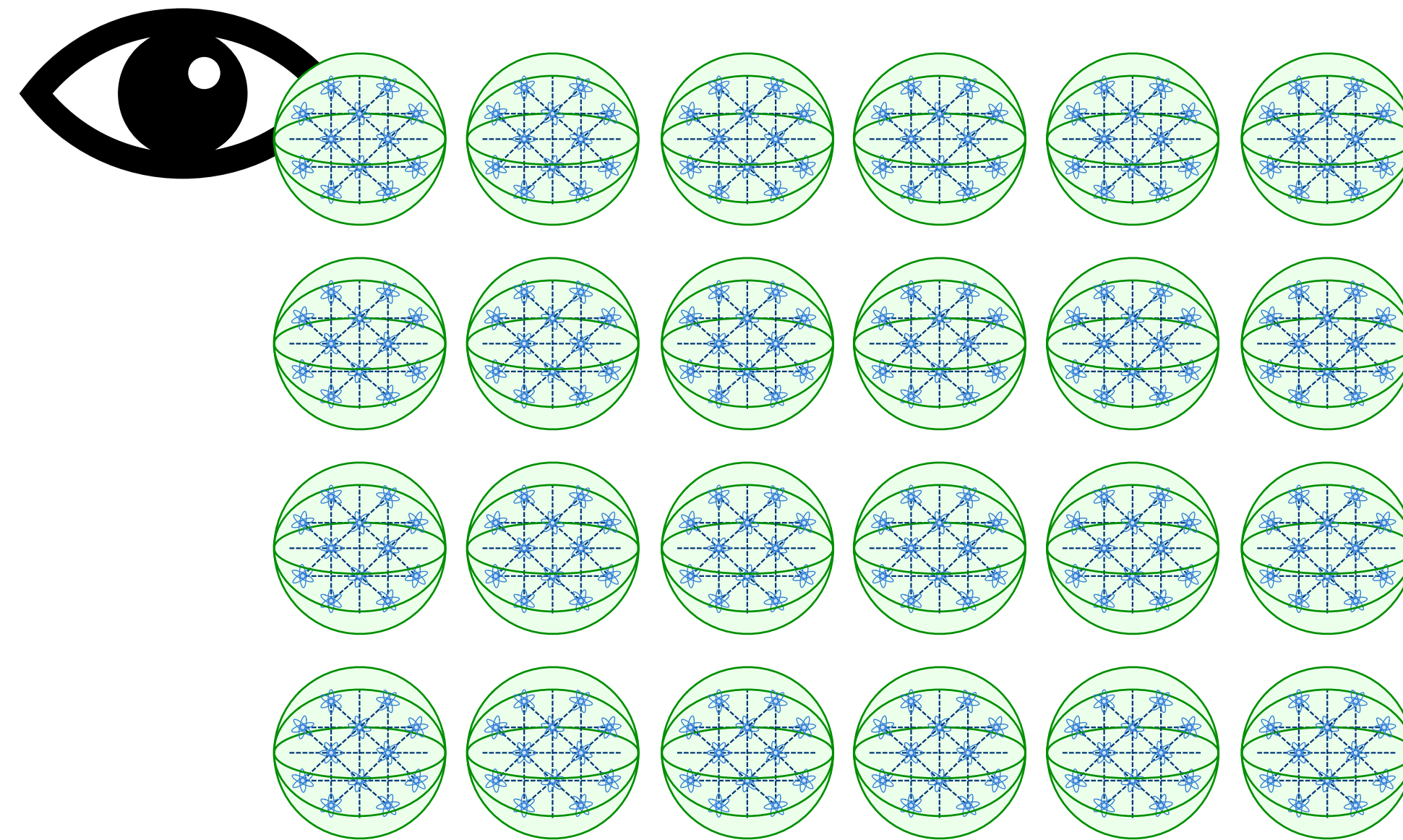
# Outline

- Review on classical shadow formalism
- Training machines to predict ground states  
(theory+experiments)
- Training machines to classify quantum phases of matter  
(theory+experiments)



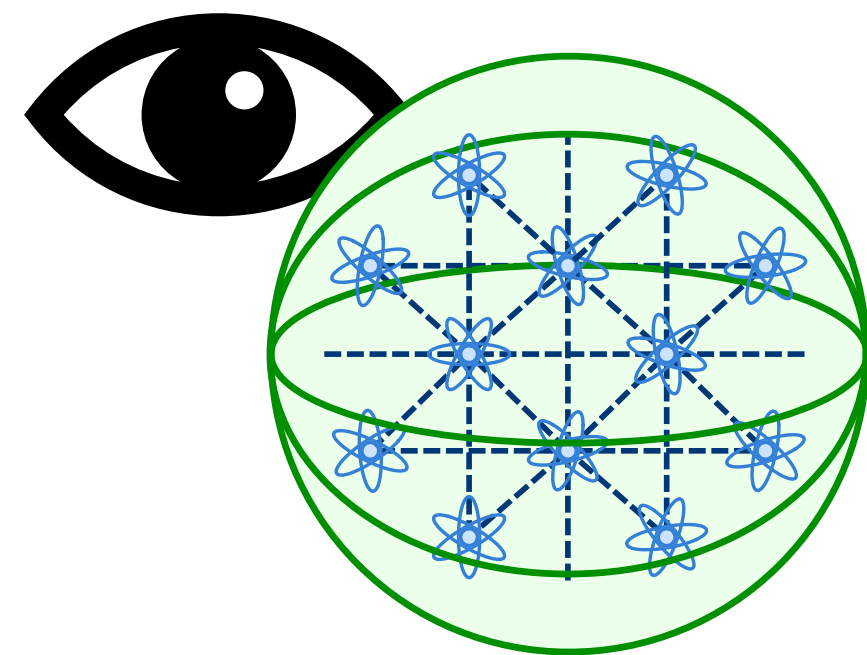
# Classical shadow formalism

- How can classical machines "see" quantum many-body systems?

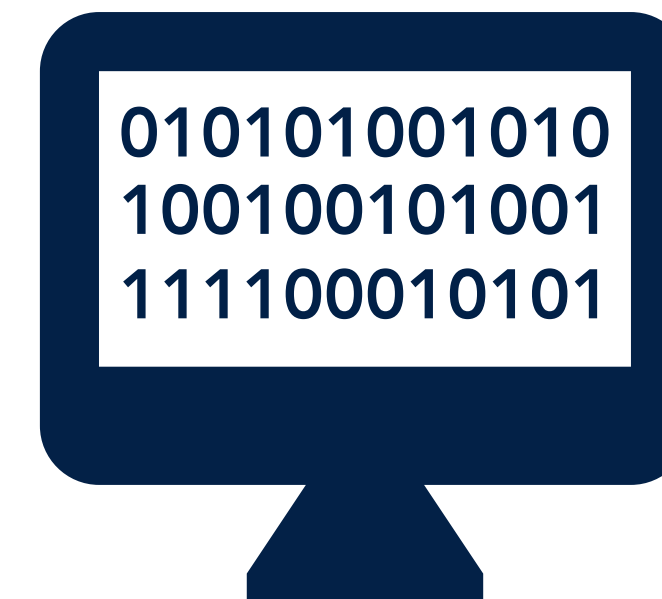
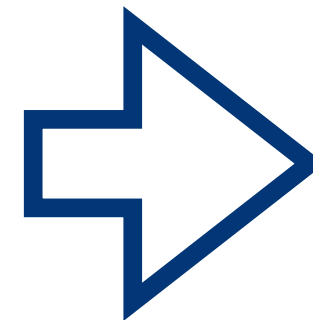


# Classical shadow formalism

- What do we mean by “seeing” a quantum system?
- Converting the quantum system to a classical form that accurately captures many properties of the quantum system.



**Unknown  
Quantum System**



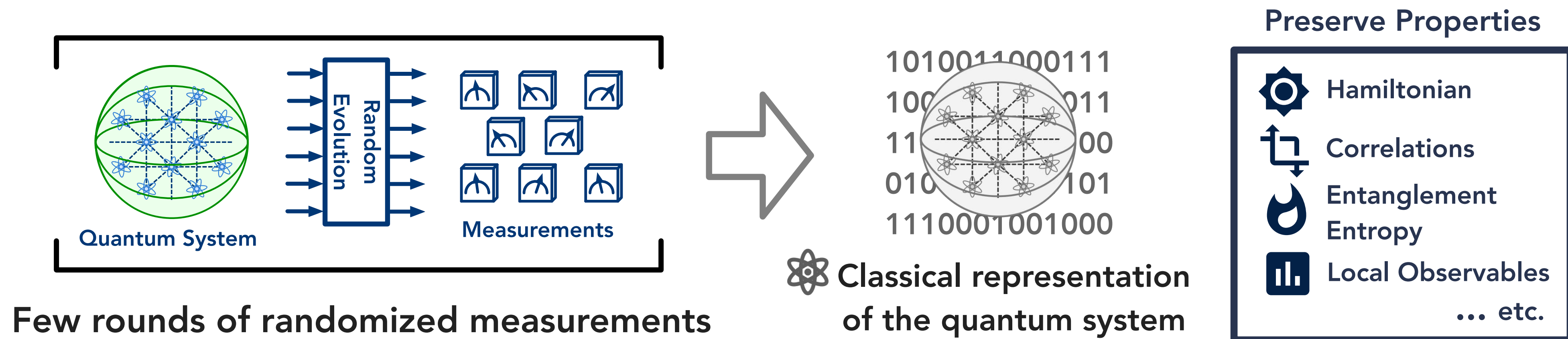
**Efficient  
Classical Representation**

# Classical shadow with randomized Pauli measurements

- After  $T$  randomized Pauli measurements, an  $n$ -qubit system  $\rho$  yields a classical shadow

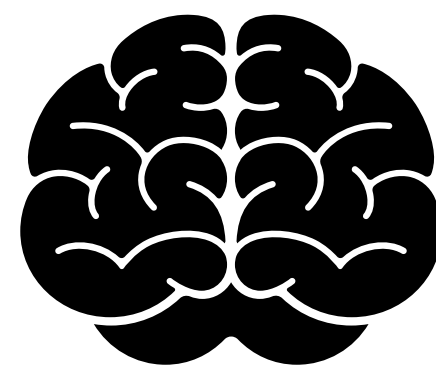
$$\sigma_T(\rho) = \frac{1}{T} \sum_{t=1}^T \sigma_1^{(t)} \otimes \dots \otimes \sigma_n^{(t)}, \text{ where } \sigma_i^{(t)} \in \mathbb{C}^{2 \times 2} \text{ is the measurement outcome for qubit } i.$$

- $\sigma_T(\rho)$  is a  $2^n \times 2^n$  random matrix with  $\mathbb{E}\sigma_T(\rho) = \rho$  and takes  $\mathcal{O}(nT)$  bits to represent.



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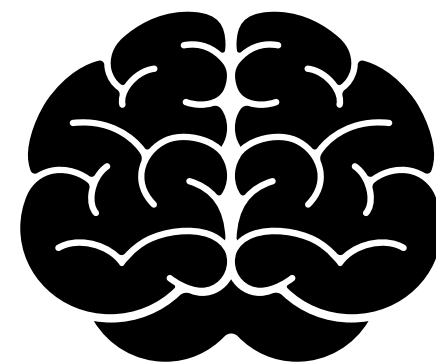
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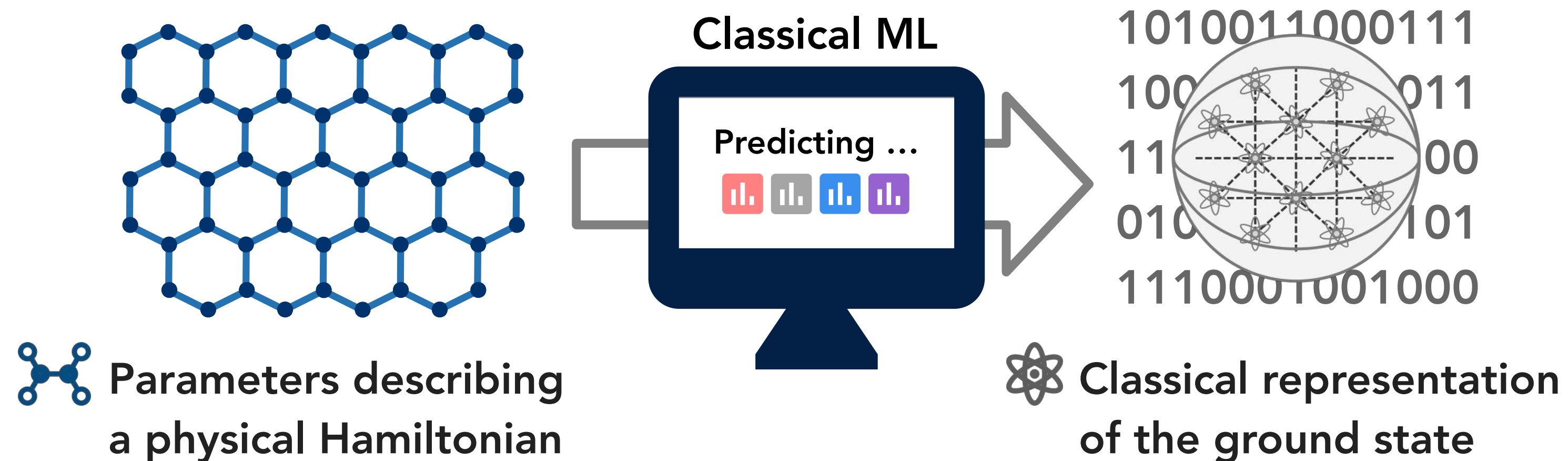
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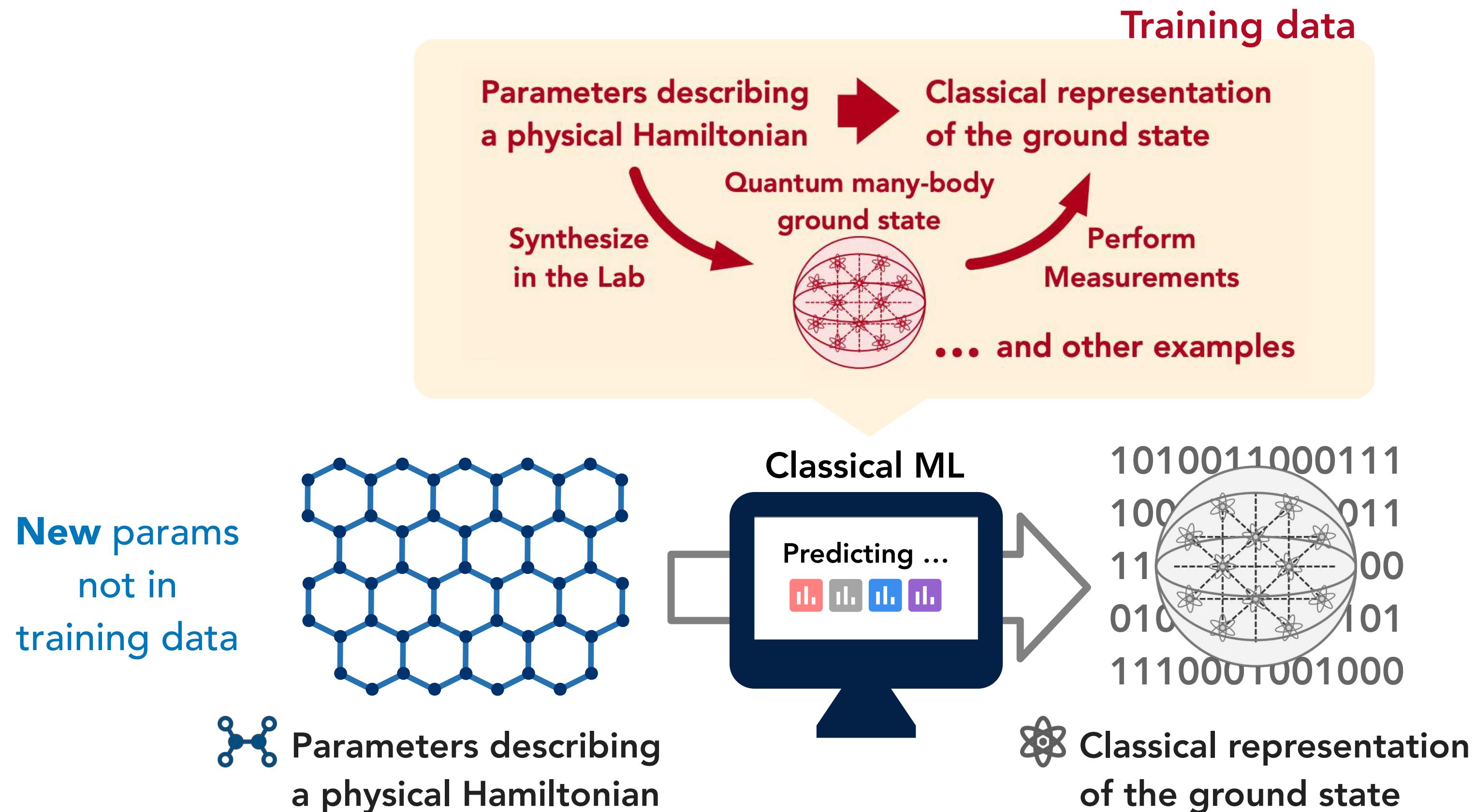
# Predicting ground states: Task

- Given parameters  $x$  that describes a Hamiltonian  $H(x)$ , the machine needs to predict a classical representation of the ground state  $\rho(x)$  of  $H(x)$ .
- $x \in \mathbb{R}^m$  describes laser intensities, few-body interactions, magnetic fields, etc.
- We assume that  $x \mapsto H(x)$  is not known exactly. And we represent  $\rho(x)$  on a classical computer using its classical shadow  $\sigma_T(\rho(x))$ .



