

# Correlations, area laws and stability of open and thermal quantum many-body systems

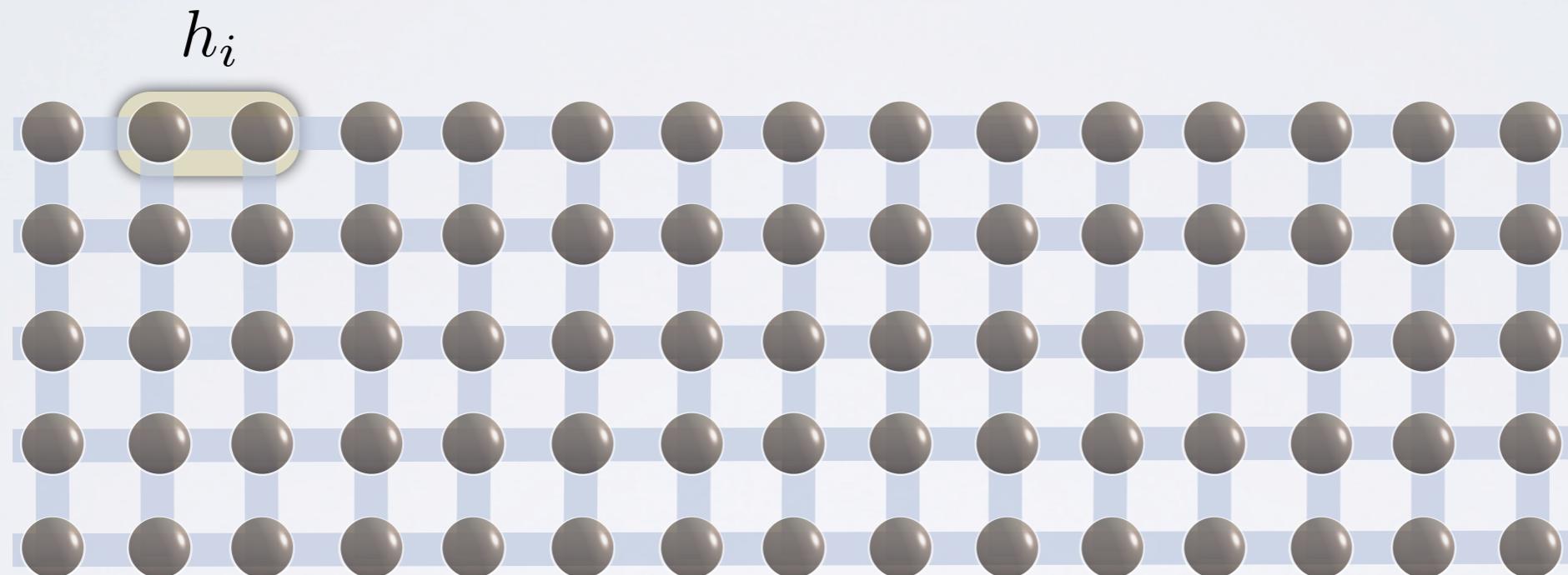


Jens Eisert, FU Berlin



Complexity meets Condensed Matter, Simons Institute for the Theory of Computing, March 2014  
Joint work with Michael Kastoryano, Martin Kliesch, Christian Gogolin, Arnau Riera

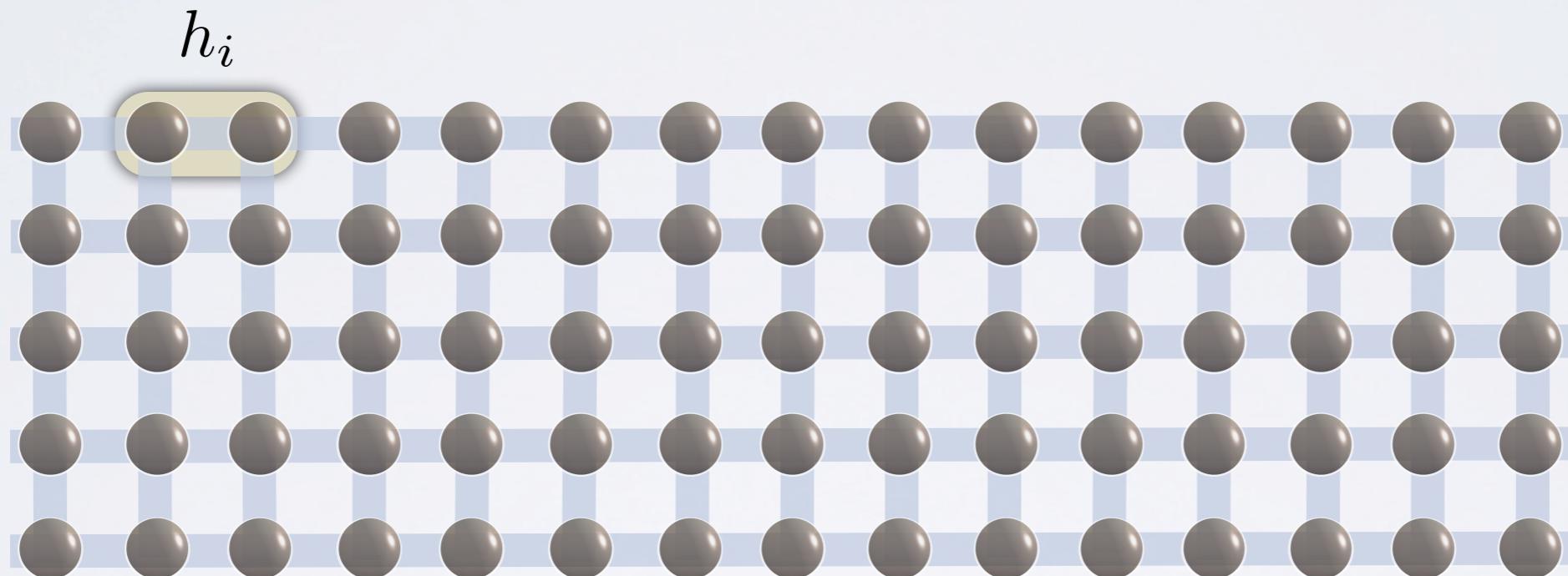
# Hamiltonian complexity



- Can ground states of “natural” quantum systems be described succinctly?
- Does the exponential complexity of general quantum systems persist at high temperature?
- Is the scientific method sufficiently powerful to understand general quantum systems?

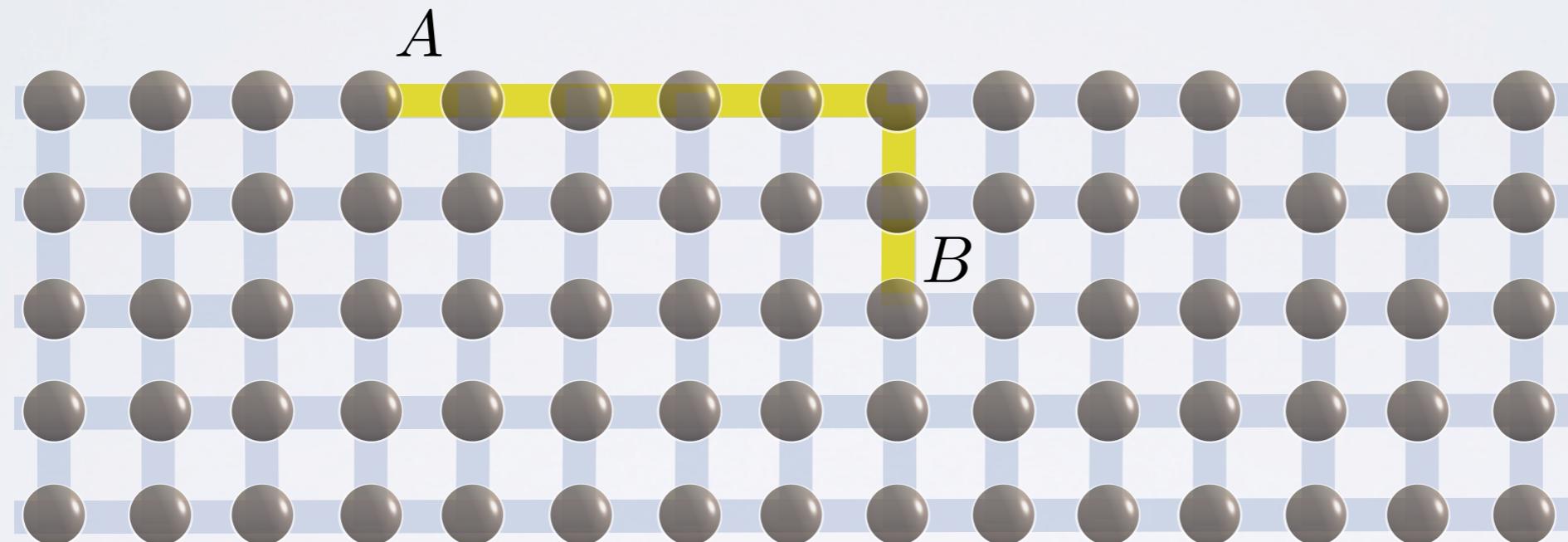
- **Local Hamiltonian** problem is QMA-complete
- Steps towards a **quantum PCP** theorem (Matt's and Dorit's talks)

