Classical algorithms for quantum mean values

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Joint work with David Gosset, Alexander Kliesch, Robert Koenig, Ramis Movassagh, Eugene Tang

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Why quantum many-body problems are hard to solve classically?

- Exponentially large Hilbert space
- Entanglement
- Sign problem
- Glassiness







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Taming the exponential scaling: variational algorithms

Minimize the energy of a Hamiltonian describing a system of n qubits over a class of variational states that depend only on poly(n) parameters.

Variational Quantum Eigensolver (VQE)



Simulation of electronic structure in molecules **Peruzzo et al 2014, Kandala et al 2017** Quantum Approximate Optimization Algorithm (QAOA) **Farhi, Goldstone, Gutmann 2014**

Robust agains systematic unitary errors; random errors can be mitigated Temme, SB, Gambetta 2017

Limitations of VQE

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Algorithmic limitations: the number of variational parameters must be small enough to enable efficient energy minimization. Large-scale VQE with an extensive number of variational parameters may give rise to intractable optimization problems.



Variational Quantum Eigensolver (VQE)

Electronic structure simulation for chemistry or material science



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Expected value of a multi-qubit Pauli operator can be inferred by measuring each qubit in X or Y or Z basis and classically multiplying the measured outcomes:

$$\langle \psi(\theta) | X \otimes Y \otimes Z \otimes \cdots | \psi(\theta) \rangle = \mathfrak{E}(-1)^{m_1 + m_2 + m_3 + \cdots}$$

$$average over many experiments$$



Quantum Mean Value problem

Suppose U is a low-depth quantum circuit on n qubits and $\varepsilon > 0$ is the error tolerance. Given a tensor product observable

 $P = P_1 \otimes P_2 \otimes \cdots \otimes P_n, \qquad ||P_i|| \le 1$

approximate $QMV \equiv \langle 0^n | U^{\dagger}PU | 0^n \rangle$ within an additive error ε .

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Do we really need a quantum computer to solve the problem ?

variational circuit	best previously known [1,2]	our algorithm
constant depth, 2D	$2^{O(n^{1/2})}$	$O(n\varepsilon^{-2})$
constant depth, 3D	$2^{O(n^{2}/3)}$	$\varepsilon^{-2} 2^{0(n^{1/3})}$

[1] Aaronson and Chen, arXiv:1612.05903

[2] Markov and Shi, SIAM J. on Comp. (2008), quant-ph/0511069

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No quantum advantage if variational circuits are 2D and constant depth !

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constant depth, all-to-all connectivity	$O(n2^n)$	$2^{\tilde{O}(\sqrt{n\log(1/\varepsilon)})}$
	Caveat: we	can only compute

aveat: we can only compute the magnitude of QMV

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general constant depth	$O(n2^n)$	$2^{\tilde{O}(\sqrt{n\log(1/\varepsilon)})}$
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Van den Nest, "Simulating quantum computers with probabilistic methods", arXiv:0911.1624

$\langle \Psi_1 | \Psi_2 \rangle = ?$

Efficient approximation algorithm for computationally tractable states such that

- Amplitudes (x|Ψ_i) are easy to compute
 Distributions |(x|Ψ_i)|² are easy to sample

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Example: Matrix Product States with a small bond dimension are computationally tractable for any order of qubits.

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The algorithm approximates the inner product with a small additive error. Computing the inner product exactly or with a small multiplicative error is #P-hard.