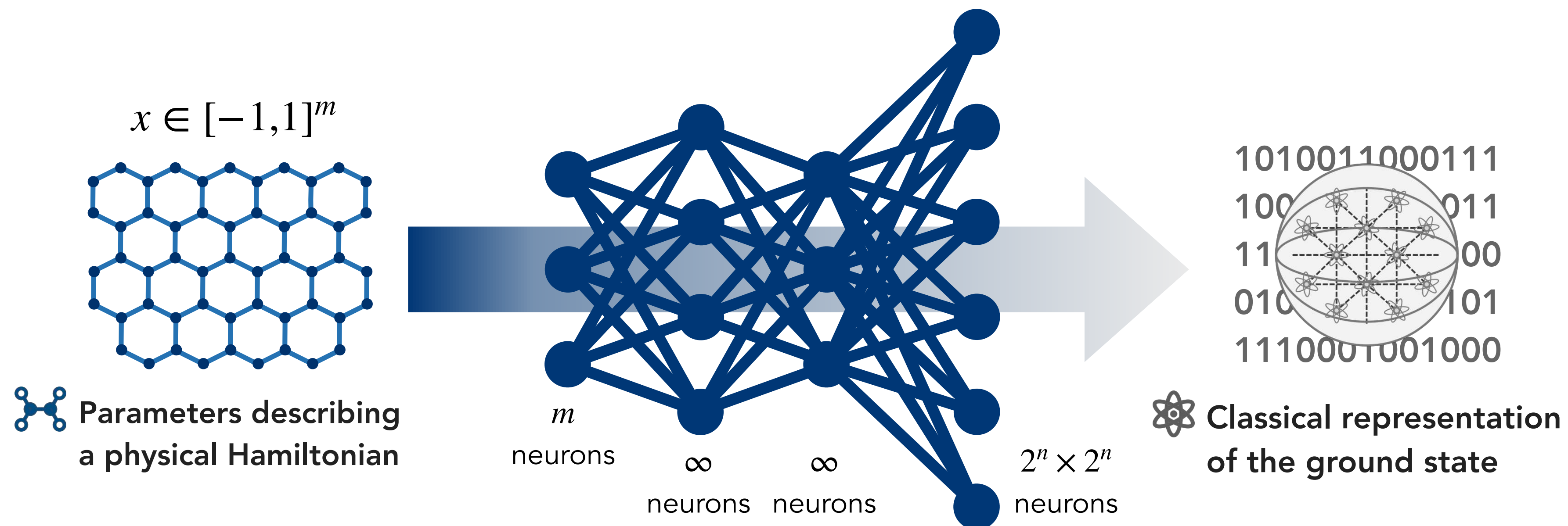
# Predicting ground states: Task

- $x \in \mathbb{R}^m$  describes laser intensities, few-body interactions, magnetic fields, etc. We will normalize such that  $x \in [-1,1]^m$ .
- Training data: examples of params and associated ground state  $\{x_\ell \rightarrow \sigma_T(\rho(x_\ell))\}_{\ell=1}^N$ .



# Predicting ground states: ML

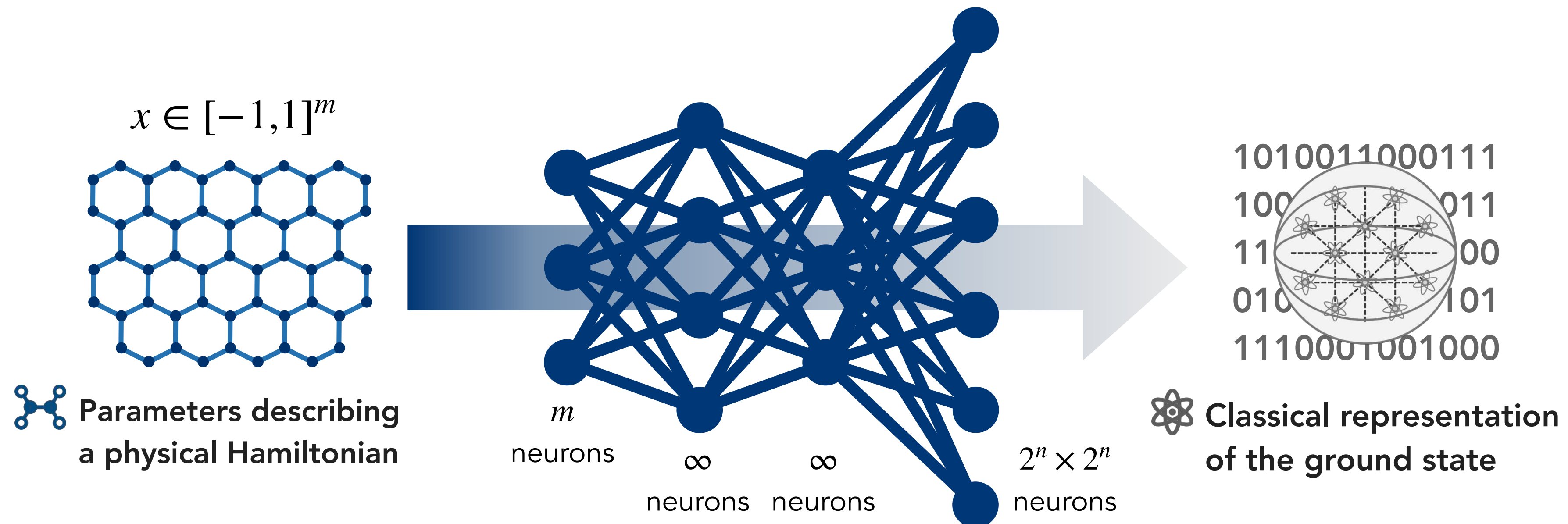
- Training data:  $\{x_\ell \rightarrow \sigma_T(\rho(x_\ell))\}_{\ell=1}^N$ , where  $x_\ell \in \mathbb{R}^m$ ,  $\sigma_T(\rho(x_\ell)) \in \mathbb{C}^{2^n \times 2^n}$ .
- We consider training an ML model that takes in an  $m$ -dim vector  $x$  and outputs a  $2^n \times 2^n$ -size matrix  $\hat{\sigma}(x)$ ; more precisely, an efficient representation of  $\hat{\sigma}(x)$ .
- The ML model needs to be trained within time polynomial in  $n, m$ .



# Predicting ground states: ML

- Suppose we train a neural network  $\hat{\sigma}_W(x)$  (illustrated below) with infinitely many neurons in the hidden layers and exponentially many neurons in the output layer.
- We can train the bizarrely large model in time polynomial in  $n, m$ .

Training data:  $\{x_\ell \rightarrow \sigma_T(\rho(x_\ell))\}_{\ell=1}^N$ , where  $x_\ell \in \mathbb{R}^m$ ,  $\sigma_T(\rho(x_\ell)) \in \mathbb{C}^{2^n \times 2^n}$ .



# Predicting ground states: ML

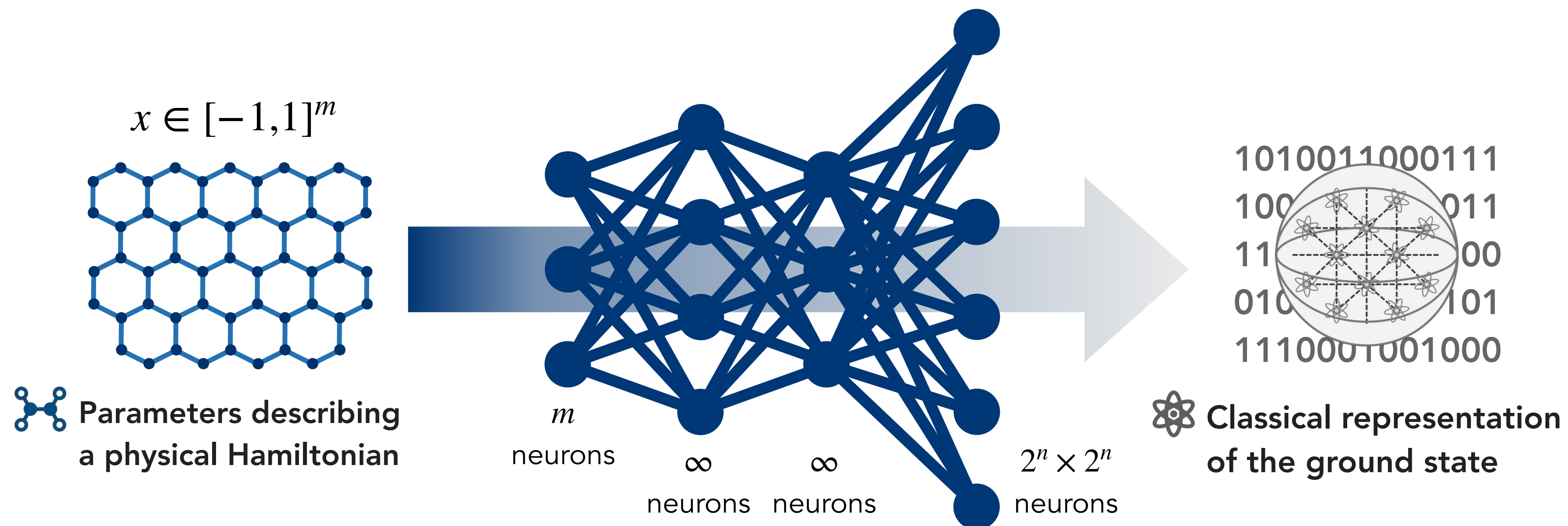
- We show that the neural network after training actually have an analytical form given by

$$\hat{\sigma}^{\text{NN}}(x) = \operatorname{argmin}_{\hat{\sigma}_W} \sum_{\ell=1}^N \|\hat{\sigma}_W(x_\ell) - \sigma_T(\rho(x_\ell))\|_2^2 = \sum_{\ell=1}^N \kappa^{\text{NN}}(x, x_\ell) \sigma_T(\rho(x_\ell))$$

where the learned function  $\kappa^{\text{NN}}(x, x_\ell) \in \mathbb{R}$  can be obtained efficiently; based on [JGH18].

[JGH18] "Neural tangent kernel: Convergence and generalization in neural networks."  
arXiv preprint arXiv:1806.07572 (2018).

Training data:  $\{x_\ell \rightarrow \sigma_T(\rho(x_\ell))\}_{\ell=1}^N$ , where  $x_\ell \in \mathbb{R}^m$ ,  $\sigma_T(\rho(x_\ell)) \in \mathbb{C}^{2^n \times 2^n}$ .



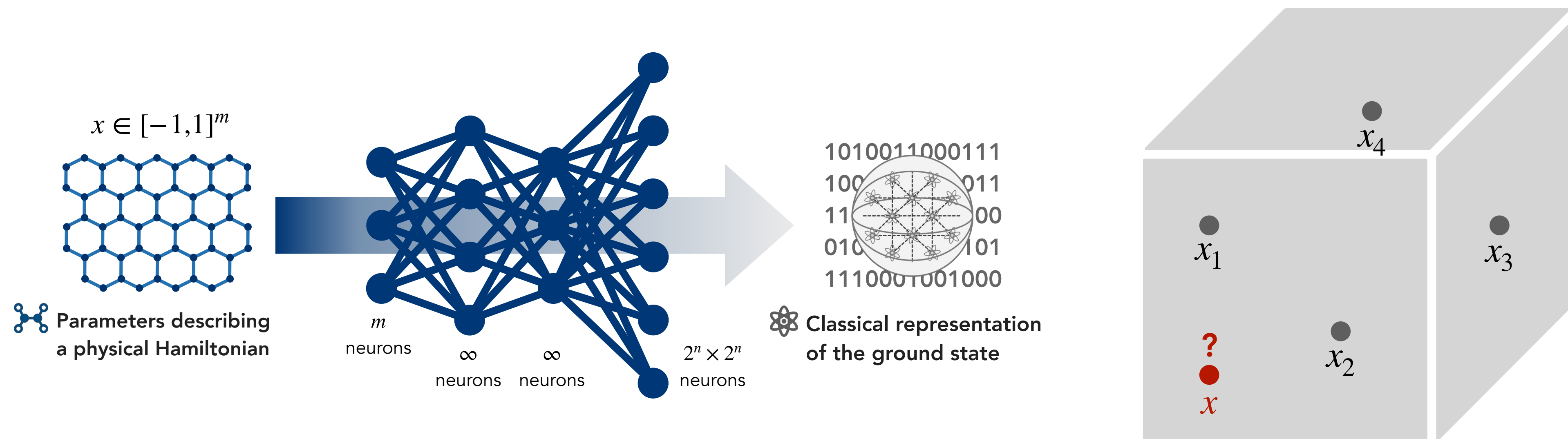
# Predicting ground states: ML

- Furthermore, various machine learning models (kernel methods, infinite-width neural networks, etc.) can be shown to yield an analytical form as the global minimum of the optimization (training):

$$\hat{\sigma}(x) = \sum_{\ell=1}^N \kappa(x, x_{\ell}) \sigma_T(\rho(x_{\ell}))$$

where  $\kappa(x, x_{\ell}) \in \mathbb{R}$  is a learned function for how to extrapolate the known examples to the full space.

Training data:  $\{x_{\ell} \rightarrow \sigma_T(\rho(x_{\ell}))\}_{\ell=1}^N$ , where  $x_{\ell} \in \mathbb{R}^m$ ,  $\sigma_T(\rho(x_{\ell})) \in \mathbb{C}^{2^n \times 2^n}$ .



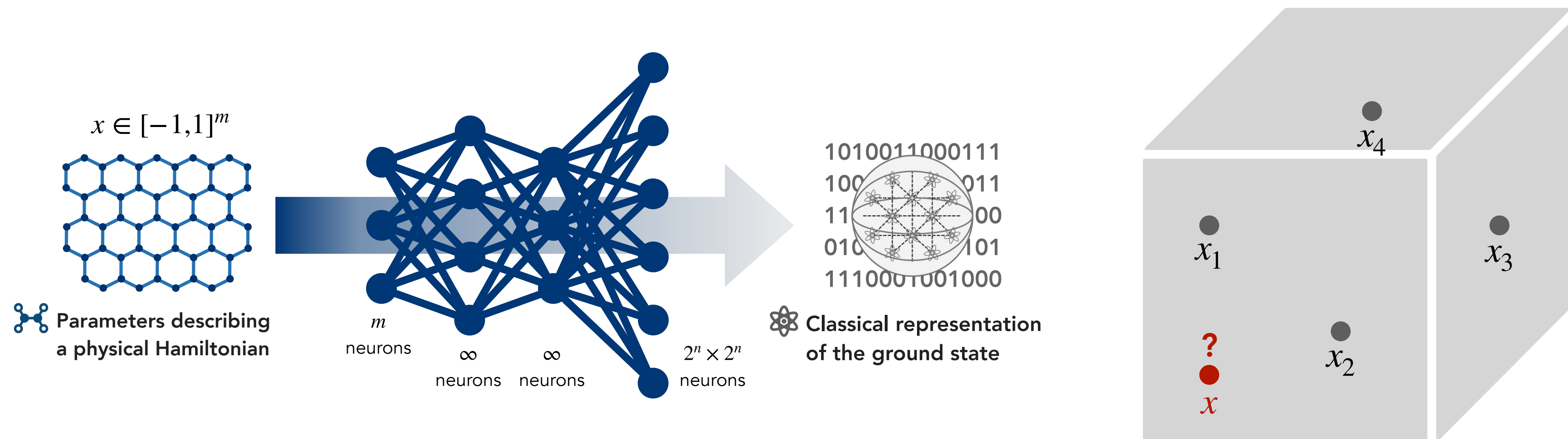
# Predicting ground states: ML

- As long as  $\kappa(x, x_\ell) \in \mathbb{R}$  is efficiently computable, the ML model's prediction

$$\hat{\sigma}(x) = \sum_{\ell=1}^N \kappa(x, x_\ell) \sigma_T(\rho(x_\ell))$$

can be represented efficiently with  $\mathcal{O}(nTN)$  bits; recall  $\sigma_T(\rho(x_\ell))$  only require  $\mathcal{O}(nT)$  bits.

Training data:  $\{x_\ell \rightarrow \sigma_T(\rho(x_\ell))\}_{\ell=1}^N$ , where  $x_\ell \in \mathbb{R}^m$ ,  $\sigma_T(\rho(x_\ell)) \in \mathbb{C}^{2^n \times 2^n}$ .