# Hamiltonian complexity



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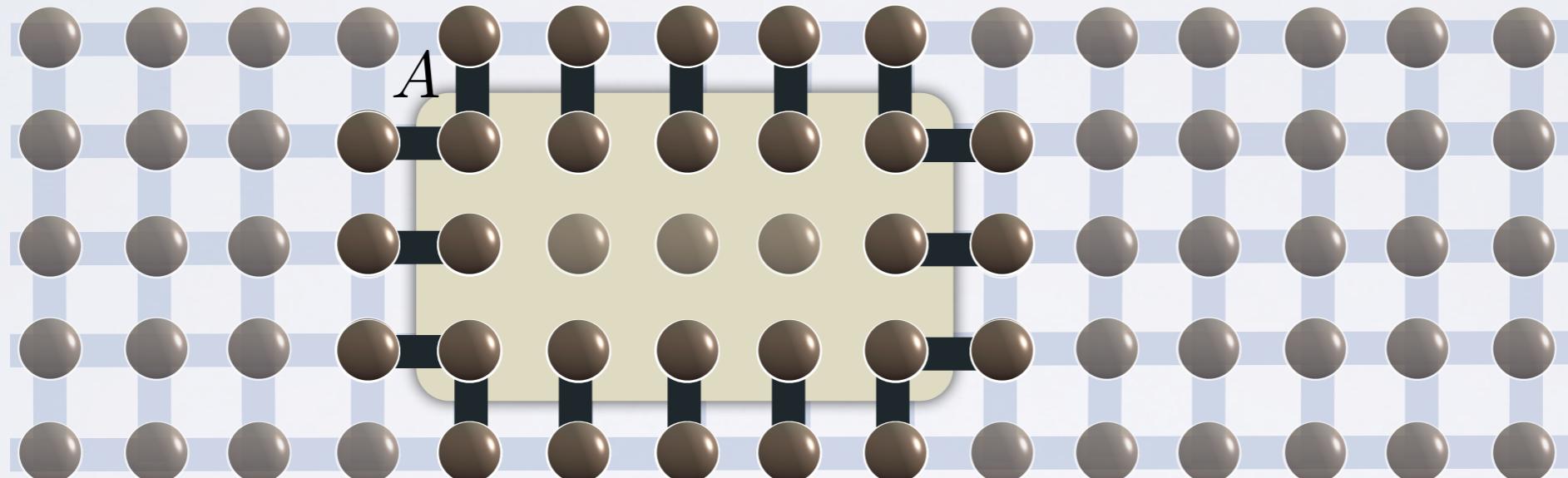
# Ground states of local gapped models



- **Energy gap**  $\Delta(H) = E_1 - E_0 > 0$
- Ground states of gapped models have **exponentially decaying correlations**

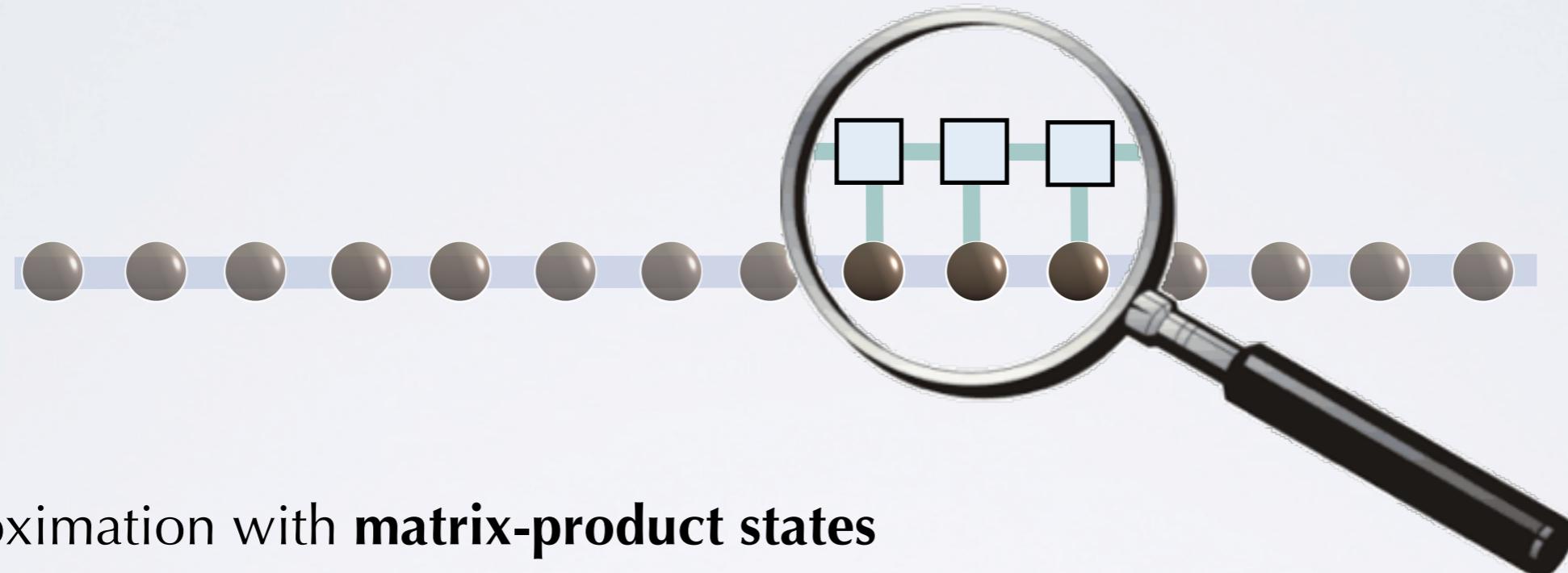
- Proof based on Lieb-Robinson bounds
  - Hastings, Koma, Commun Math Phys 265, 781 (2006)
  - Nachtergael, Sims, Commun Math Phys 265, 119 (2006)
- Combinatorial proof (detectability lemma)

# Area laws



- **Area laws for the entanglement entropy**  $S(\rho_A) = O(|\partial A|)$
- Proven for gapped quasi-free **bosonic and fermionic systems** in any dimension,  
**1D gapped local models** and ones with exponentially **decaying correlations**

# Matrix-product states and efficient descriptions

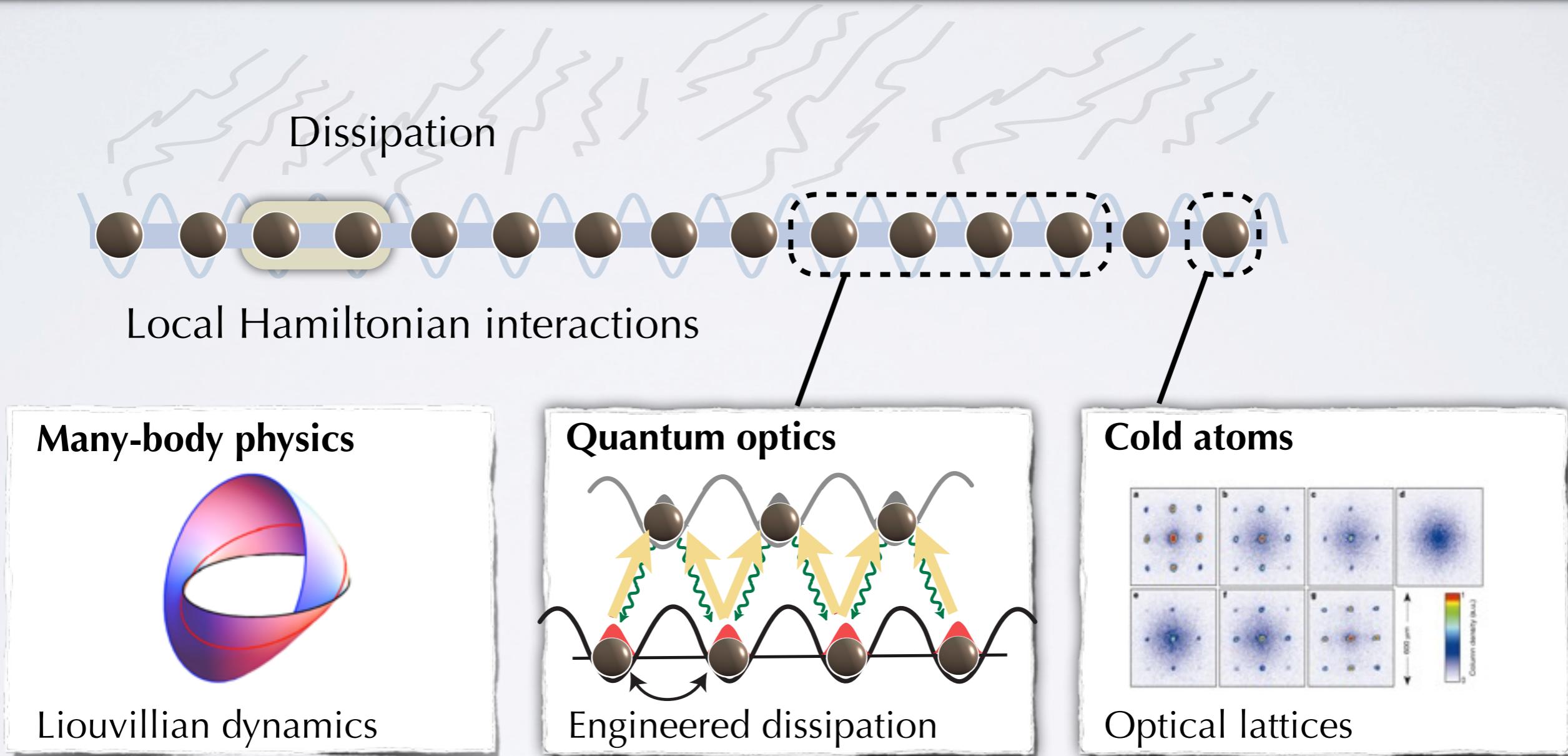


- Approximation with **matrix-product states**
- Polynomial-time algorithm for ground states of 1D gapped local Hamiltonians