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Monte Carlo approach:

$$\langle \Psi_1 | \Psi_2 \rangle = \sum_{x} |\langle x | \Psi_1 \rangle|^2 \cdot \frac{\langle x | \Psi_2 \rangle}{\langle x | \Psi_1 \rangle}$$

mean value of a function $f(x)$
over a distribution $p(x)$
 $Var(f) \le 1$

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$$\langle \Psi_1 | \Psi_2 \rangle \approx \frac{1}{M} \sum_{i=1}^{M} f(x^i)$$

empirical mean value of $f(x)$
over M samples from $p(x)$
Approximation error: $\epsilon \sim M^{-1/2}$

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dressed observable

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Locality



Simulation within a single light cone



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 $Q_i Q_j = Q_j Q_i$

Locality

 $\mathbf{QMV} = \left\langle 0^n \middle| U^{\dagger}(P_1 \otimes P_2 \otimes \cdots \otimes P_n) U \middle| 0^n \right\rangle = \left\langle 0^n \middle| Q_1 Q_2 \cdots Q_n \middle| 0^n \right\rangle$

Step 2: coarse grain the lattice such that each dressed observable Q_i acts on a 2x2 block.



Now each lattice site has local dimension $D = 2^{O(d^2)}$

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permutation of *n* qubits





MPS with bond dimension $\chi \le D^2 = 2^{O(d^2)}$

 $\mathbf{QMV} = \langle 0^n | Q_1 Q_2 \cdots Q_n | 0^n \rangle = \langle \Psi_A | W | \Psi_B \rangle$ permutation of *n* qubits Inner product of computationally tractable states. Apply Van den Nest algorithm.





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Maximize a classical cost function $C : \{1, -1\}^n \to \mathbb{R}$

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$$G = (V, E)$$

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1) Promote the cost function to a quantum Hamiltonian: $C = \sum_{z} C(z) |z\rangle \langle z |$

2) Maximize expected energy $\langle \psi | \mathcal{C} | \psi \rangle$ over variational states ψ with a few parameters.

3) Measure the optimal state ψ to obtain a classical solution $z \in \{1, -1\}^n$. Mean value of C(z) equals $\langle \psi | C | \psi \rangle$.

Level-*p* variational state:

$$|\psi(\beta,\gamma)\rangle = \prod_{j=1}^{p} \exp[-i\beta_{j}(X_{1} + \dots + X_{n})]\exp[-i\gamma_{j}C]|+ + \dots + \rangle$$

 $|+\rangle \sim |0\rangle + |1\rangle$

Variational parameters: $\beta, \gamma \in \mathbb{R}^p$

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Trotterized version of the Adiabatic Quantum Computation for p = poly(n)
Quantum Approximate Optimization Algorithm [Farhi, Goldstone, Gutmann 2014]

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Trotterized version of the Adiabatic Quantum Computation for p = poly(n)

Reasons for keeping the level p small:

- Non-linear optimization over β and γ is hard
- Need to keep the circuit depth small for near-term implementation

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No-go theorems:

Level-1 QAOA is inferior to local classical optimizers for bounded-degree graphs [Hastings arXiv:1905.07047]

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> Can we overcome these limitations ? New idea: variable elimination and recursive QAOA [SB, Kliesch, Koenig, Tang, arXiv:1910.08980]

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- 4. Impose a constraint $z_b = \operatorname{sign}(M_{a,b})z_a$ and eliminate z_b from the cost function $J_{b,c}z_bz_c \leftarrow J_{b,c}\operatorname{sign}(M_{a,b})z_az_c$

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We get a new Ising-like cost function C'(z) that depends on n-1 variables.















Numerical simulation of level-1 QAOA and RQAOA for 100 qubits. Cost function: random-bond Ising model on the complete graph. SDP: Goemans-Williamson semidefinite programming relaxation with 100 rounding trials.



$$C(z) = \sum_{a < b} J_{a,b} z_a z_b$$



Variables: vertex colors $z_a \in \{0,1,2\}$

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SDP relaxation algorithm: $C(z) \ge 0.836 \cdot C_{max}$ [Klerk, Pasechnik, Warners 2004] Numerical simulations of level-1 QAOA and RQAOA for 50 qutrits and the MAX 3-CUT cost function. We consider a random ensemble of 3-colorable dense graphs with 50 vertices randomly partitioned into red/blue/green.





add random red-blue, red-green, blue-green edges with probability ½ for each pair of vertices Bad news: no quantum advantage for level-1 RQAOA