# Predicting ground states: Theorem

- We consider an ML model  $\hat{\sigma}(x) = \sum_{\ell=1}^N \kappa(x, x_{\ell}) \sigma_T(\rho(x_{\ell}))$  with  $l_2$ -Dirichlet kernel.
- The learned model  $\hat{\sigma}(x)$  captures the ground state properties accurately (on average).

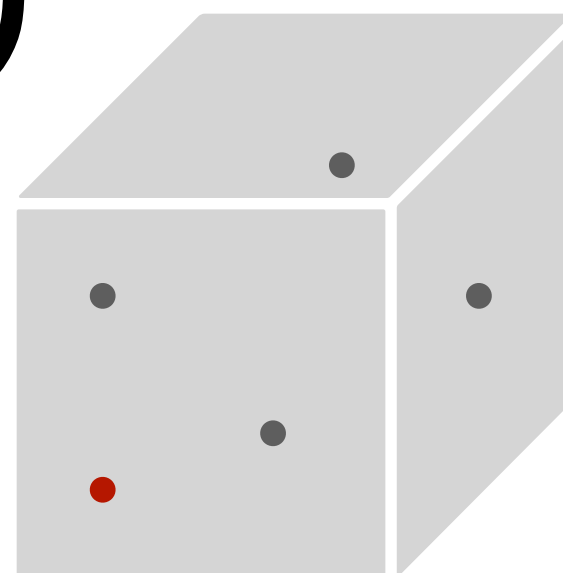
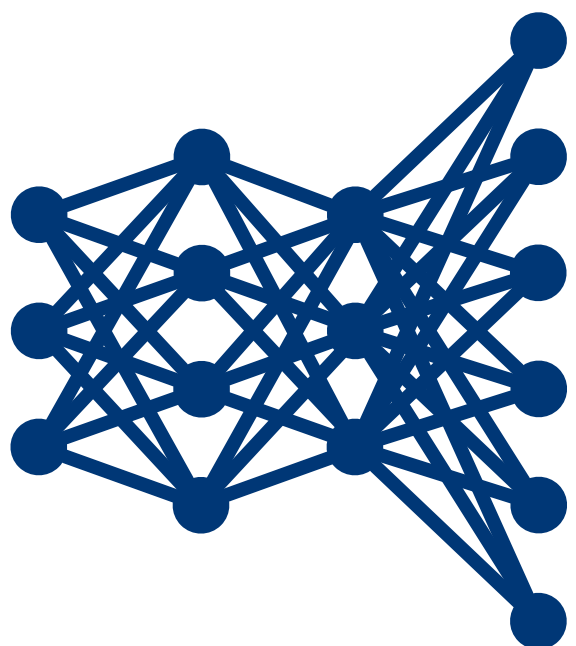
## Theorem 1

For any smooth class of local Hamiltonians  $H(x)$  in a **finite spatial dimension** with a **const. spectral gap**, given the number of training data  $N = \text{poly}(m)$  and  $T = 1$  (one randomized Pauli measurements each),

$$\mathbb{E}_{x \sim [-1,1]^m} |\text{Tr}(O\hat{\sigma}(x)) - \text{Tr}(O\rho(x))|^2 \leq \epsilon,$$

for any sum of local observables  $O = \sum_{j=1}^L O_j$  with  $\sum_{j=1}^L \|O_j\| = \mathcal{O}(1)$  and  $\epsilon$ : const. Training and prediction

time are polynomial in  $m$  and linear in system size  $n$ .



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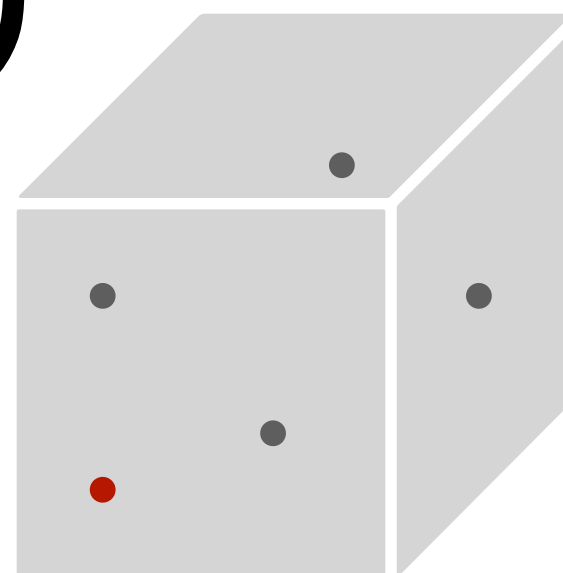
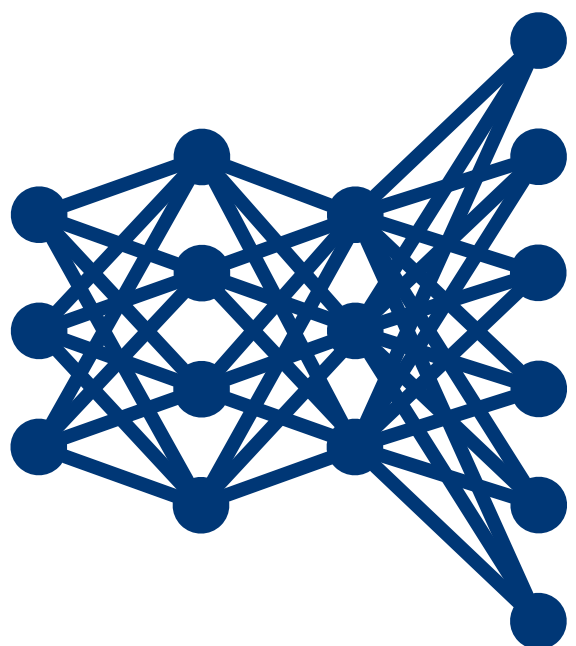
Intuitively, in a quantum phase

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# Predicting ground states: Theorem

- Key steps in the proof:
  1. Constant spectral gap implies some “smoothness” condition in ground state space (spectral flow + Lieb-Robinson bounds).
  2. Generalization error bounds for the proposed ML with  $\ell_2$ -Dirichlet kernel trained on randomized measurement data under the “smoothness” guarantee (statistical analysis + #lattices in a  $m$ -dim. sphere).

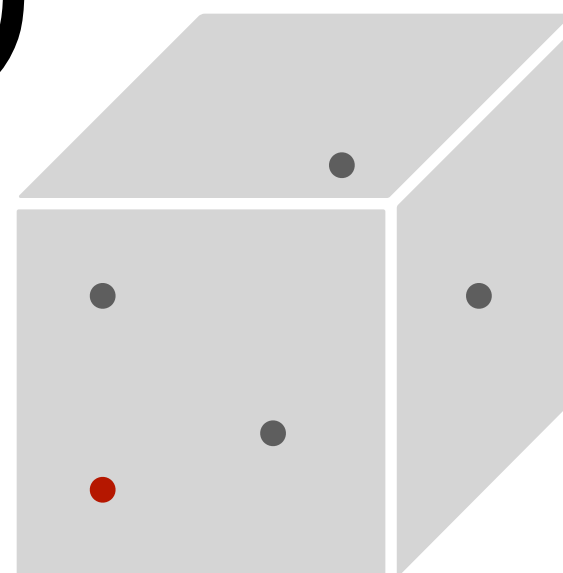
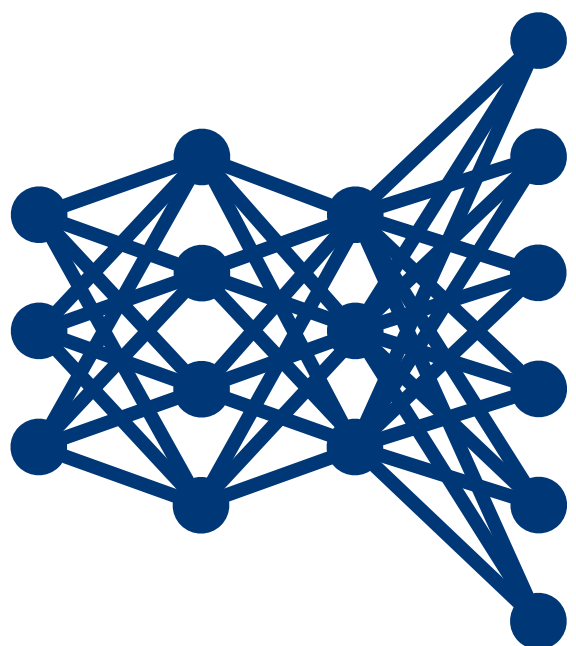
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# Predicting ground states: Theorem

- A limitation:  $\epsilon$  can only be a constant. In particular  $N = m^{\mathcal{O}(1/\epsilon)}$ .
- One may wonder if quantum ML algorithm could overcome this limitation.
- We prove in the appendix that any quantum (classical) ML algorithm require  $N = m^{\Omega(1/\epsilon)}$ , so the advantage of quantum ML can only be polynomial.

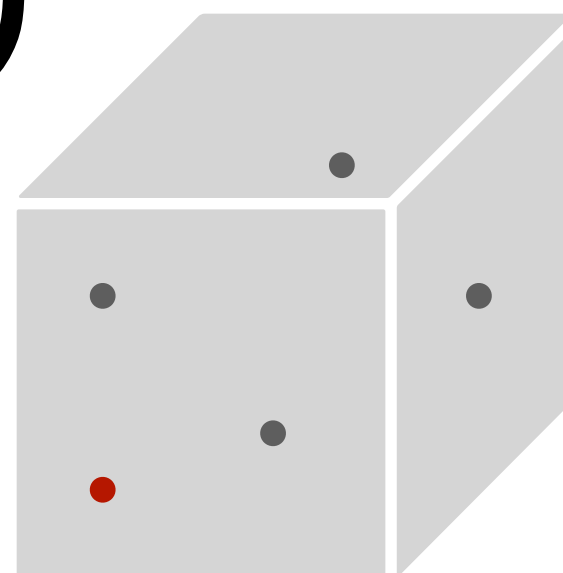
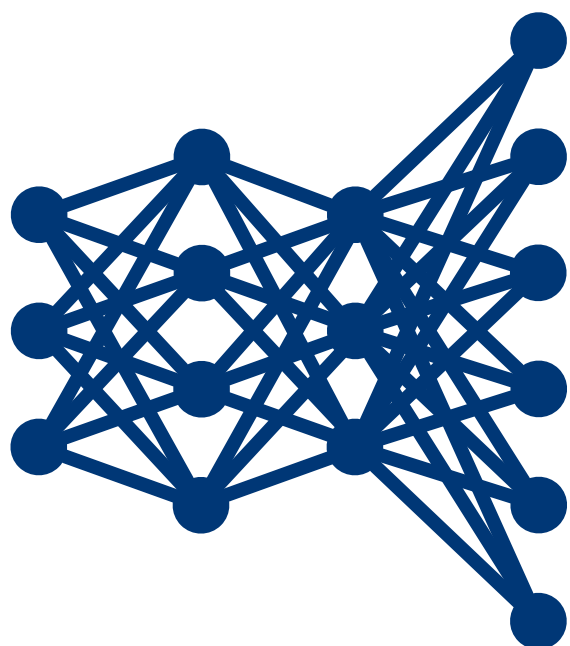
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# Provable advantage of learning algorithm with data

## Proposition 1

If a classical polynomial-time randomized algorithm  $\mathcal{A}$  can achieve

$$\mathbb{E}_{x \sim [-1,1]^m} |\mathcal{A}(x, O) - \text{Tr}(O\rho(x))|^2 \leq 1/4,$$

for any one-local observables  $O$  and any smooth class of local Hamiltonians in a two spatial dimension with a constant spectral gap, then  $\text{RP} = \text{NP}$ .

$\text{RP} = \text{NP}$ : NP-complete problems can be solved in randomized polynomial time.

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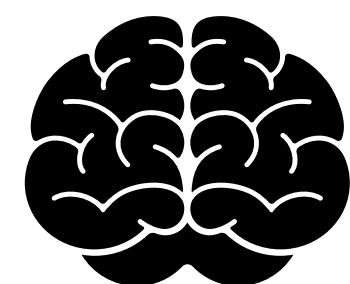
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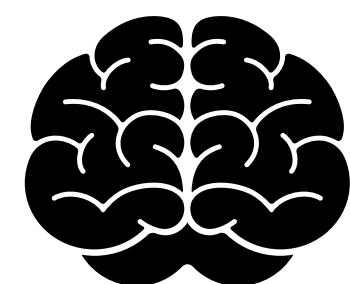
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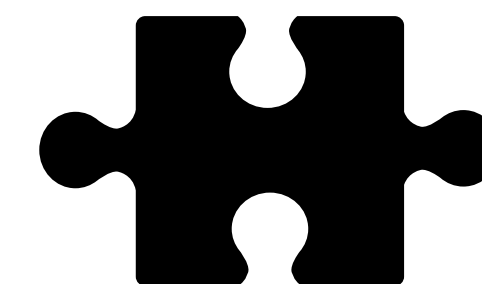
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The answer ⚡:  
Yes, generalizing from data can be easier than computing everything






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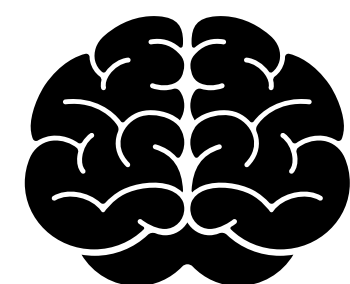
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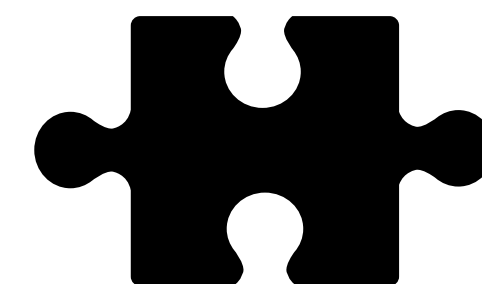
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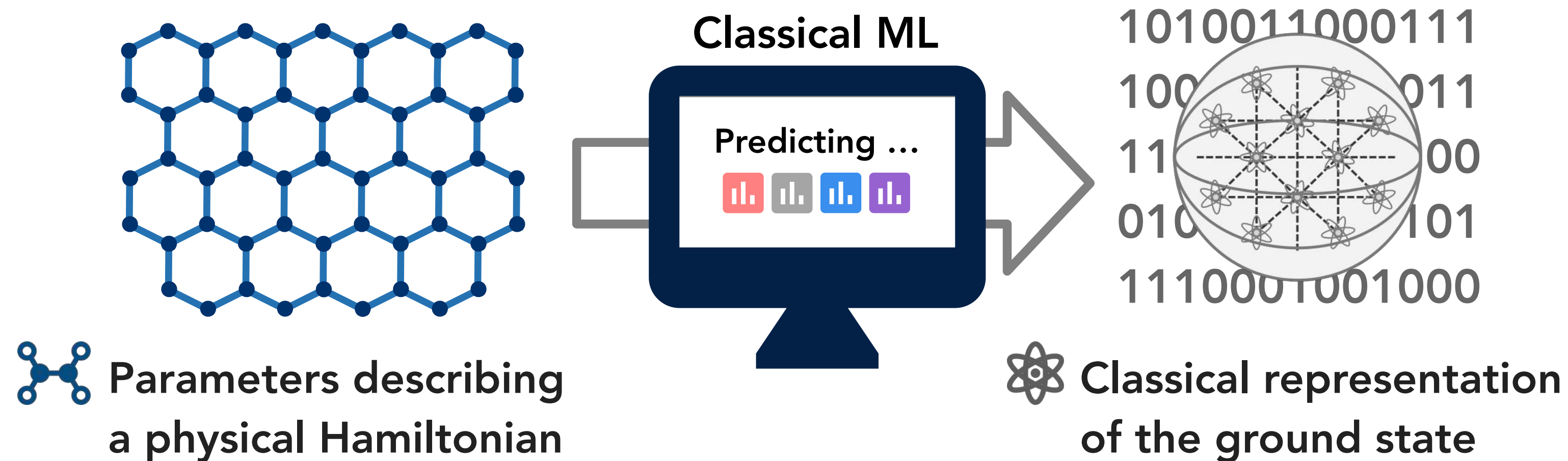
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Data contain computational power  
(e.g., nature operates quantumly)

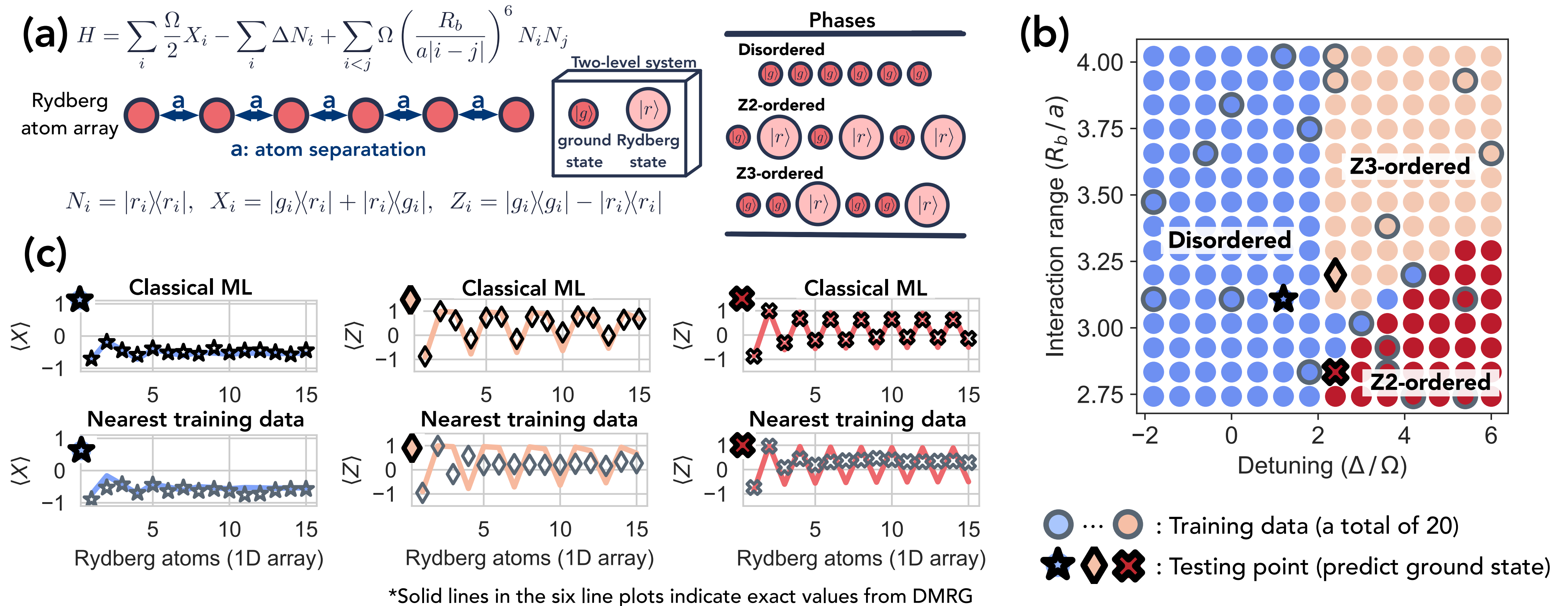
# Predicting ground states: Numerics

- How well does classical ML algorithm perform in actual physical systems?