Landau, Vazirani, Vidick, arXiv:1307.5143

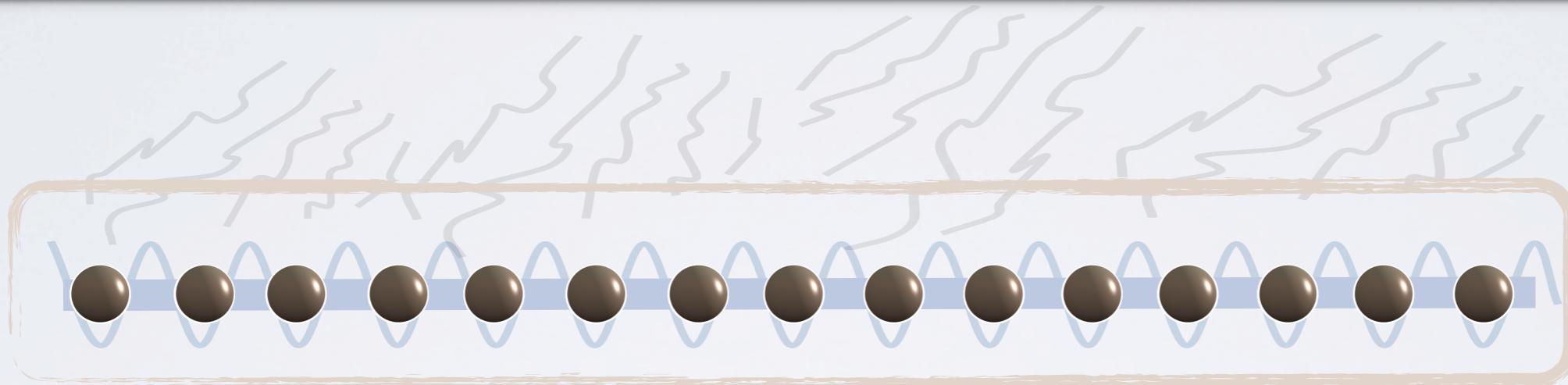
- **This talk:** Correlations in **thermal** and **open** quantum many-body systems

# Open quantum many-body systems



Diehl, Micheli, Kantian, Kraus, Buechler, Zoller, Nature Physics 4, 878 (2008)  
Kraus, Diehl, Micheli, Kantian, Buechler, Zoller, Phys Rev A 78, 042307 (2008)  
Verstraete, Wolf, Cirac, Nature Physics 5, 633 (2009)  
Eisert, Prosen, arXiv:1012.5013  
Bravyi, Chesi, Loss, Terhal, New J Phys 12, 025013 (2010)  
Kastoryano, Wolf, Eisert, Phys Rev Lett 110, 110501 (2013)

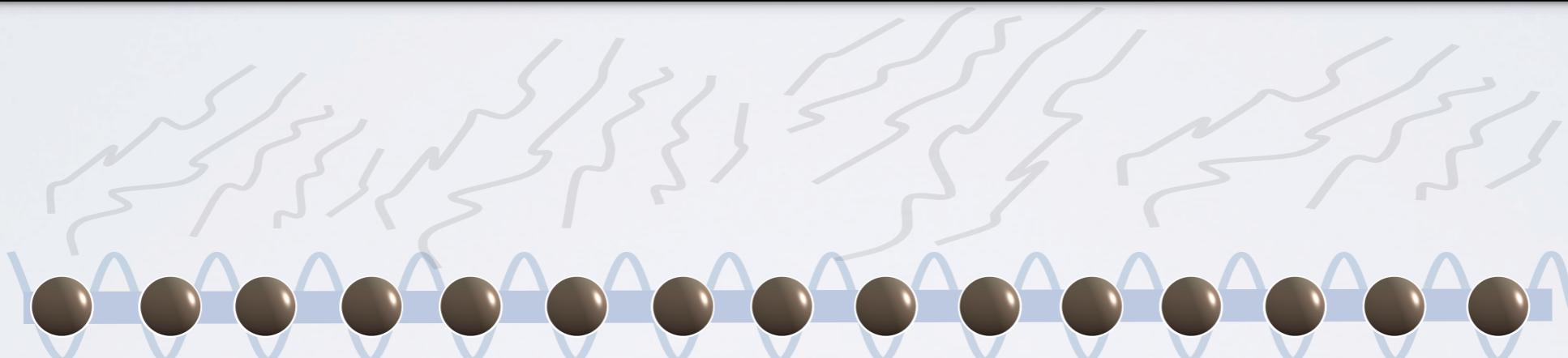
# Open quantum many-body systems



- Dissipative quantum phase transitions, noise-driven criticality, topological order
- Dissipative quantum computing
- Dissipative passive quantum memories?

Diehl, Micheli, Kantian, Kraus, Buechler, Zoller, Nature Physics 4, 878 (2008)  
Kraus, Diehl, Micheli, Kantian, Buechler, Zoller, Phys Rev A 78, 042307 (2008)  
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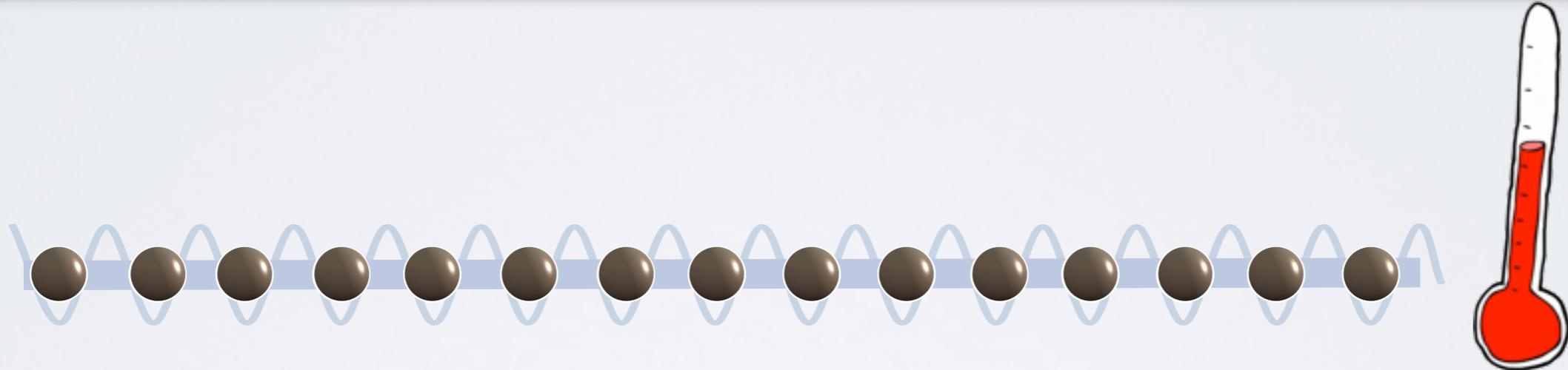
# Questions of the rest of talk



- "Liouvillian complexity", resembling Hamiltonian complexity

- How is closing of Liouvillian gaps related to clustering of correlations?
- Area laws in dissipative systems? Stability? Topological dissipative memories?

# Questions of the rest of talk



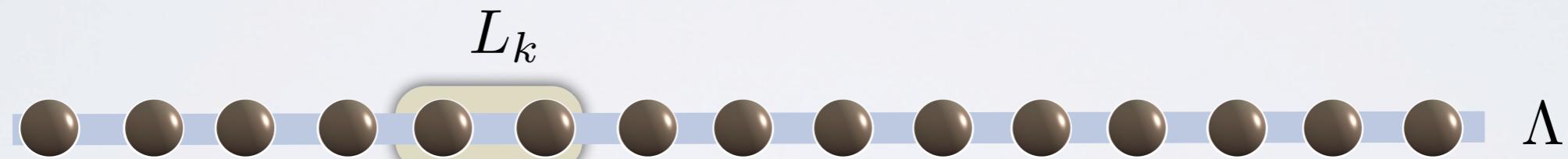
- Thermal states at high temperatures

- Is temperature intensive/local?
- Correlations in thermal many-body states?
- Computational complexity of computing expectation values?