Bad news: no quantum advantage for level-1 RQAOA

The only "quantum" step: computing quantum mean values

 $M_{a,b} = \langle \psi(\beta,\gamma) | Z_a Z_b | \psi(\beta,\gamma) \rangle$

Efficient classical algorithm for level-1 QAOA

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Quantum approximate optimization algorithm for MaxCut: A fermionic view

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 $QMV = \left\langle \psi(\beta, \gamma) \middle| O_{1,2} \middle| \psi(\beta, \gamma) \right\rangle =$



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Cancel all gates that do not touch qudits 1,2

 $\mathrm{QMV} = \left\langle \psi(\beta,\gamma) \left| \mathcal{O}_{1,2} \right| \psi(\beta,\gamma) \right\rangle =$



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Change the order of gates such that qudits 1,2 are coupled to qudits 3,4,...,n sequentially.

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$$QMV = \langle \psi(\beta,\gamma) | O_{1,2} | \psi(\beta,\gamma) \rangle = \langle + + | \Phi_n \circ \Phi_{n-1} \circ \cdots \circ \Phi_3 (\tilde{O}) | + + \rangle$$



Sequential coupling: need to compute a product of two-qudit quantum channels. Classical simulation time O(n) assuming that qudit dimension is O(1).




Our algorithm:

Arbitrary interaction graph: runtime $2^{n/2}\epsilon^{-2}$

Planar interaction graph: runtime $O(n^2 \epsilon^{-2})$. Works for qudits of constant dimension.



B-layers: single-qubit X-rotations

C-layers: diagonal two-qubit gates



Step 1: conjugate the observable O by the inner B-layers. This gives a modified two-qubit observable \tilde{O} .



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tensor product observables



 $P = P_1 \otimes P_1 \otimes \cdots \otimes P_n$



Step 3: conjugate the observable *P* by the inner B-layers.



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Step 4: express the QMV as the inner product of computationally tractable states. Approximate the inner product $\langle \Psi_1 | \Psi_2 \rangle$ using Van den Nest algorithm.



Claim: any amplitude $\langle \Psi_1 | x \rangle$ can be computed in time roughly $2^{n/2}$.

 Ψ_1



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- The C-layer includes only diagonal two-qubit gates
- It cannot entangle a qubit initialized in a basis state
- Half of all qubits remains unentangled



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Approximating the inner product $QMV = \langle \Psi_1 | \Psi_2 \rangle$ by Monte-Carlo takes time $2^{n/2} \epsilon^{-2}$



Suppose the C-layer only includes nearest-neighbor gates on a planar graph. Can we pick a good partition of qubits making Ψ_1 and Ψ_2 less entangled ? **Theorem**. Suppose G = (V, E) is a planar graph. There exists a partition $V = V_1 V_2$ such that the subgraphs of G induced by V_1 and V_2 have treewidth at most 2. Such partition can be efficiently computed.



DeVos et al, "Excluding any graph as a minor allows a low tree-width 2-coloring" (2004)

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Low-treewidth tensor networks are easy to contract. [Markov and Shi 2004] **Recap:** we obtained a polynomial-time classical algorithm for approximating

 $\mathbf{QMV} = \langle \psi(\beta, \gamma) | Z_a Z_b | \psi(\beta, \gamma) \rangle$

for level-2 QAOA states with the Ising-like cost function on any planar graph.

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Bonus feature: RQAOA preserves planarity



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Corollary: level-2 RQAOA on planar graphs can be simulated classically in polynomial time.

Summary

- Variational quantum algorithms based on constant-depth geometrically local circuits in 2D can be simulated classically in linear time.
- Large-scale classical simulation of level-1 RQAOA is reported. Classical simulation of level-2 RQAOA is a work in progress.

Open problems

- Establish classical hardnes of approximating quantum mean values for low-depth circuits or low-level QAOA
- Rigorous bounds on the performance of RQAOA. More general variable elimination methods.