# 1D Rydberg atom array

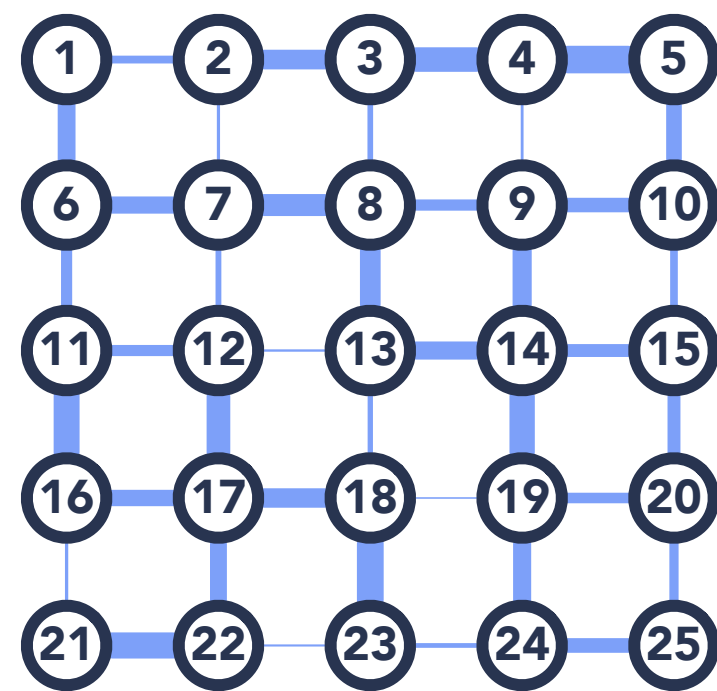
We consider training data size  $N = 20$ ,  $T = 500$  randomized measurements for constructing classical shadows. The best ML model is chosen from Gaussian kernel method, infinite-width neural networks, and  $l_2$ -Dirichlet kernel.



# 2D random Heisenberg model

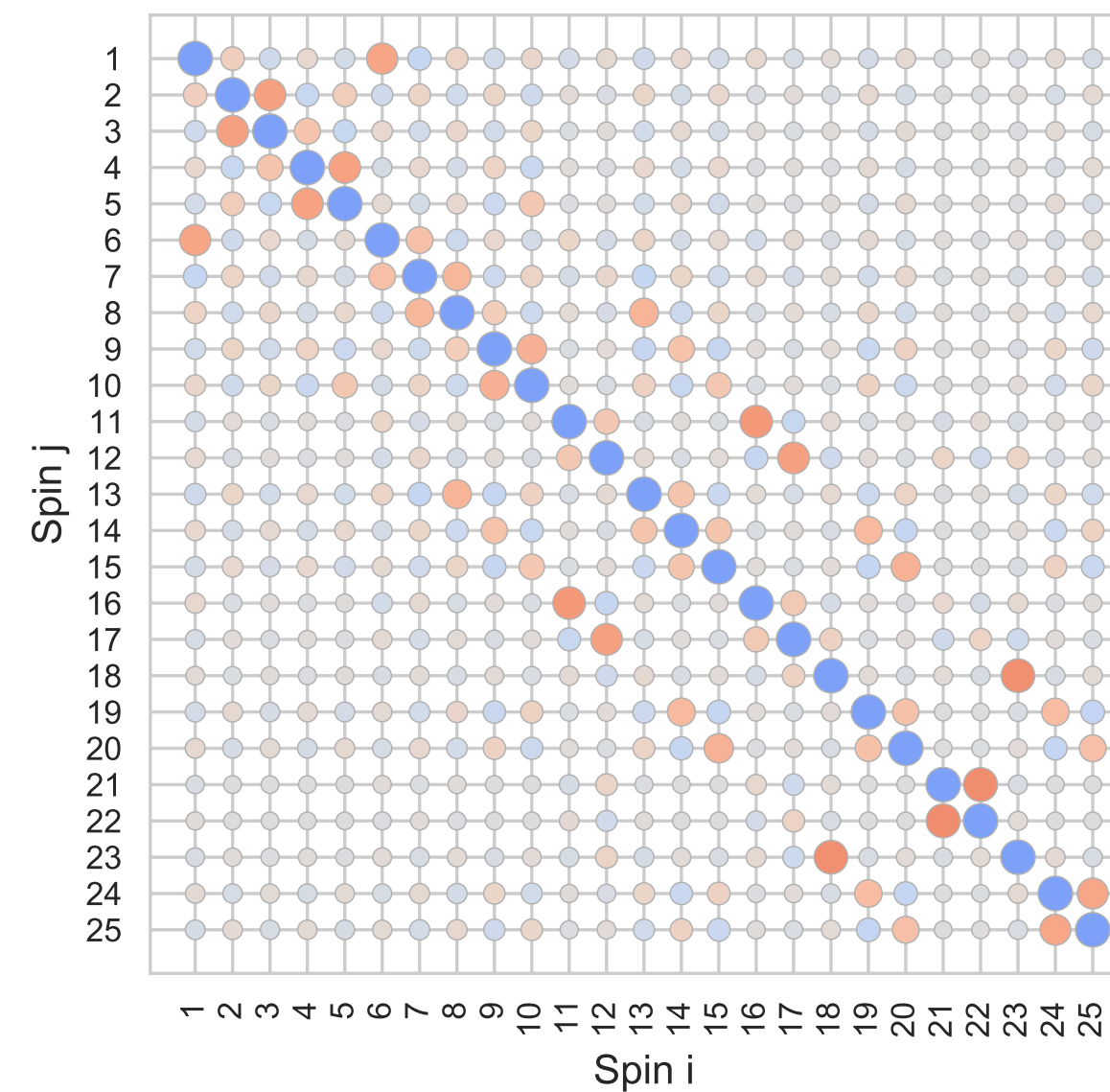
We consider training data size  $N = 100$ ,  $T = 500$  randomized measurements for constructing classical shadows. The best ML model is chosen from Gaussian kernel method, infinite-width neural networks, and  $l_2$ -Dirichlet kernel.

(a) 2D anti-ferromagnetic random Heisenberg model  
$$H = \sum_{\langle ij \rangle} J_{ij}(X_i X_j + Y_i Y_j + Z_i Z_j)$$

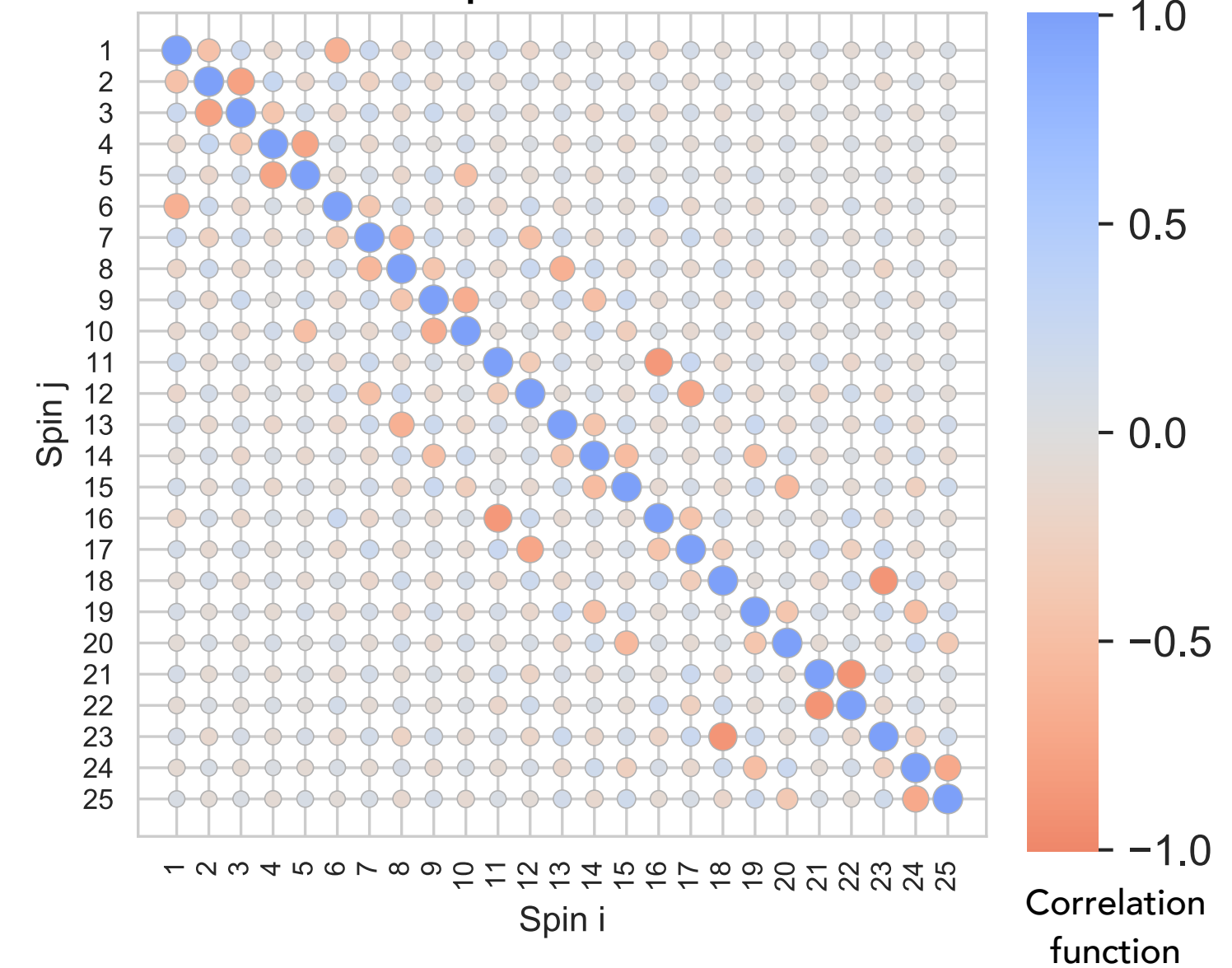


\*The random J considered in (c)

(b) Exact values from DMRG

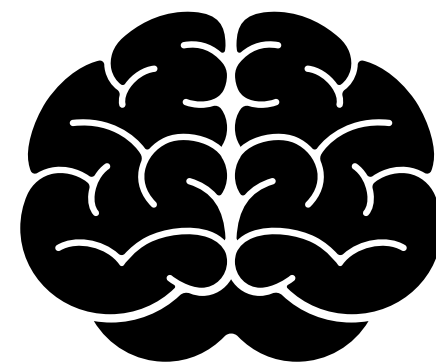


ML predictions



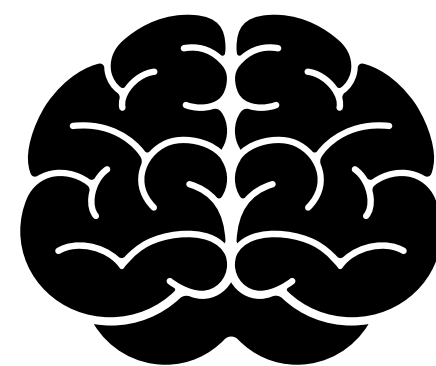
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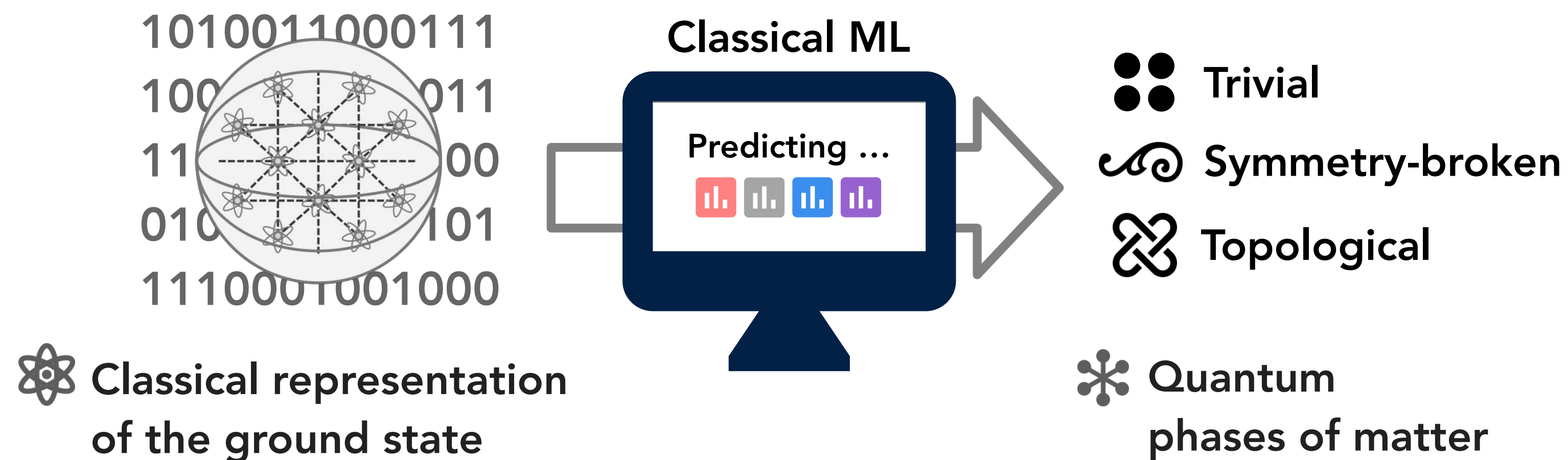
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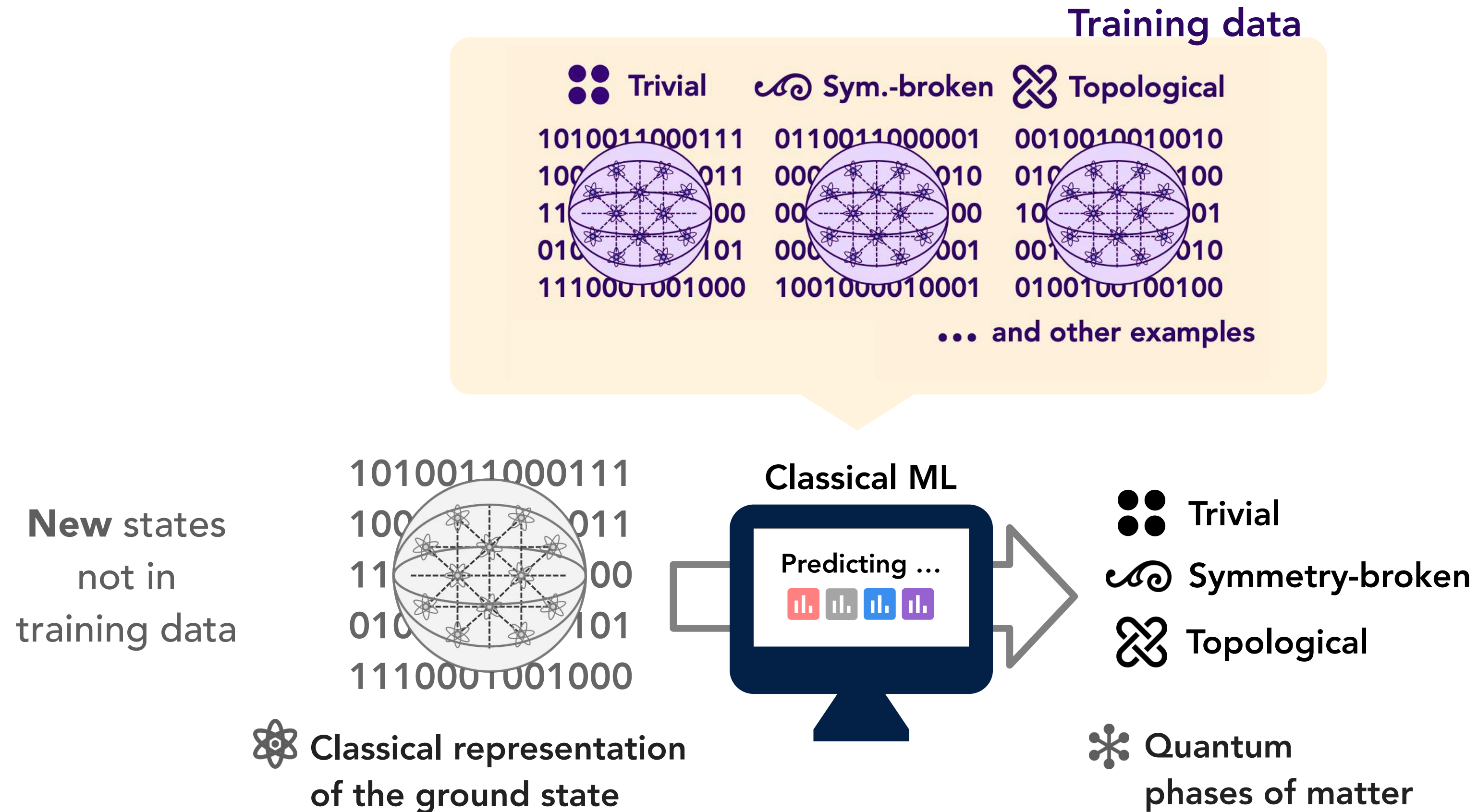
# Classifying quantum phases: Task