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# Correlations in open many-body systems

# Local Liouvillian setting

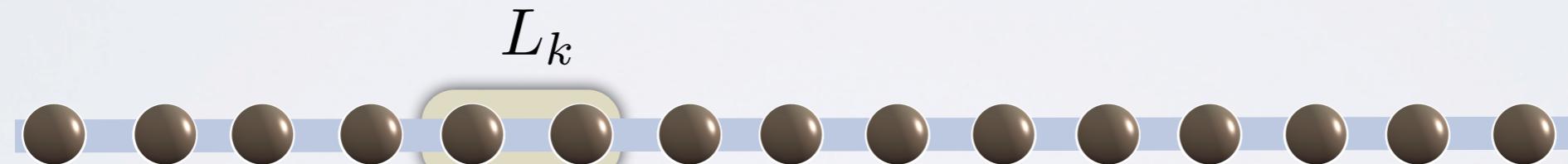


- **Liouvillian setting**, reflecting Markovian dynamics

$$\frac{d}{dt}\rho = \mathcal{L}(\rho) = i[H, \rho] + \sum_k \left( L_k \rho L_k^\dagger - \frac{1}{2} \{L_k^\dagger L_k, \rho\} \right)$$

- Effective system dynamics
- Lindblad operators are **geometrically local on some graph**  $\Lambda$
- **Bounded** interactions, i.e.,  $\|L_k\| < K$  for all  $k$

# Stationary states and mixing properties



- Role of ground state taken over by **stationary state**  $\sigma$ , satisfying

$$\mathcal{L}(\sigma) = 0$$

here often taken to be full rank (primitive), with detailed balance

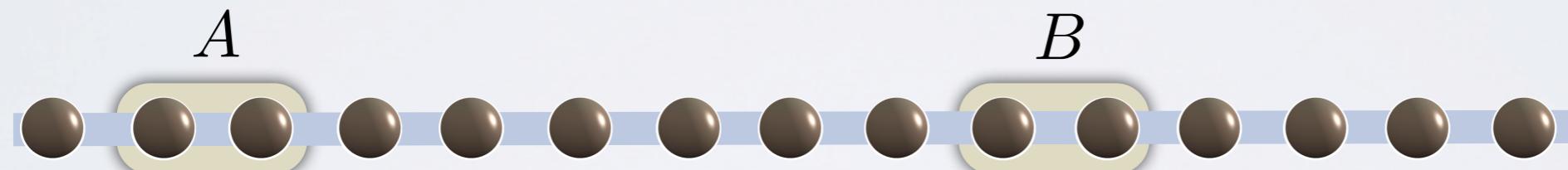
- **Mixing properties:** For any primitive local Liouvillian,

$$\|e^{t\mathcal{L}}(\rho_0) - \sigma\|_1^2 \leq \|\sigma^{-1}\| e^{-2\lambda t}$$

for any initial state  $\rho_0$

- $\lambda$  is the **Liouvillian gap**, resembling the Hamiltonian gap in closed systems

# Correlation measures



- **Covariance:** For arbitrary regions  $A, B \subset \Lambda$

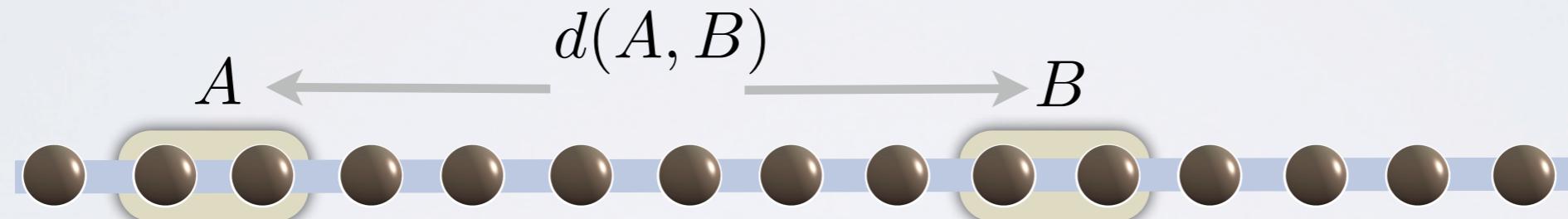
$$C_\rho(A, B) = \sup_{\|f\|=\|g\|=1} |\text{tr}((f \otimes g)(\rho_{A,B} - \rho_A \otimes \rho_B))|$$

"largest connected correlation function"

- Related to other **standard correlation measures**

- *trace distance*  $T_\rho(A, B) := \|\rho_{A,B} - \rho_A \otimes \rho_B\|_1$
- *mutual information*  $I_\rho(A, B) := S(\rho_{A,B} \| \rho_A \otimes \rho_B)$

# Clustering of correlations and gaps?

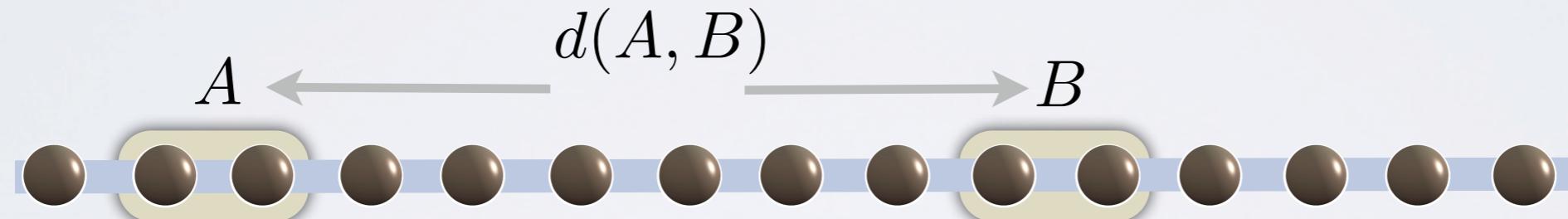


- Gapped Hamiltonians, away from **phase transitions**, show **clustering of correlations**

$$C_\rho(A, B) \leq C e^{-d(A, B)/\xi}$$

- How about gapped Liouvillians?

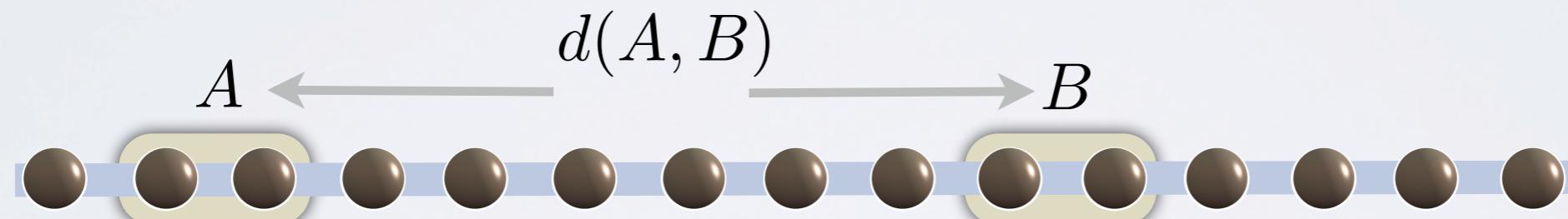
# Clustering of correlations and gaps



- **Clustering of correlations:**  $A, B \subset \Lambda$  non-overlapping subsets, consider local, bounded Liouvillian with stationary state  $\sigma$ , gap  $\lambda$ , and Lieb-Robinson velocity  $v$ . Then there ex constant  $c > 0$ , such that

$$C_\sigma(A, B) \leq c d(A, B)^{\mathcal{D}-1} e^{-\frac{\lambda d(A, B)}{v+2\lambda}}$$

# Flavour of (simple) proof



Hölder's inequality and mixing time tools:  
Variational characterisation of gap, ...  $\leq \|f\| \|g\| e^{-2t\lambda}$

Set  $f_t = e^{t\mathcal{L}^*}(f)$ , then

$$|\text{Cov}_\sigma(f, g)| \leq |\text{Cov}_\sigma(f_t, g_t)| + |\text{Cov}_\sigma(f_t, g_t) - \text{Cov}_\sigma(f, g)|$$

Choose suitable  $t$

- **Dissipative Lieb-Robinson bound:** For observables  $f, g$  supported on  $A, B \subset \Lambda$ , respectively,

$$\|(fg)_t - f_t g_t\| \leq C d(A, B)^{\mathcal{D}-1} \|f\| \|g\| e^{vt - d(A, B)/2}$$

for all  $t \geq 0$ , where  $v$  is the Lieb-Robinson velocity and  $C > 0$  constant

Kliesch, Barthel, Phys Rev Lett 108, 230504 (2012)

Nachtergaele, Vershynina, Zagrebnov, AMS Cont Math 552, 161 (2011)

Kliesch, Gogolin, Eisert, arXiv:1306.0716

Kastoryano, Eisert, J Math Phys 54, 102201 (2013)

# Slightly stronger mixing tools



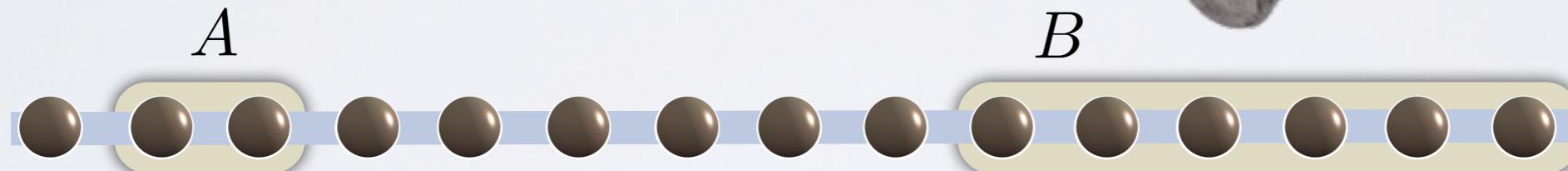
- Stronger concept of mixing, based on **Log-Sobolev constant**
- Log-Sobolev-constant  $\alpha$  bounded from above by Liouvillian gap  $\lambda$
- Variational characterisation of  $\alpha$ , related to hypercontractivity

- **Mixing properties:** For any primitive local Liouvillian,

$$\|e^{t\mathcal{L}}(\rho_0) - \sigma\|_1^2 \leq 2 \log(\|\sigma^{-1}\|) e^{-2\alpha t}$$

for any initial state  $\rho_0$

# Stability results

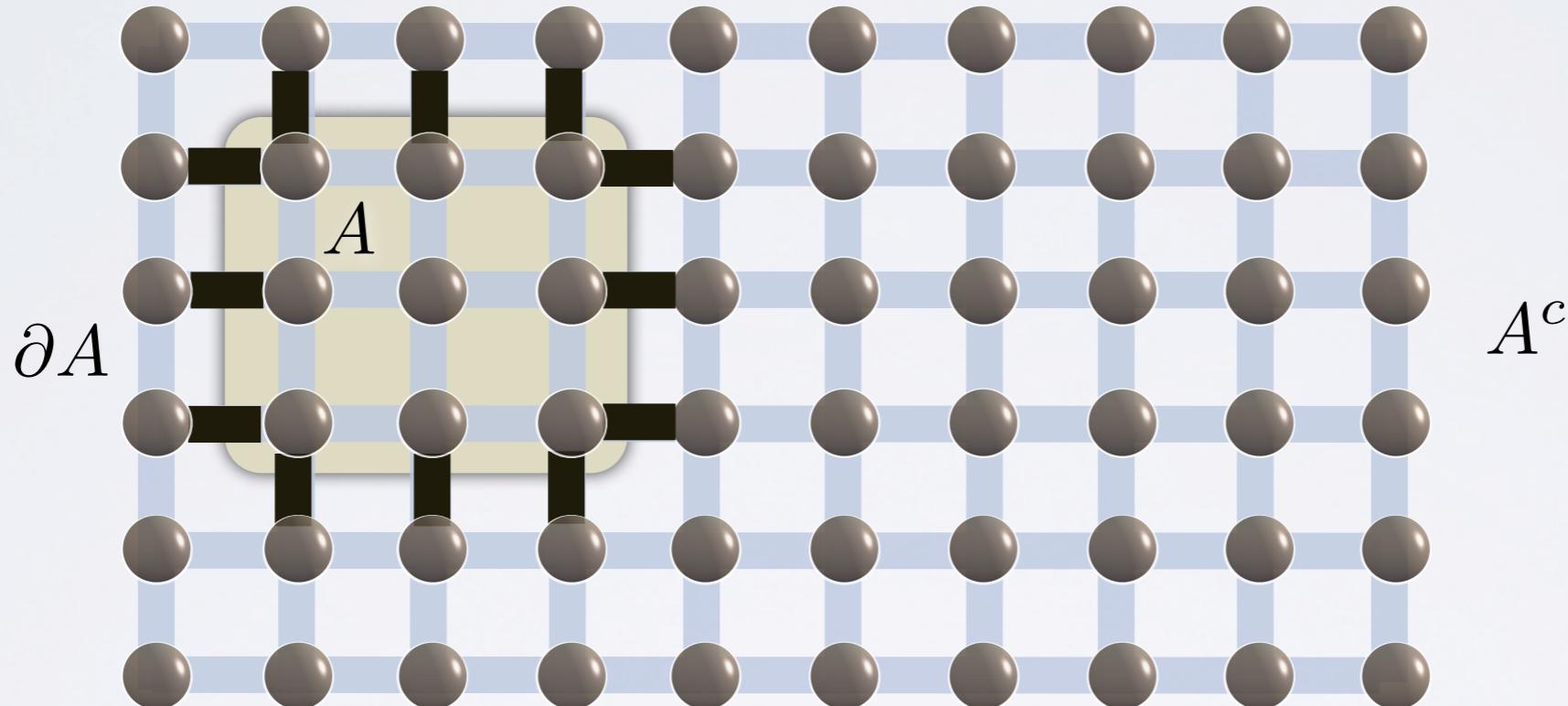


- **Local perturbations perturb locally:** Let  $\mathcal{L}$  be a local Liouvillian with Log-Sobolev-constant  $\alpha$  and stationary state  $\rho$ , let  $\mathcal{Q}_B$  be a perturbation on  $B$  only, with stationary state  $\sigma$  of  $\mathcal{L} + \mathcal{Q}_B$ , (...), then

$$\|\rho_A - \sigma_A\|_1 \leq C e^{-\alpha d(A,B)/(v+\alpha)}$$

Kastoryano, Eisert, J Math Phys 54, 102201 (2013)  
Cubitt, Lucia, Michalakis, Perez-Garcia, arXiv:1303.4744

# Area laws



- **Area law for mutual information:** (...)

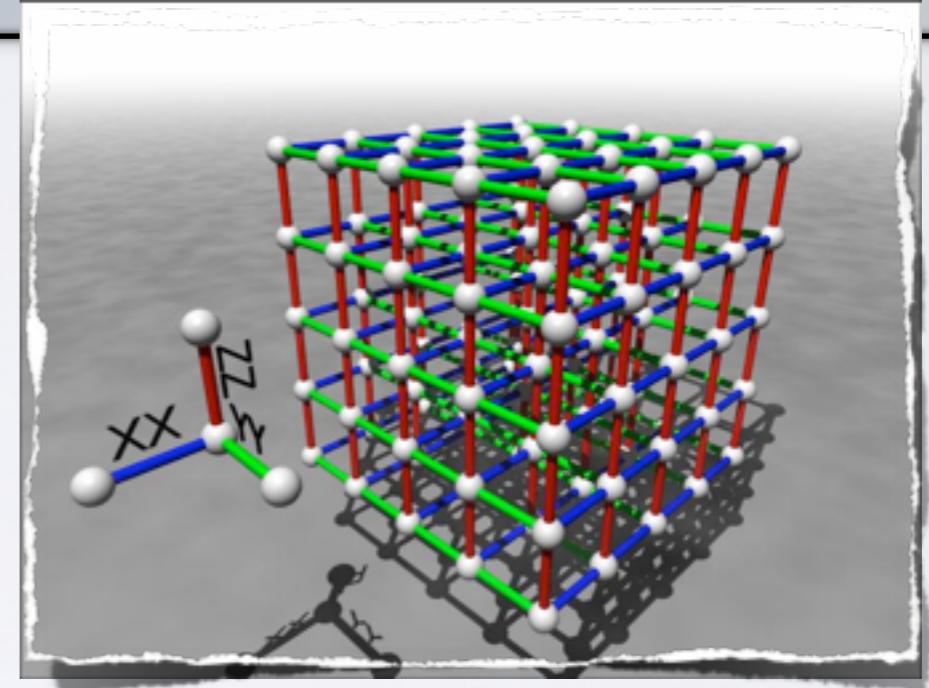
$$I_\rho(A, A^c) \leq (\gamma_1 + \gamma_2 \log \log \|\rho^{-1}\|) |\partial A| + \epsilon$$

# Mixing times and clustering of correlations

- **Lesson:** Rapidly mixing systems exhibit **exponentially clustering correlations**
- "Mixing in time related to mixing in space"
- **Liouvillian gap** (log-Sobolev constant) reminds of **Hamiltonian gap**
- Two **different regimes**, with quite different implications
- Quantum feature, difference **absent classically**