- Given a quantum state  $\rho$ , predict which quantum phases of matter the state  $\rho$  is in.
- We represent the quantum state  $\rho$  using classical shadow, which is a 2D array of measurement outcomes  $S_T(\rho) = \{\sigma_i^{(t)}\}_{i=1 \sim n, t=1 \sim T}$  with  $\sigma_T(\rho) = \frac{1}{T} \sum_{t=1}^T \sigma_1^{(t)} \otimes \dots \otimes \sigma_n^{(t)} \approx \rho$ .



# Classifying quantum phases: Task

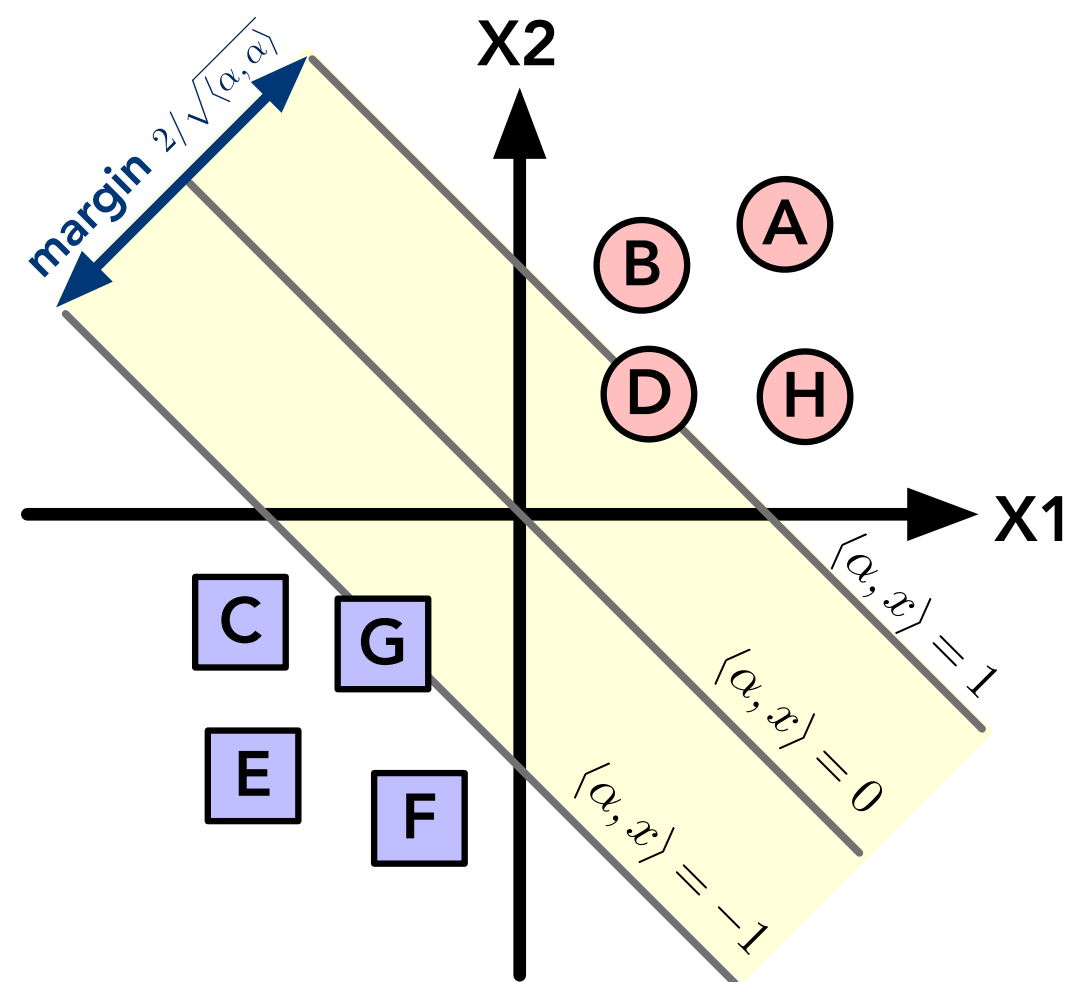
- Given a quantum state  $\rho$ , classify which quantum phases of matter the state  $\rho$  is in.
- Training data: examples of states and associated phase.





# Classifying quantum phases: ML

- The ML model tries to find a classifying function that separates the phases of matter well.
- For symmetry-broken phases, there is typically a local observable  $O$  with
$$\text{Tr}(O\rho_A) > 0, \forall \rho_A \in \text{phase A}, \quad \text{Tr}(O\rho_B) \leq 0, \forall \rho_B \in \text{phase B}.$$
- Then the classical ML model only need to learn a linear function (easy with linear classifiers).
- But Proposition 2 shows that it is not possible to classify topological phases.



## Proposition 2

Consider two distinct topological phases A and B.

No (local/global) observable  $O$  exists such that

$$\text{Tr}(O\rho_A) > 0, \forall \rho_A \in \text{phase A}, \quad \text{Tr}(O\rho_B) \leq 0, \forall \rho_B \in \text{phase B}.$$

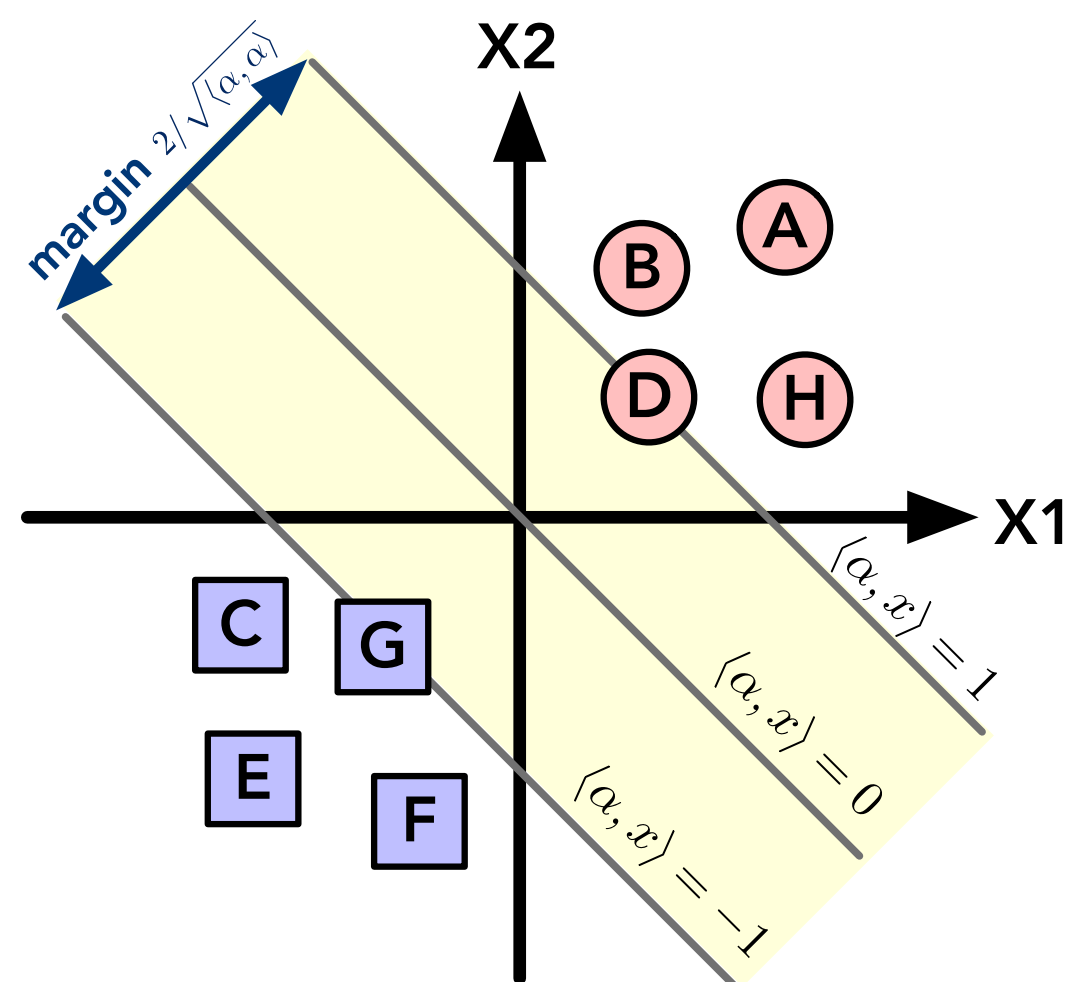
# Classifying quantum phases: ML

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## Proposition 2

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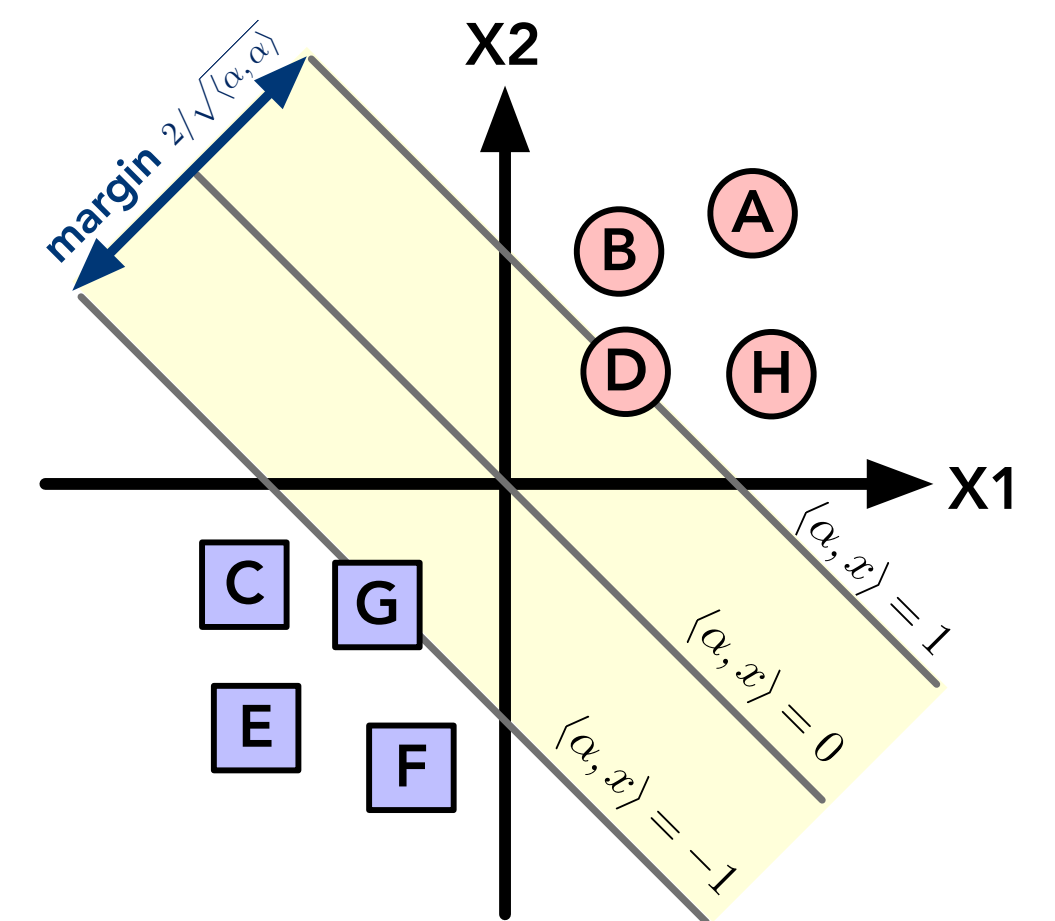
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# Classifying quantum phases: ML

- Classical ML model: Learn a linear function in  $\phi^{(\text{shadow})}(S_T(\rho))$  equiv. a nonlinear function in  $\rho$ .
- All we need is to efficiently compute the inner product (referred to as shadow kernel)

$$\langle \phi^{(\text{shadow})}(S_T(\rho)), \phi^{(\text{shadow})}(S_T(\tilde{\rho})) \rangle = \exp \left( \frac{\tau}{T^2} \sum_{t,t'=1}^T \exp \left( \frac{\gamma}{n} \sum_{i=1}^n \text{Tr} \left( \sigma_i^{(t)} \tilde{\sigma}_i^{(t')} \right) \right) \right) \equiv k^{(\text{shadow})}(S_T(\rho), S_T(\tilde{\rho})).$$

- Computing shadow kernels only take time  $\mathcal{O}(nT^2)$ .
- Training the classical ML model only take time polynomial in  $n, T, N$  (and extremely efficient in practice).



# Classifying quantum phases: Theorem

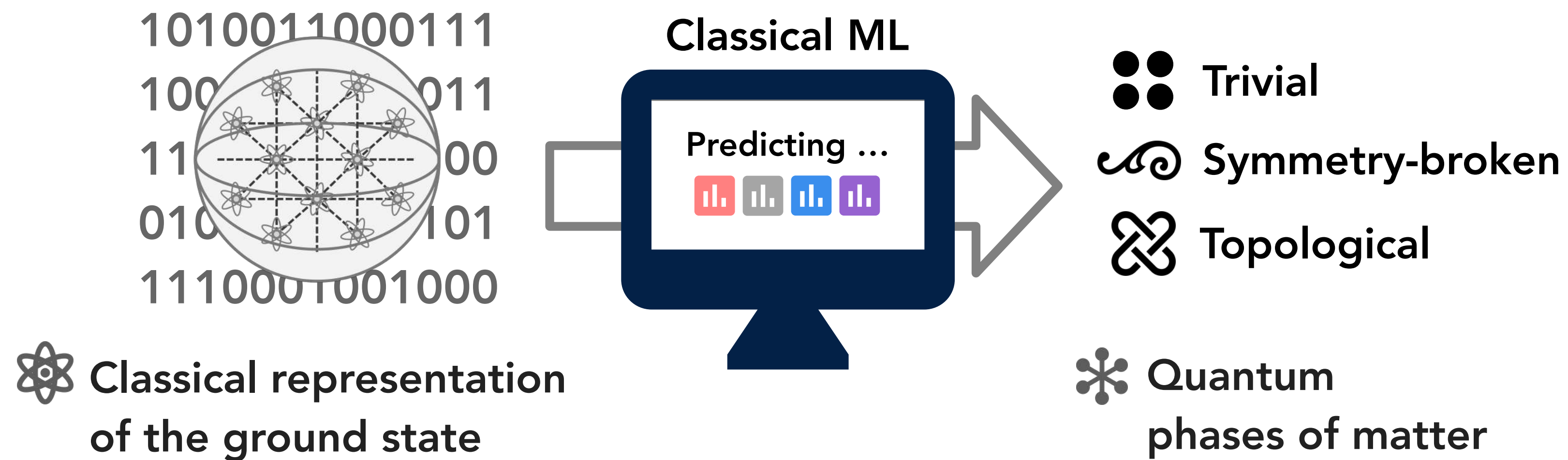
## Theorem 2

If there is a **nonlinear** function of **few-body** reduced density matrices that classifies phases, then the classical algorithm can learn to classify these phases accurately. The amount of training data and computation time scales polynomially in system size.

- The ML model constructs the classifying function explicitly.
- Examples of classifying functions on few-body reduced density matrices (assuming const. spectral gap) include:
  1. Twist operators for 1D Haldane phase with  $O(2)$ -symmetry (linear function)
  2. Hall conductivity for systems adiabatically connected to free fermion (low-degree polynomial)
  3. Topological entanglement entropy in a constant region (nonlinear function)
- As long as the classifying function exists, the ML model with shadow kernel is guaranteed to find it.

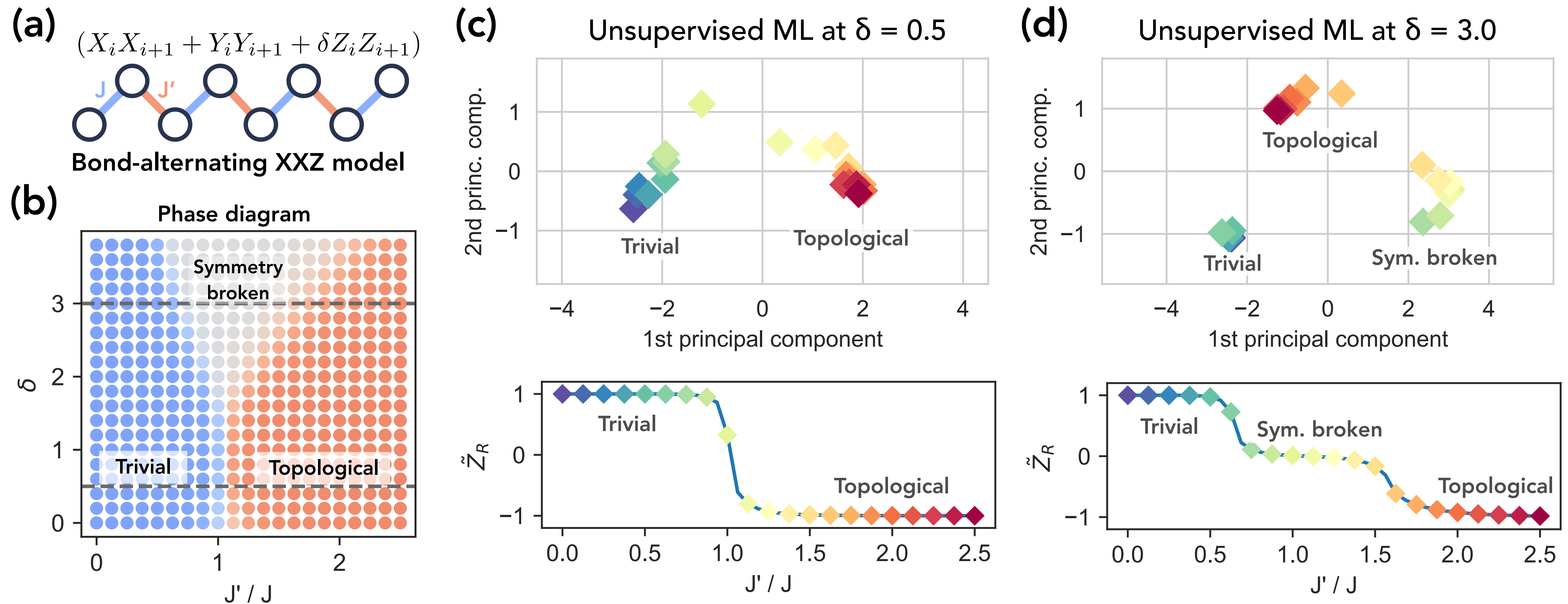
# Classifying quantum phases: Numerics

- How well does the classical ML algorithm perform in actual physical systems?



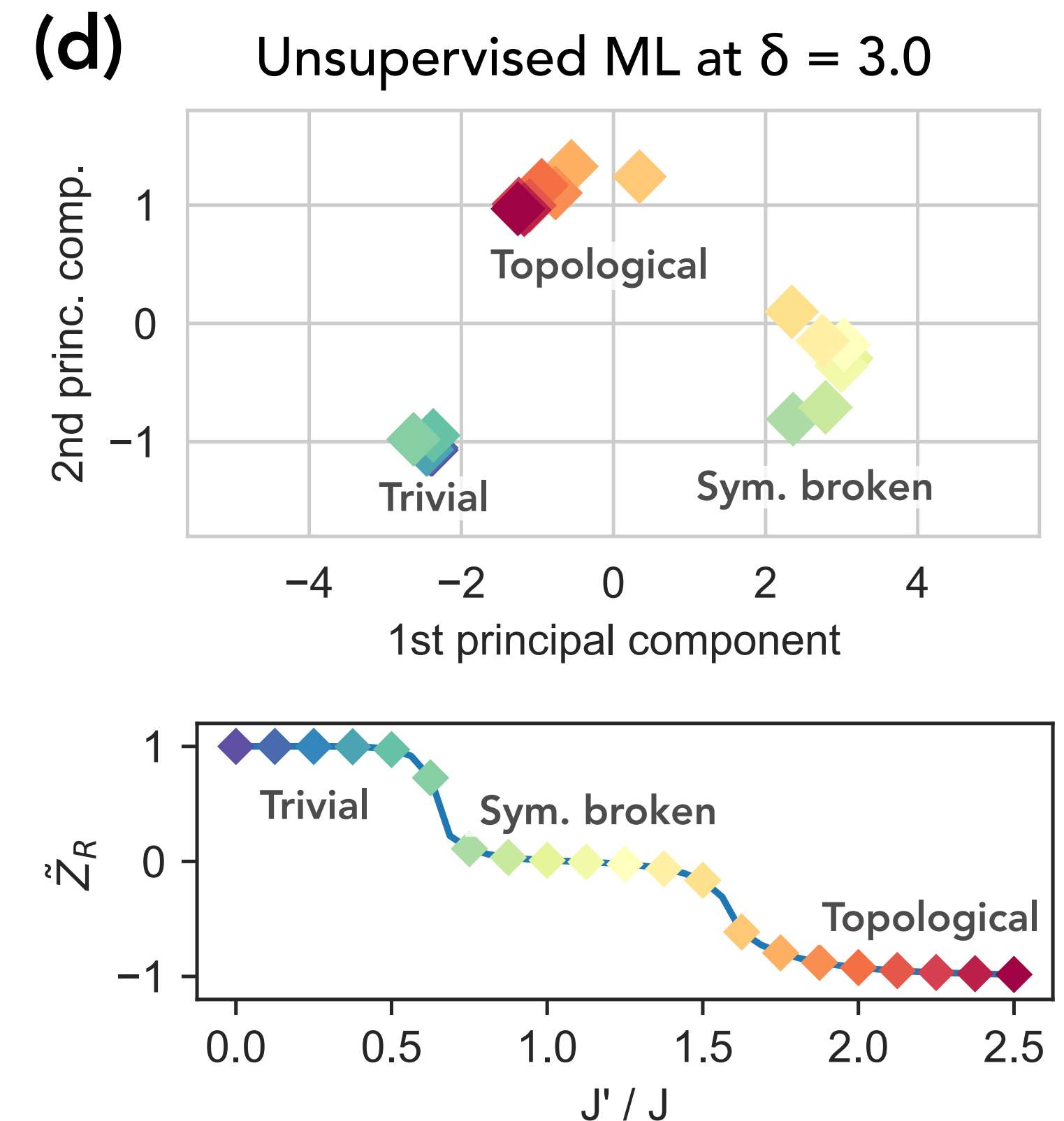
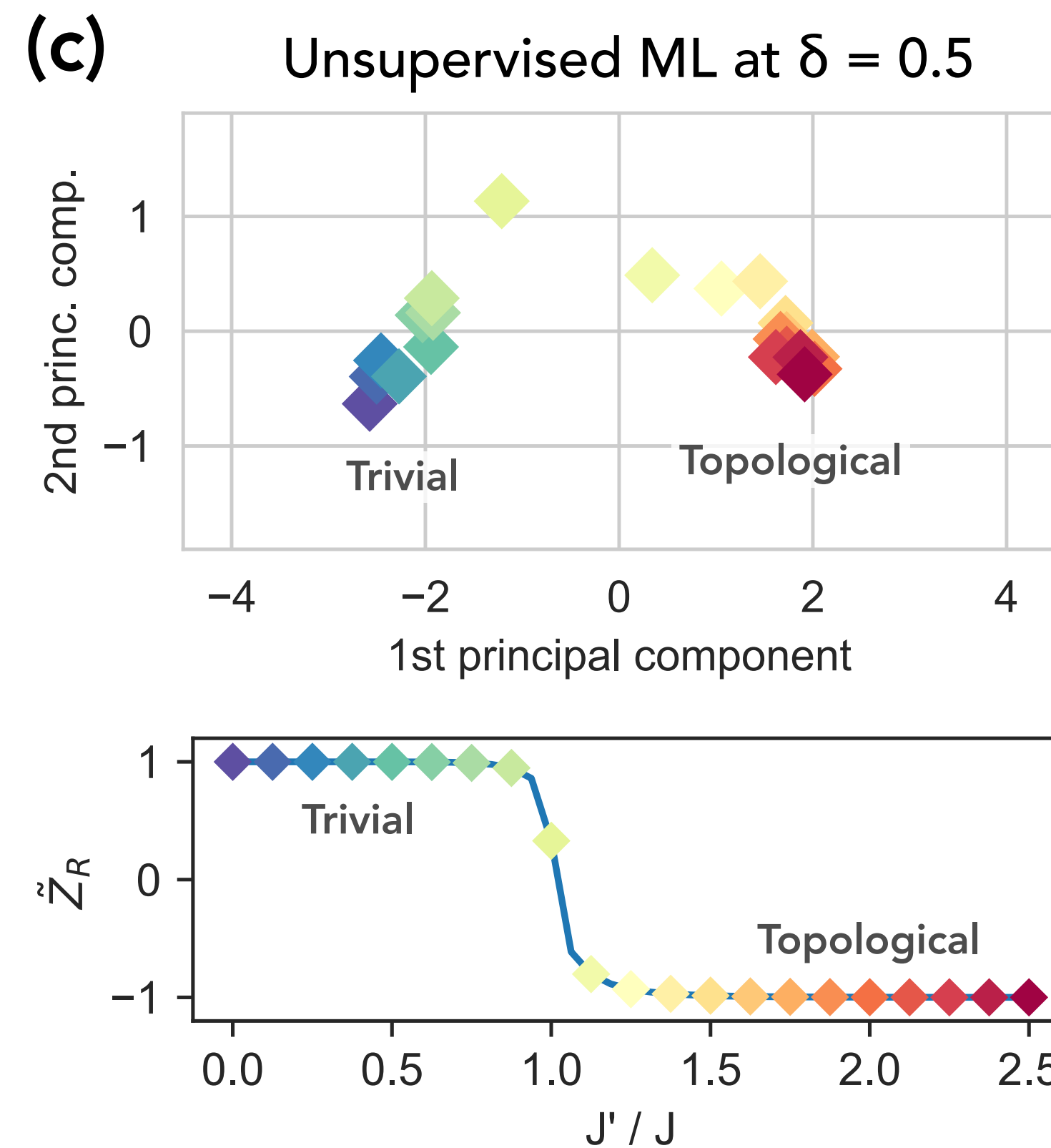
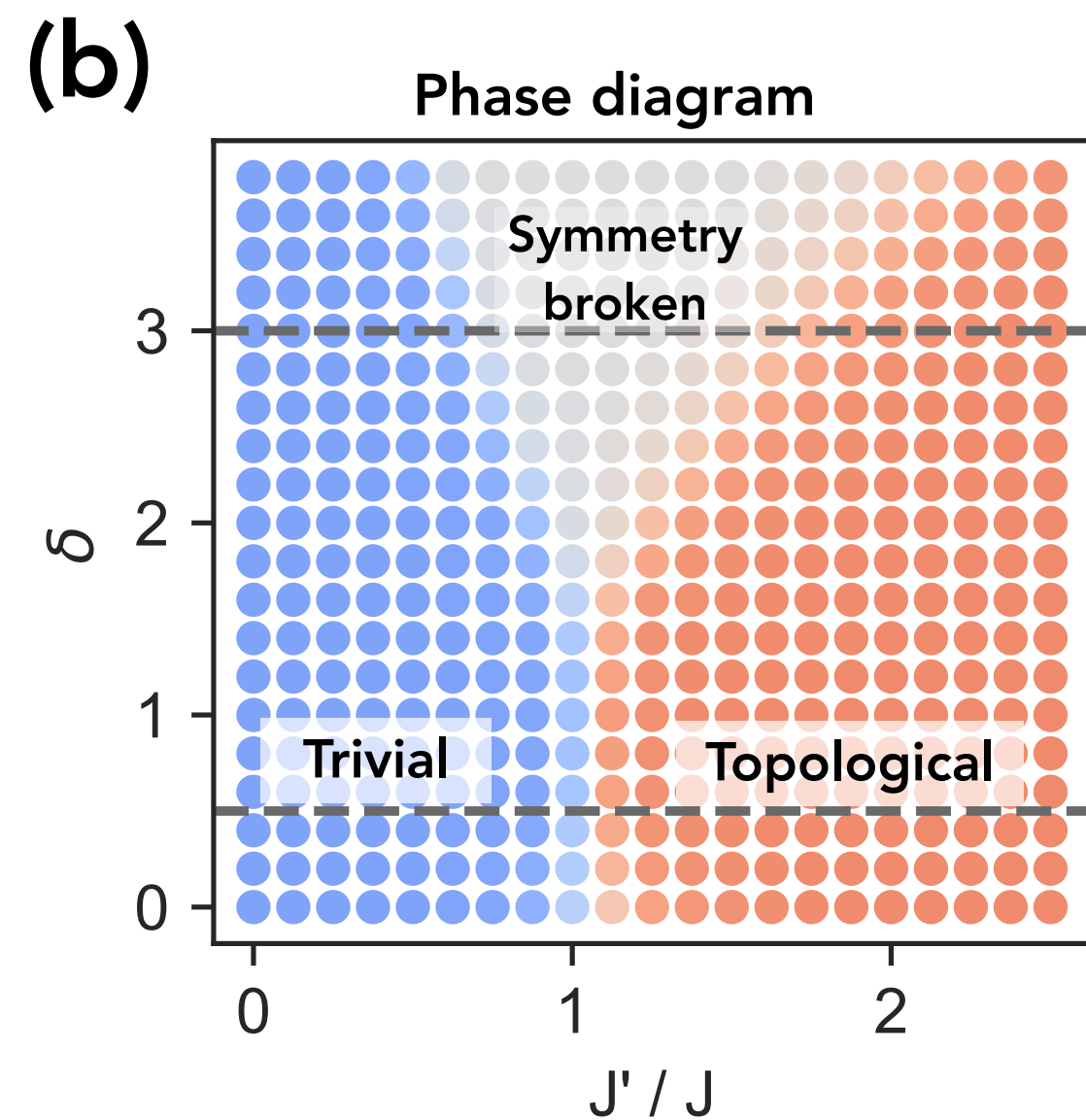
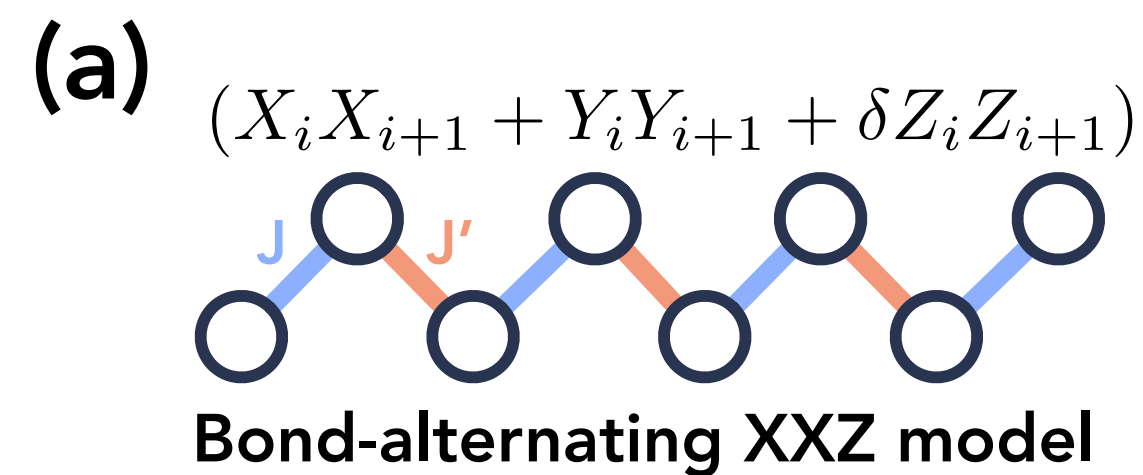
# 1D Symmetry protected topological phases

We consider  $T = 500$  randomized measurements to construct classical shadows for each state.  
The classical unsupervised ML model is a kernel PCA using the shadow kernel.