# Quantum memories, topological order and mixing times

- Optimal **dissipative encoders for toric codes**

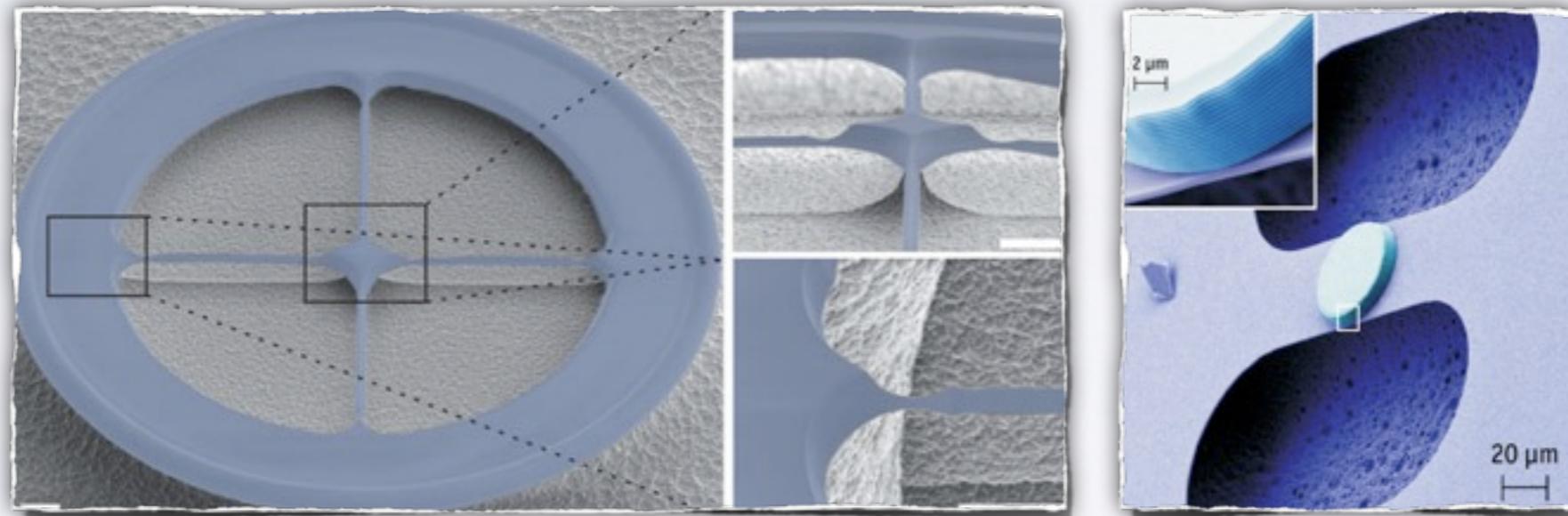


- **Interesting challenge:** Time to prepare topologically ordered states  $O(L)$  for  $L \times L$  lattice can be achieved, ...  
Dengis, Koenig, Pastawski, arXiv:1310.1036
- ... stability relies on log-Sobolev-type clustering not allowing for topological order
- How to reconcile that? **Dissipative stable passive quantum memories?**

# Clustering of correlations in thermal states

# Locality of temperature?

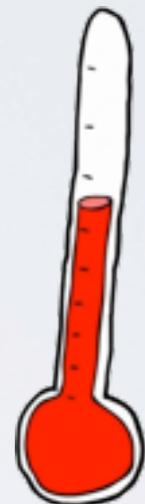
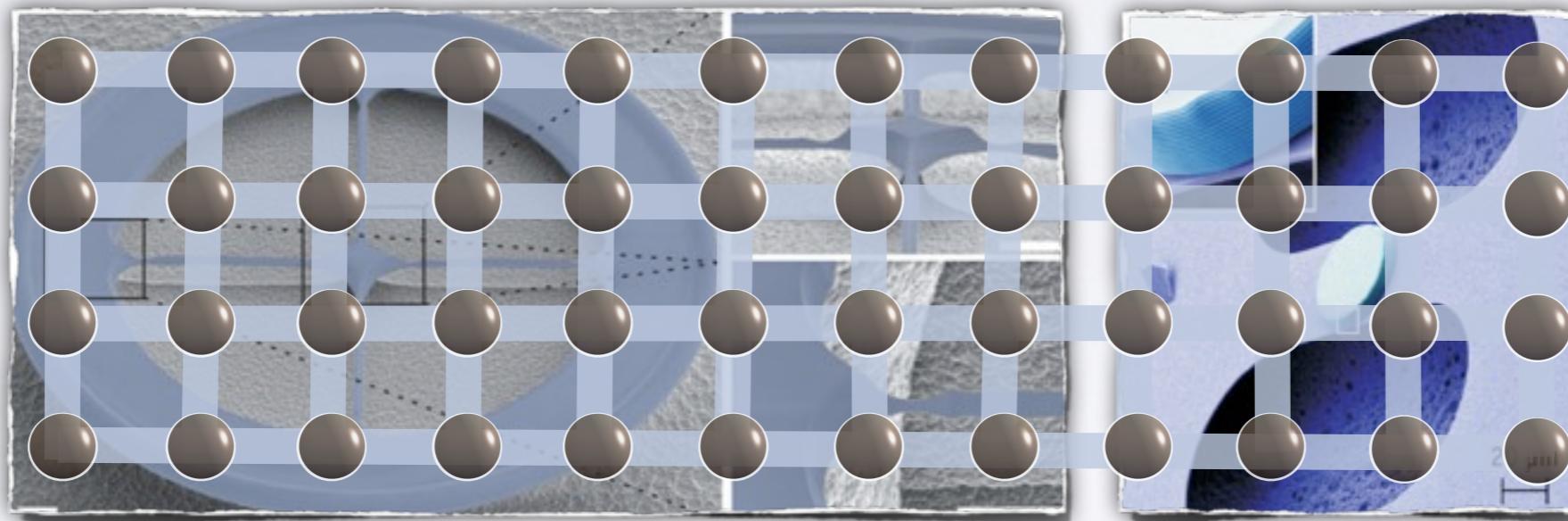
- At what **length scales** is temperature well-defined?



Pothier, Gueron, Birge, Esteve, Devoret, Phys Rev Lett 79, 3490 (1997)  
Peng, Su, Liu, Yu, Cheng, Bao, Nanoscale 5, 9532 (2013)  
Hartmann, Mahler, Hess, Phys Rev Lett 93, 080402 (2004)  
Ferraro, Garcia-Saez, Acin, Europhys Lett 98, 10009 (2012)