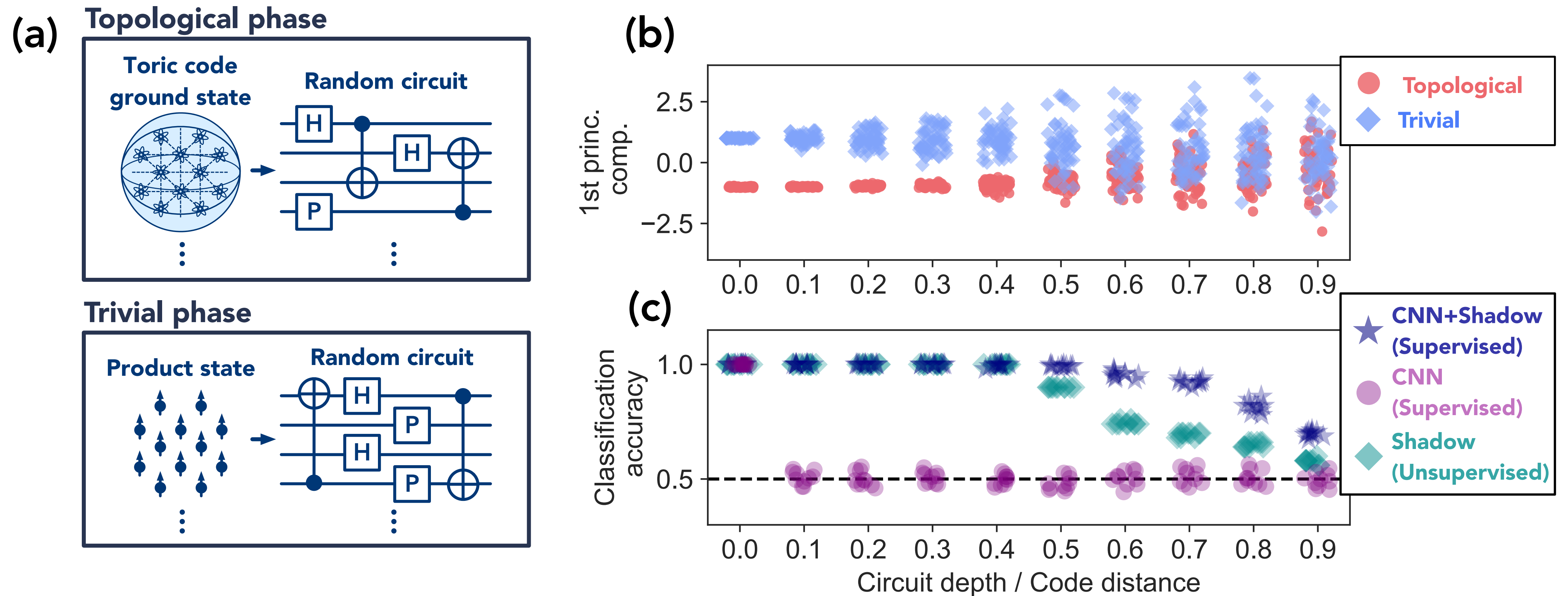
# 1D Symmetry protected topological phases

We consider  $T = 500$  random **No labeled training data** construct classical shadows for each state.  
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# 2D topologically-ordered phases

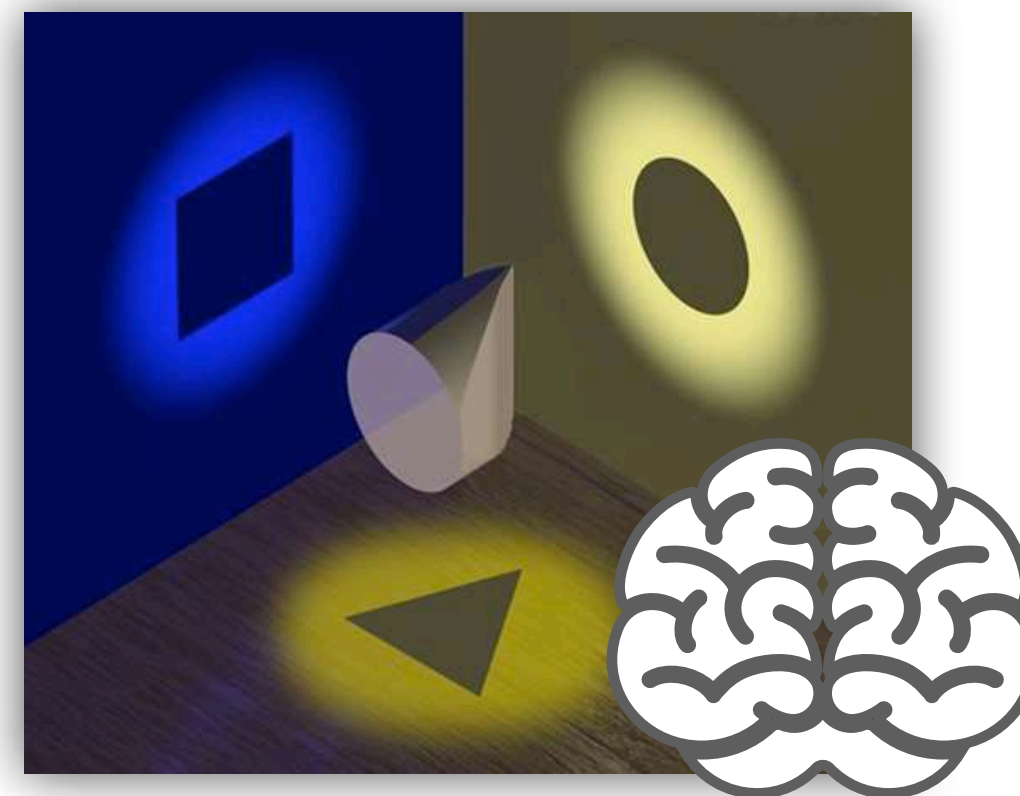
We consider  $T = 500$  randomized measurements to construct classical shadows for each state.  
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# Conclusion

- We prove that classical ML algorithms, informed by data from physical experiments, can effectively address some quantum many-body problems.
- As a consequence, we rigorously establish the advantage of **classical ML models** over classical non-ML algorithms.
- **Open questions:**
  - Advantage of ML over non-ML algorithms in other tasks?
  - Rigorous guarantee for other quantum problems with classical ML?
  - Useful class of quantum learning problems with exponential quantum advantage?



Classical shadows enhanced with ML

**Additional slides**

# Classical shadow formalism

## Theorem 1 [HKP20]

There exists procedure that guarantees the following.

1. Given  $B, \epsilon > 0$ , the procedure learns a classical representation of an unknown quantum state  $\rho$  from

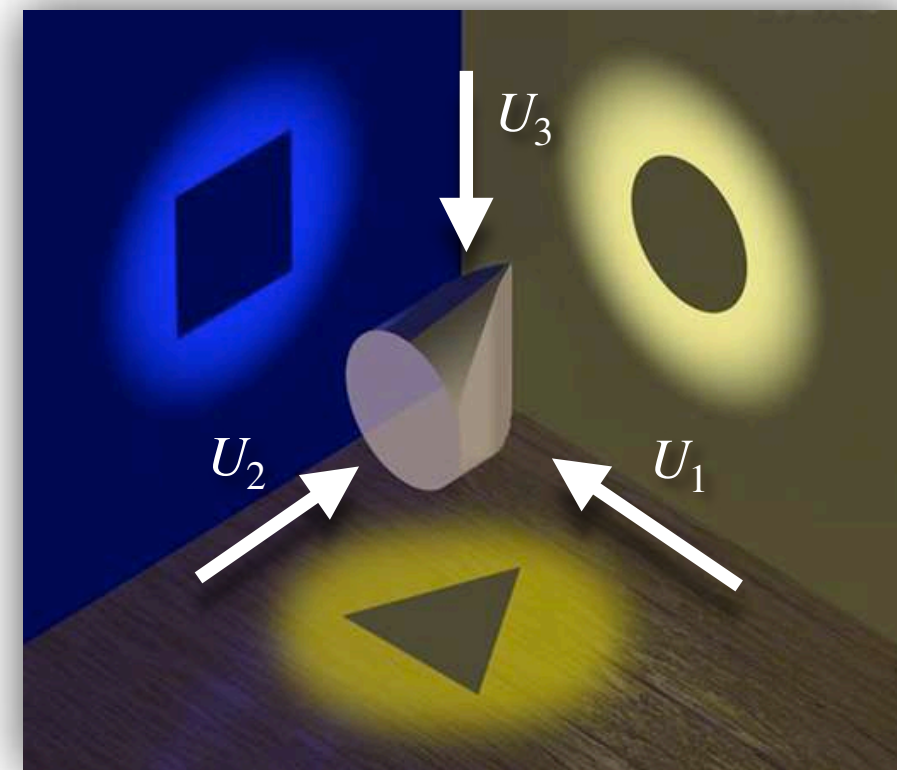
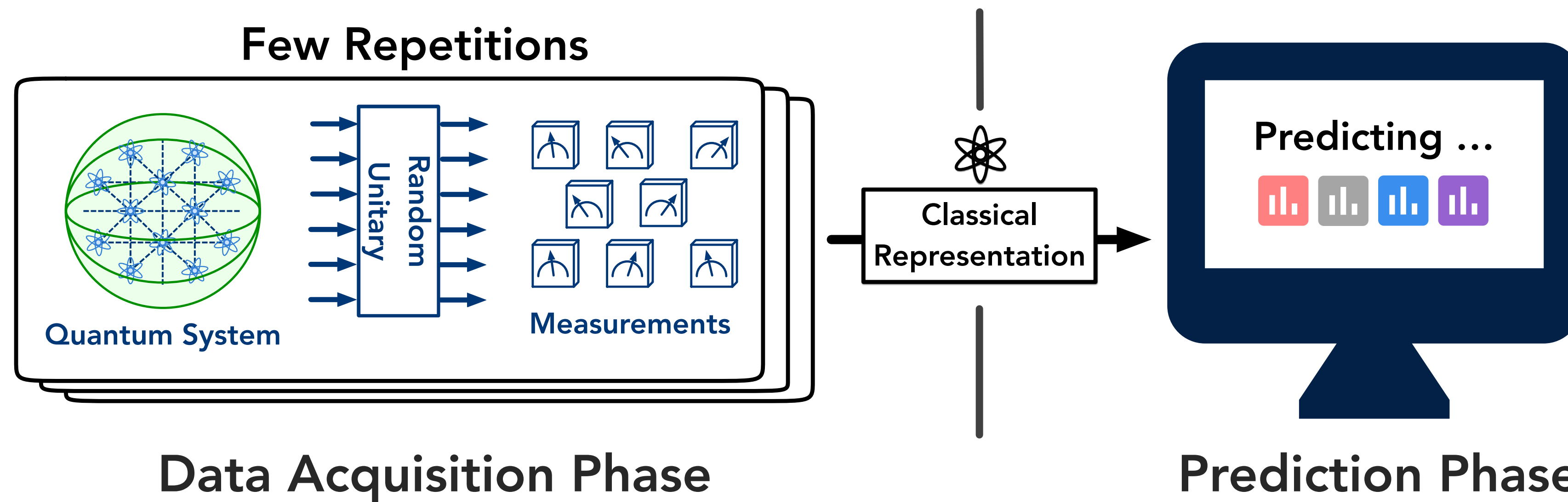
$$N = \mathcal{O}(B \log(M)/\epsilon^2) \text{ measurements.}$$

2. Subsequently, given any  $O_1, \dots, O_M$  with  $B \geq \max \|O_i\|_{\text{shadow}}^2$ , the procedure can use the classical representation to predict  $\hat{o}_1, \dots, \hat{o}_M$ , where  $|\hat{o}_i - \text{tr}(O_i \rho)| < \epsilon$ , for all  $i$ .

For example:

- $M = 10^6$ ,  $B = 1$ , then naively we need  $10^6/\epsilon^2$  measurements.
- This theorem shows that we only need  $6 \log(10)/\epsilon^2$  measurements.

# Classical shadow formalism



Algorithm for predicting  $\text{tr}(O\rho)$ : (median-of-means)

Compute  $X_i = \text{tr}(O\mathcal{M}^{-1}(|s_i\rangle\langle s_i|))$ ,  $\forall i = 1, \dots, N$ .

Predict  $\hat{o} = \text{median} \left( \frac{1}{N/K} \sum_{i=1}^{N/K} X_i, \dots, \frac{1}{N/K} \sum_{i=N-N/K+1}^N X_i \right)$ .

# Classifying quantum phases: ML

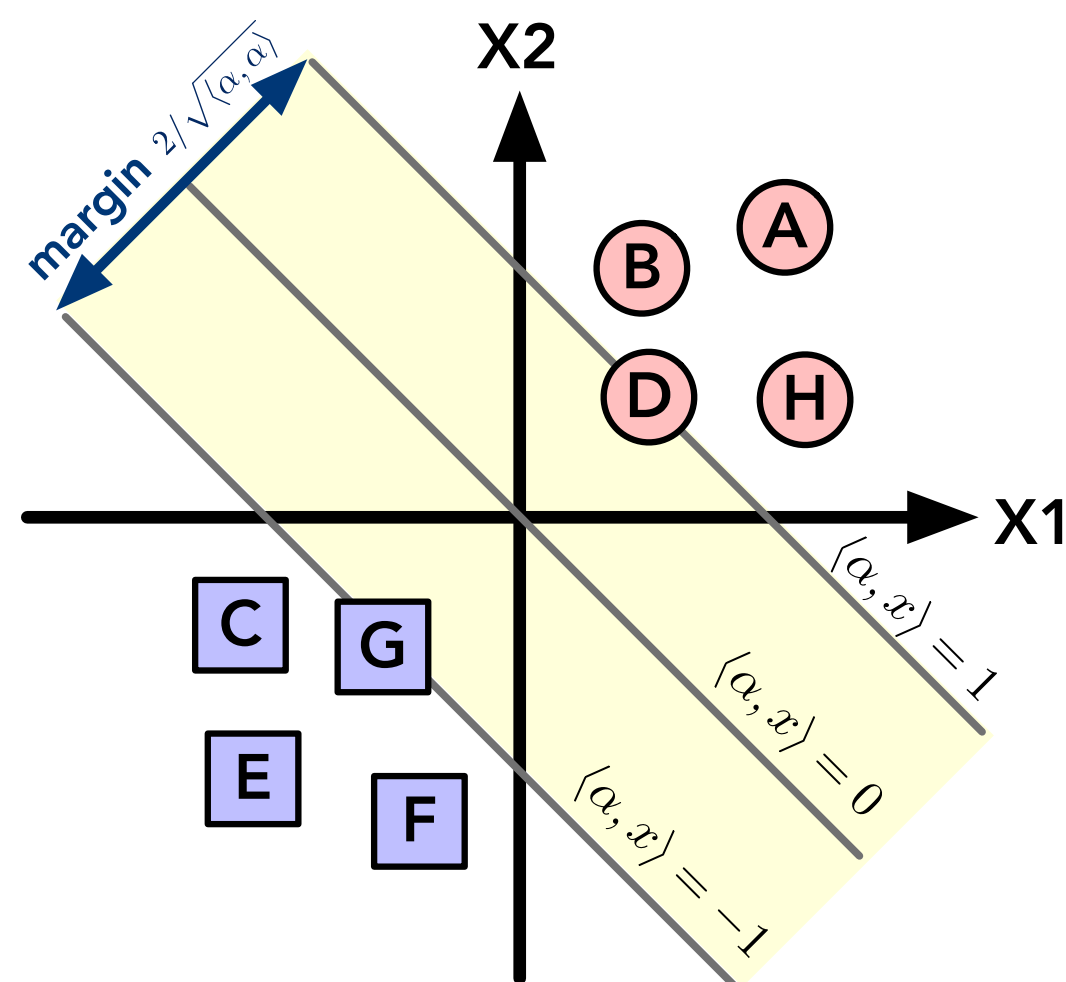
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Classical shadow formalism

- To do so, we consider learning a linear function  $\phi$  that approximates the reduced density matrix for subsystem  $i_1, \dots, i_r$ .

$$\phi^{(\text{shadow})}(S_T(\rho)) = \lim_{D,R \rightarrow \infty} \bigoplus_{d=0}^D \sqrt{\frac{\tau^d}{d!}} \left( \bigoplus_{r=0}^R \sqrt{\frac{1}{r!} \left(\frac{\gamma}{n}\right)^r} \bigoplus_{i_1=1}^n \dots \bigoplus_{i_r=1}^n \text{vec} \left[ \frac{1}{T} \sum_{t=1}^T \bigotimes_{\ell=1}^r \sigma_{i_\ell}^{(t)} \right] \right)^{\otimes d}.$$

- It consists of arbitrarily-large  $r$ -body reduced density matrices and arbitrarily-high-degree expansion.



## Proposition 2

Consider two distinct topological phases A and B.

No (local/global) observable  $O$  exists such that

$$\text{Tr}(O\rho_A) > 0, \forall \rho_A \in \text{phase A}, \quad \text{Tr}(O\rho_B) \leq 0, \forall \rho_B \in \text{phase B}.$$

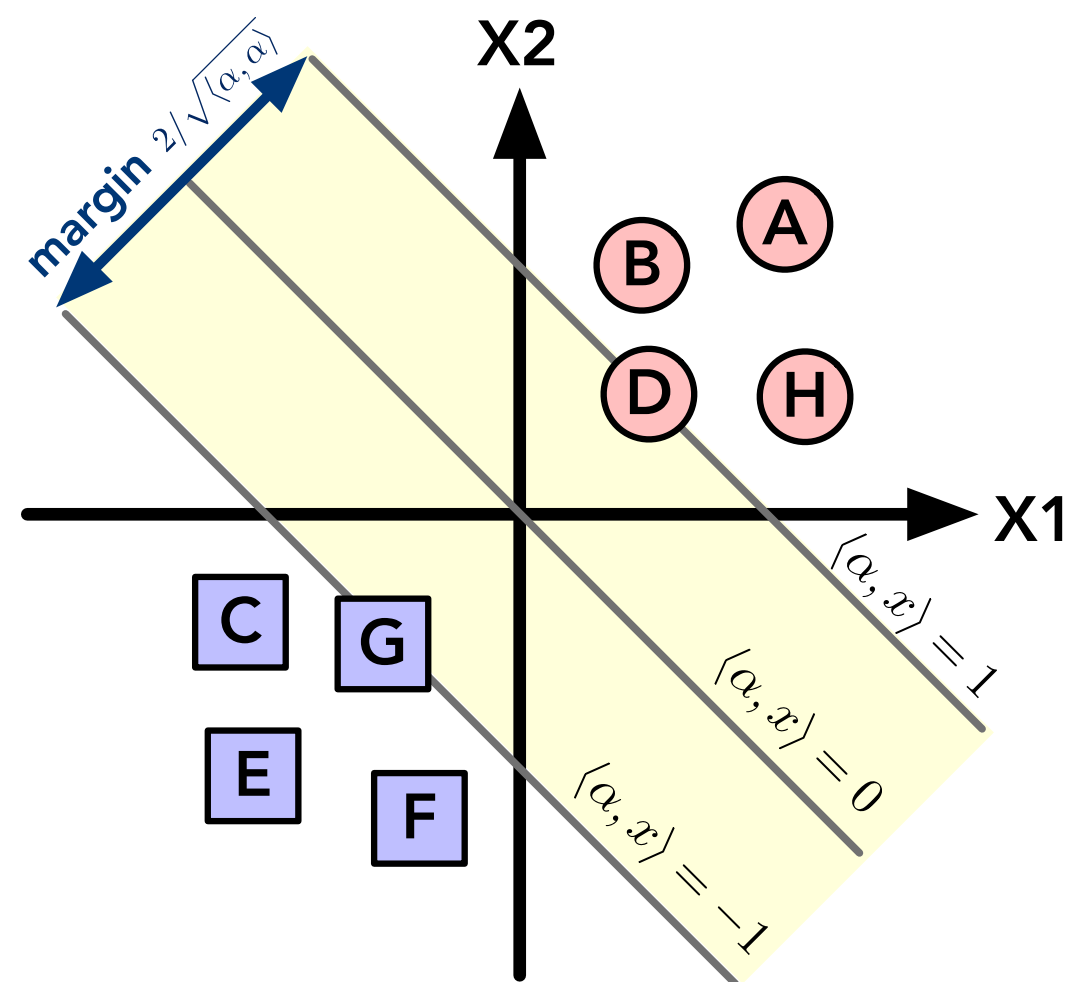
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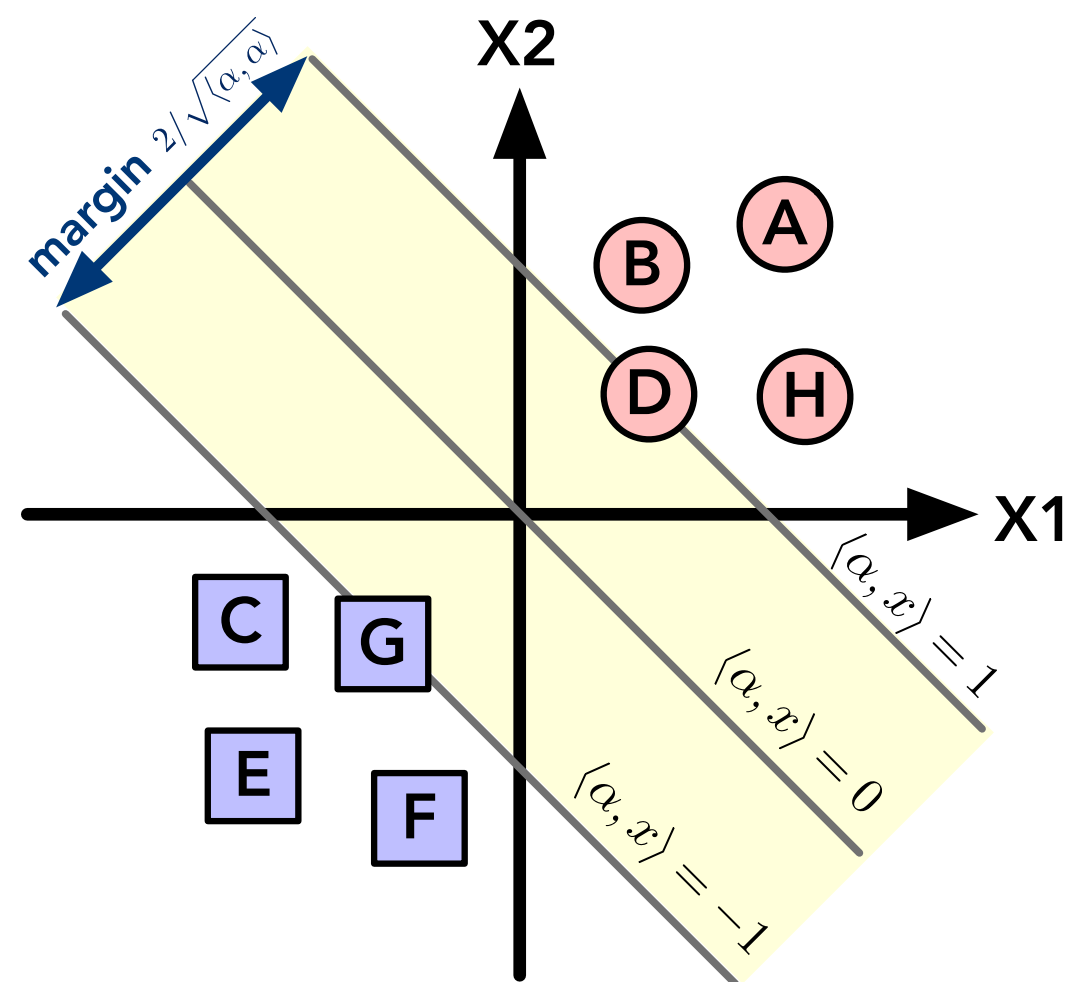
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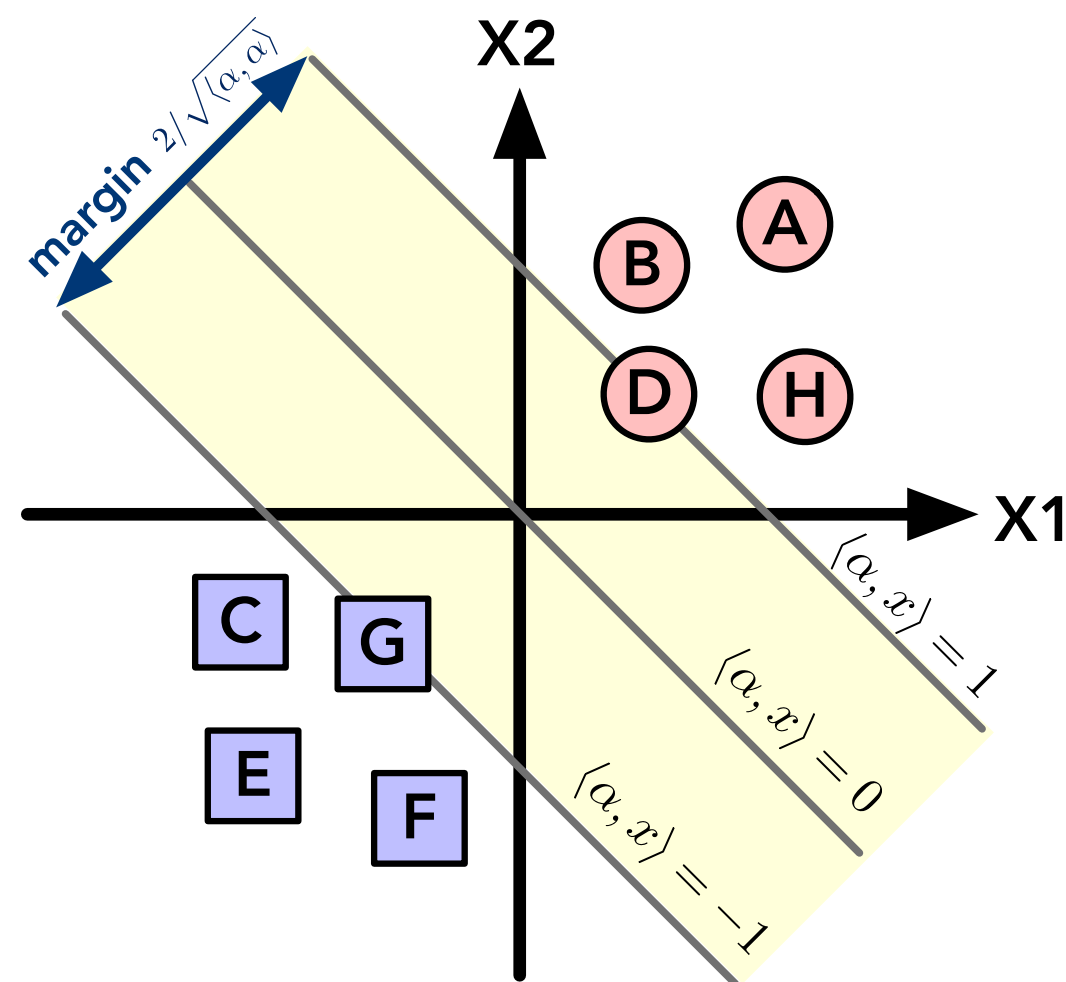
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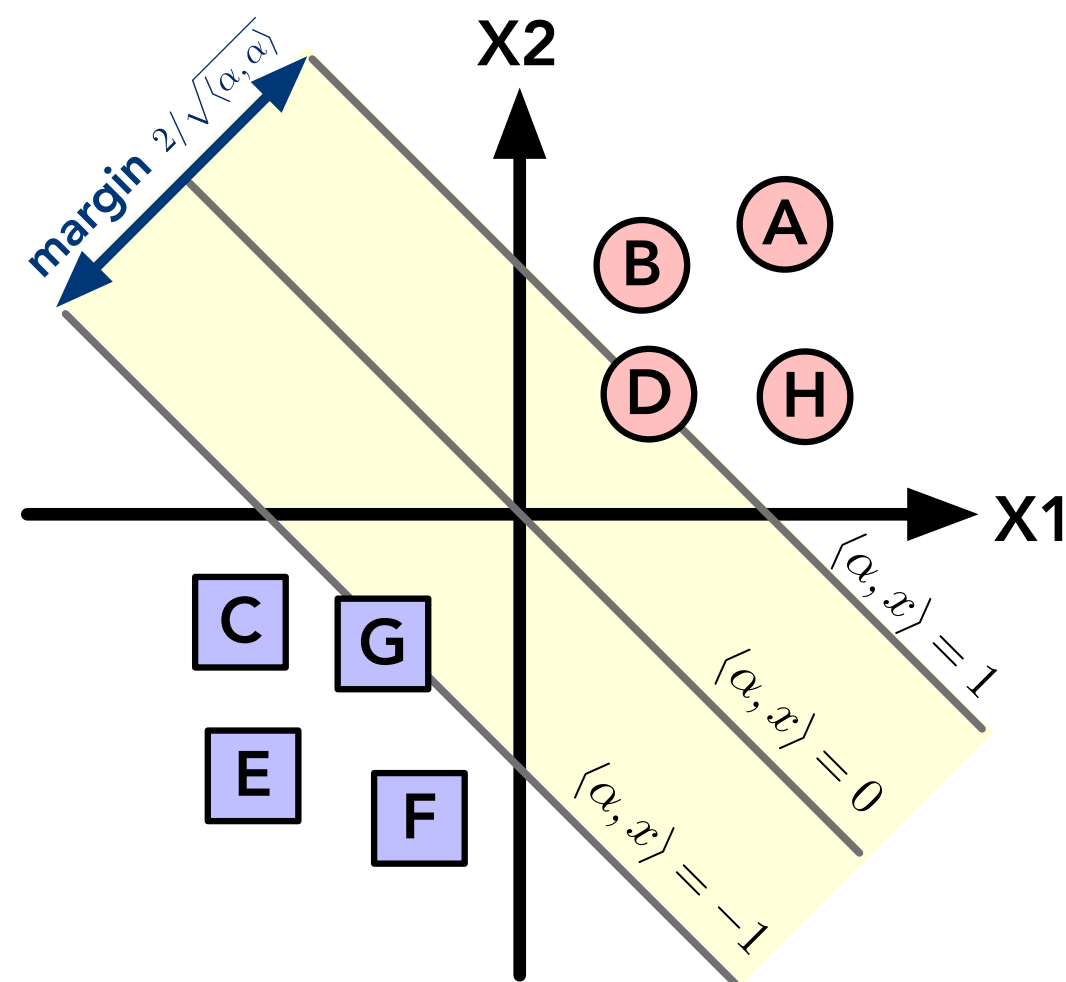
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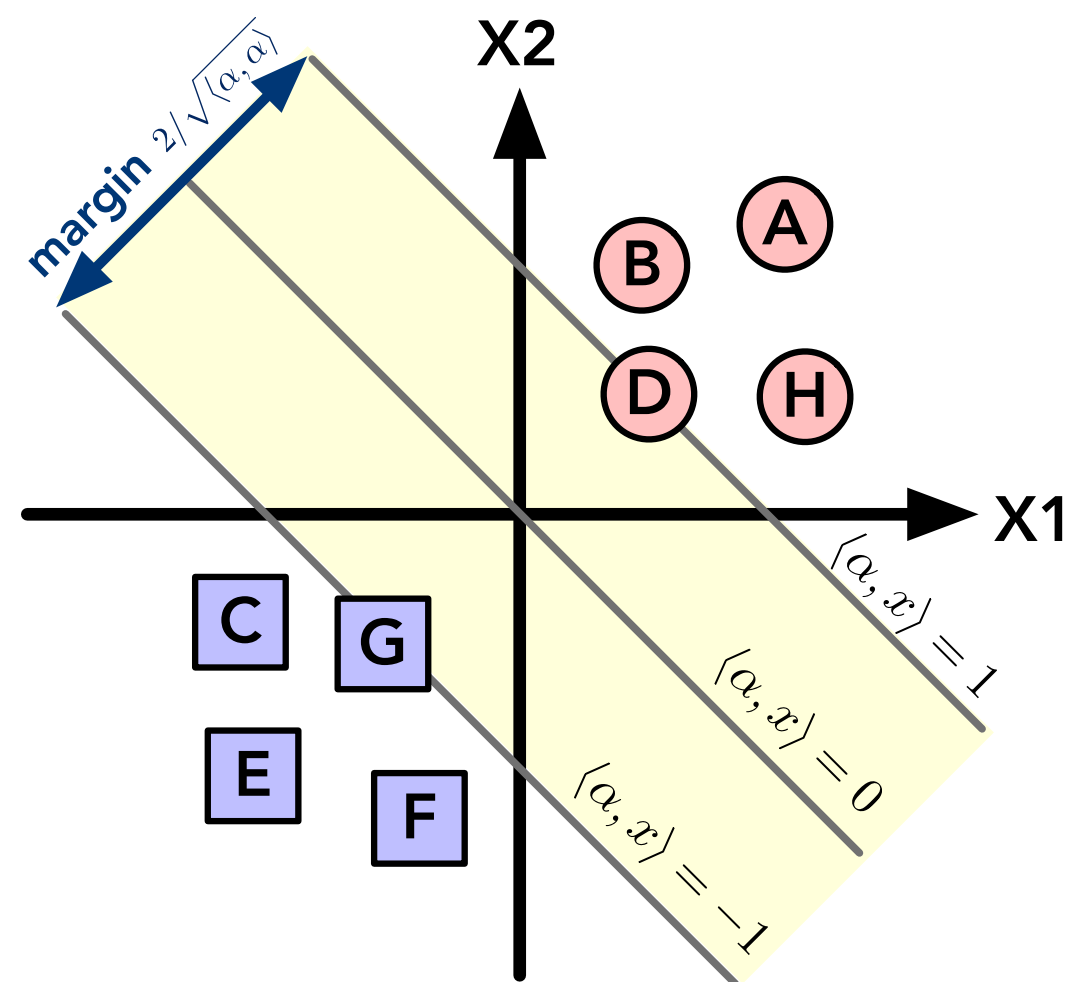
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