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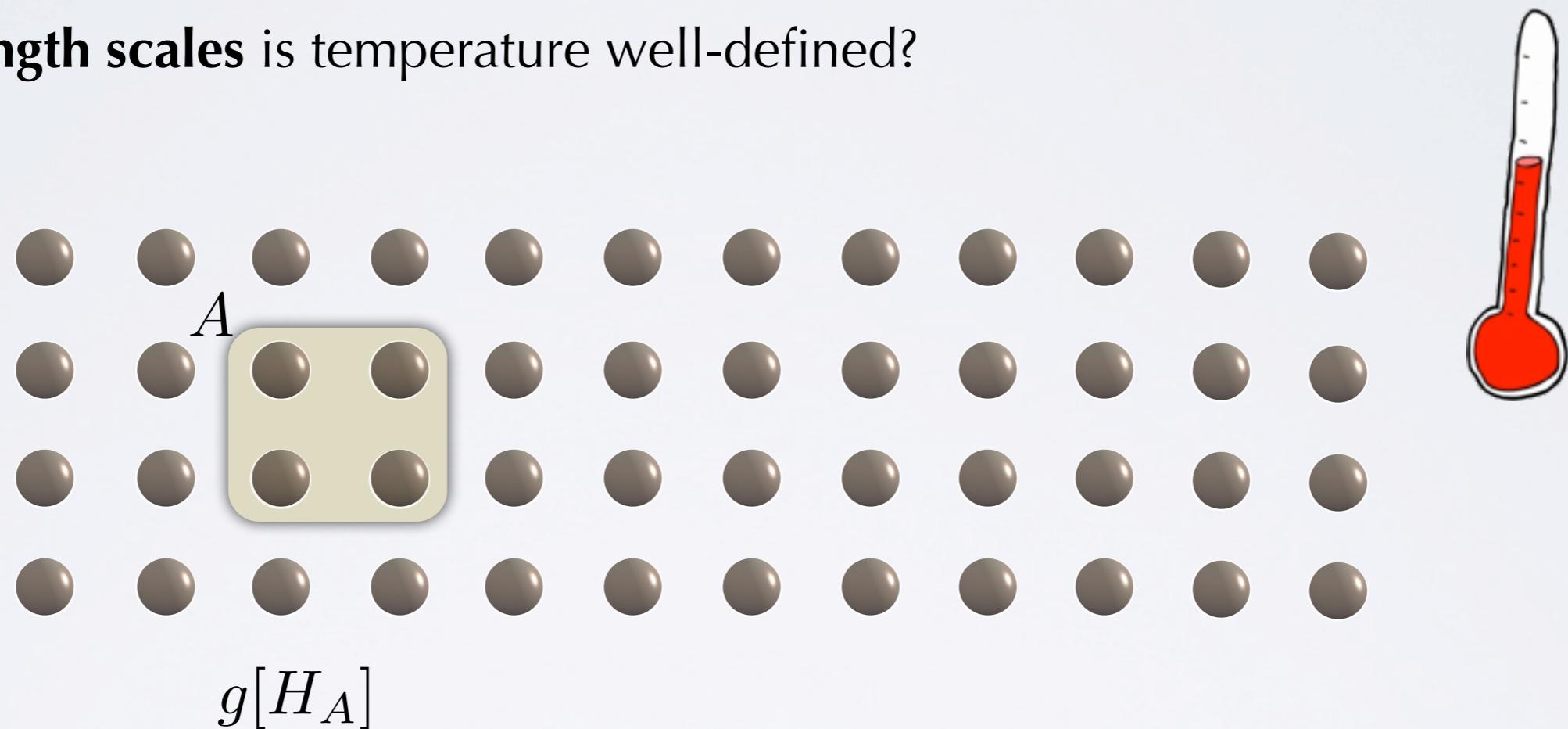
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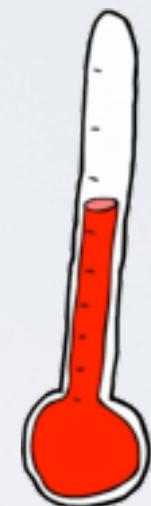
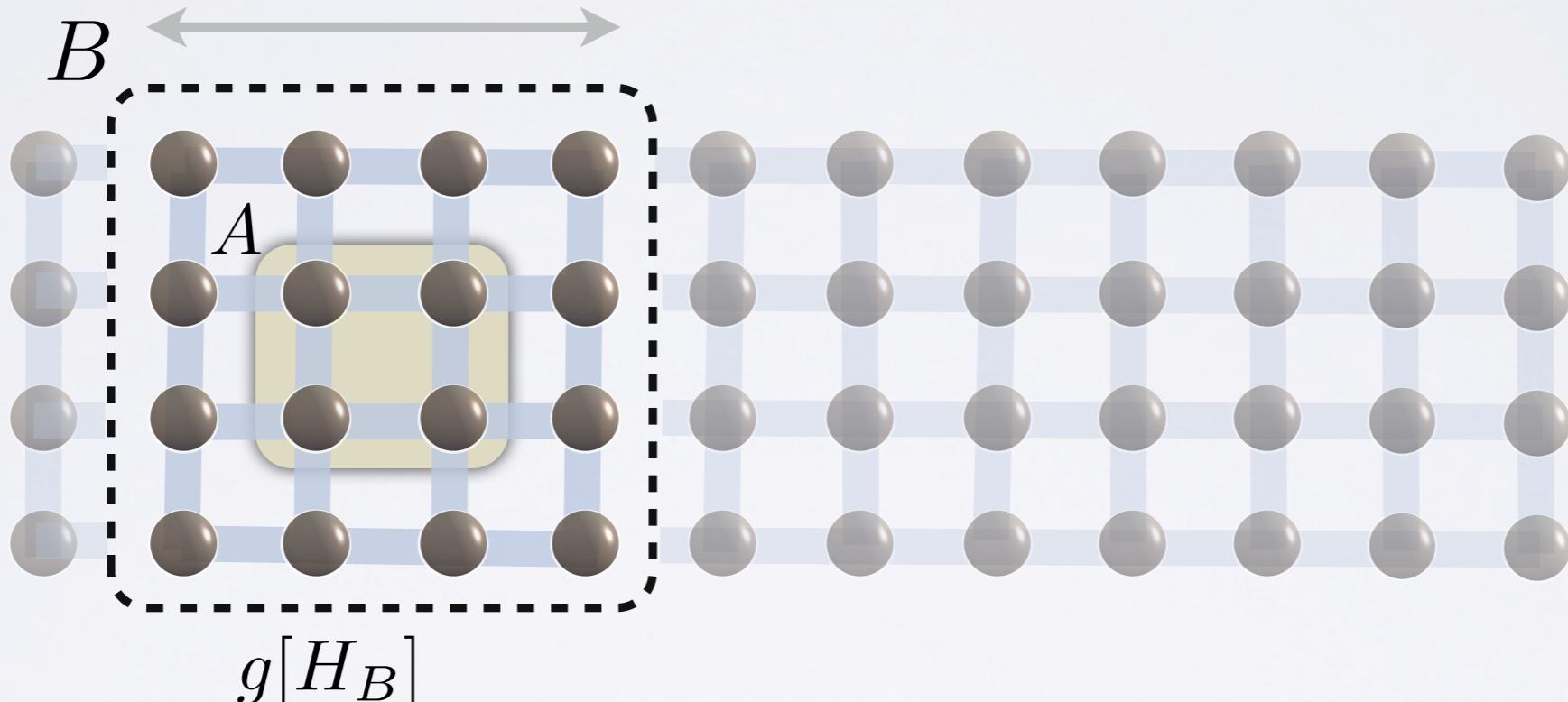
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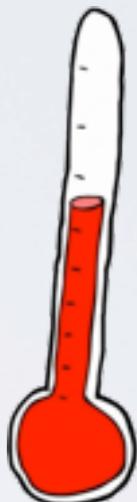
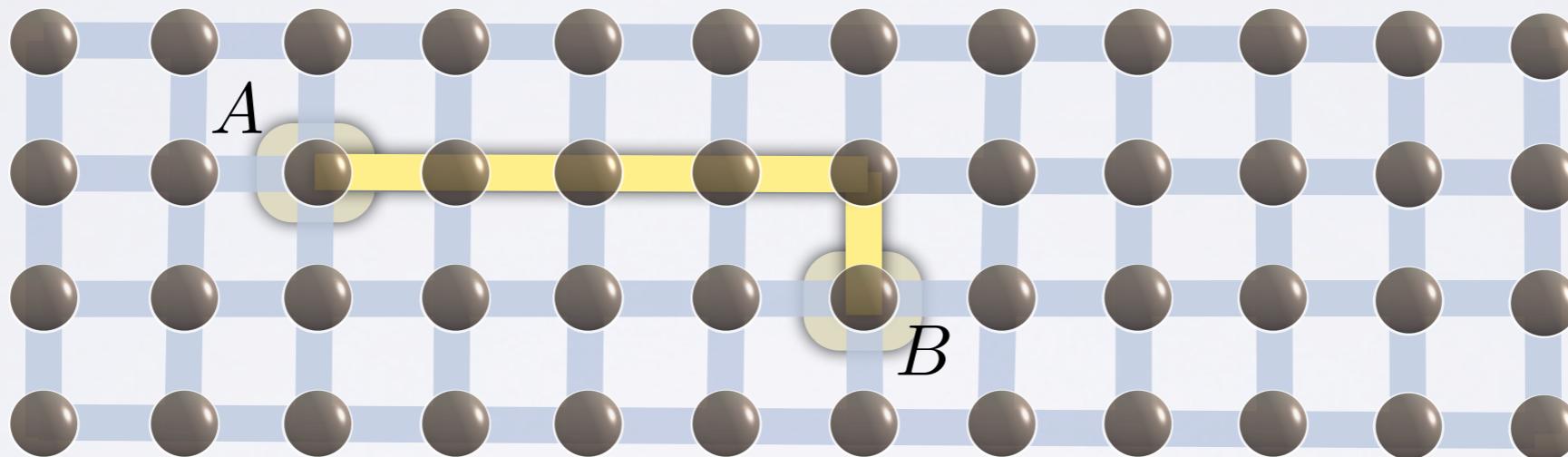
- Gibbs states  $g[H] = \frac{e^{-\beta H}}{\text{tr}(e^{-\beta H})}$

# Locality of temperature?

- At what **length scales** is temperature well-defined?



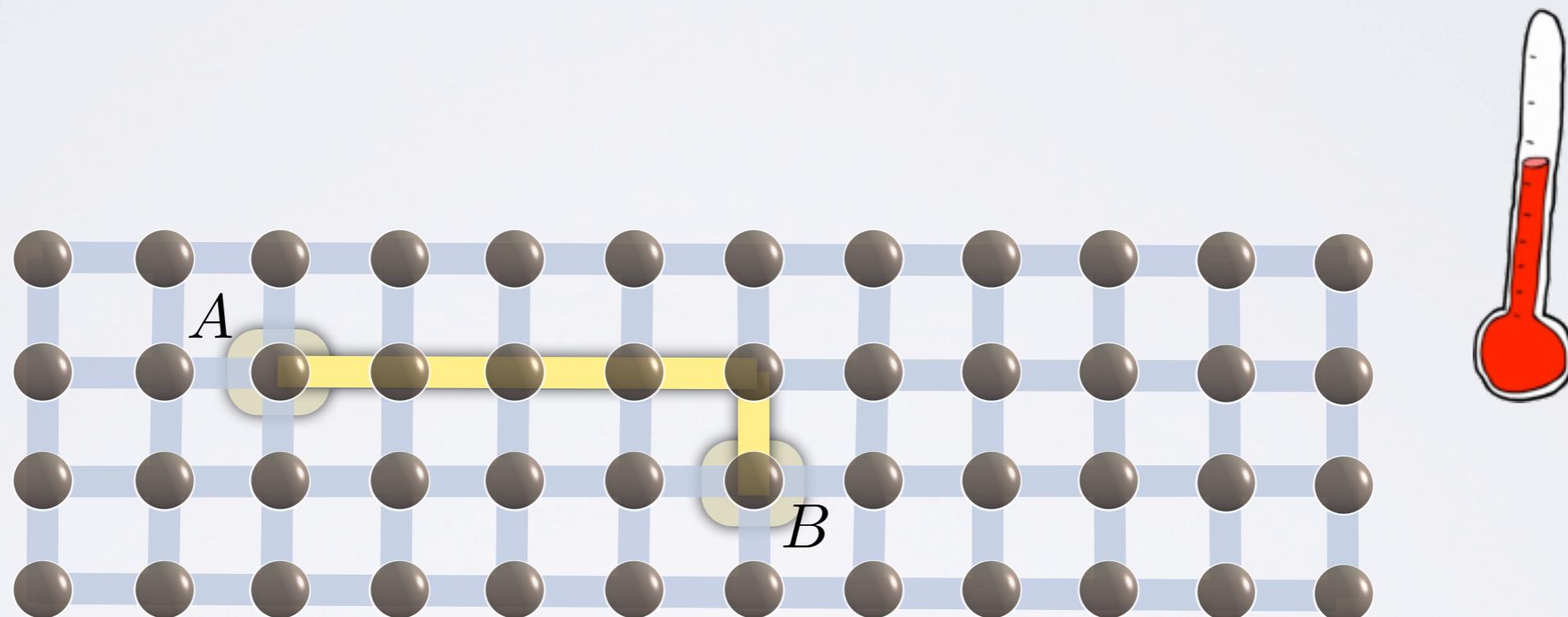
# Thermal states of quantum many-body systems



- Again, GS of gapped Hamiltonians have clustering correlations

• Is there a thermal analogue?

# Thermal states of quantum many-body systems



- **Critical temperature**, dependent only on crude properties of graph (+ coupling strength), above which correlations cluster?
- Long-standing **open question**, results known for **classical and continuum models**, some (few) insights into quantum lattice models

Araki, Commun Math Phys 38,1 (1974)

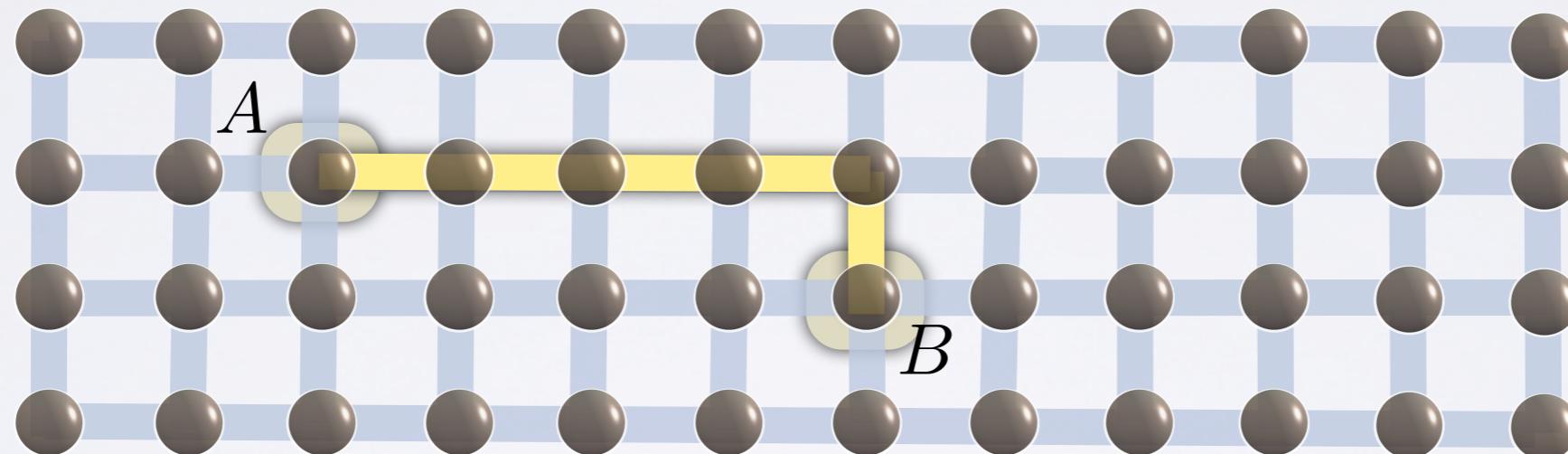
Ruelle, Rev Mod Phys 36, 580 (1964)

Ginibre, J Math Phys 6, 252 (1965)

Greenberg, Commun Math Phys 13, 335 (1969)

Brattelli, Robinson, Operator algebras in quantum statistical mechanics (Springer, 1981)

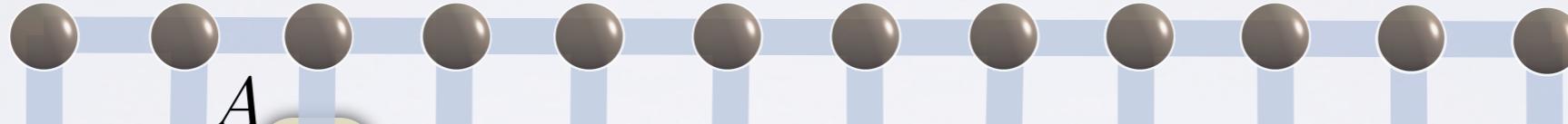
# Thermal states of quantum many-body systems



• Yes :)

# General clustering of correlations at high temperatures

$$\xi(\beta) = \left| 1 / \ln(\alpha e^{2|\beta|J} (e^{2|\beta|J} - 1)) \right|$$



- **Clustering of correlations in thermal states:** Consider local Hamiltonian on arbitrary regular lattice,  $J := \max \|h_k\|$  coupling strength, then exists critical inverse temperature

$$\beta^* := \log((1 + \sqrt{1 + 4/\alpha}/2)/(2J))$$

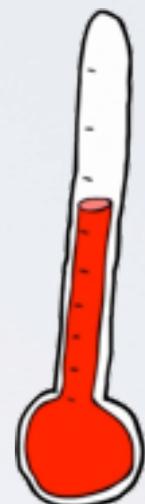
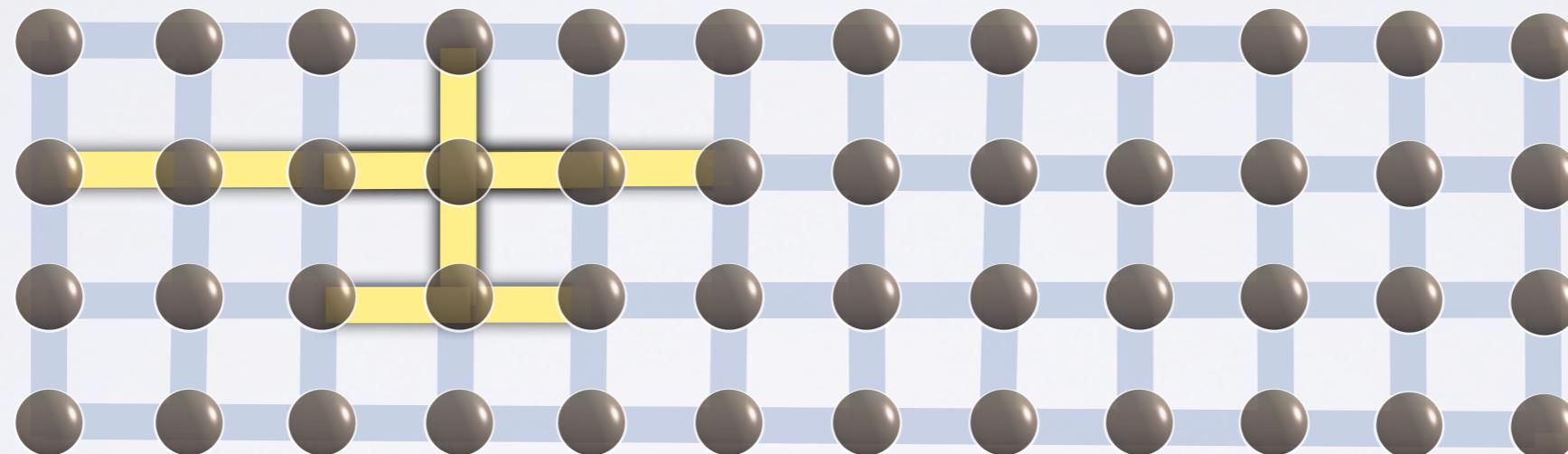
such that for all  $\beta < \beta^*$  and  $d(A, B) \geq L_0$

$$C_{g[H]}(A, B) \leq \frac{4 \min\{|\partial A|, |\partial B|\}}{\log(3)} \frac{\|f\| \|g\|}{1 - e^{-1/\xi(\beta)}} e^{-d(A, B)\xi(\beta)}$$

$$\xi(\beta) := \left| 1 / \ln(\alpha e^{2|\beta|J} (e^{2|\beta|J} - 1)) \right|$$

- **$\alpha$  lattice animal constant**
- General statement for **arbitrary lattices** and **covariances**

# Lattice animal constants



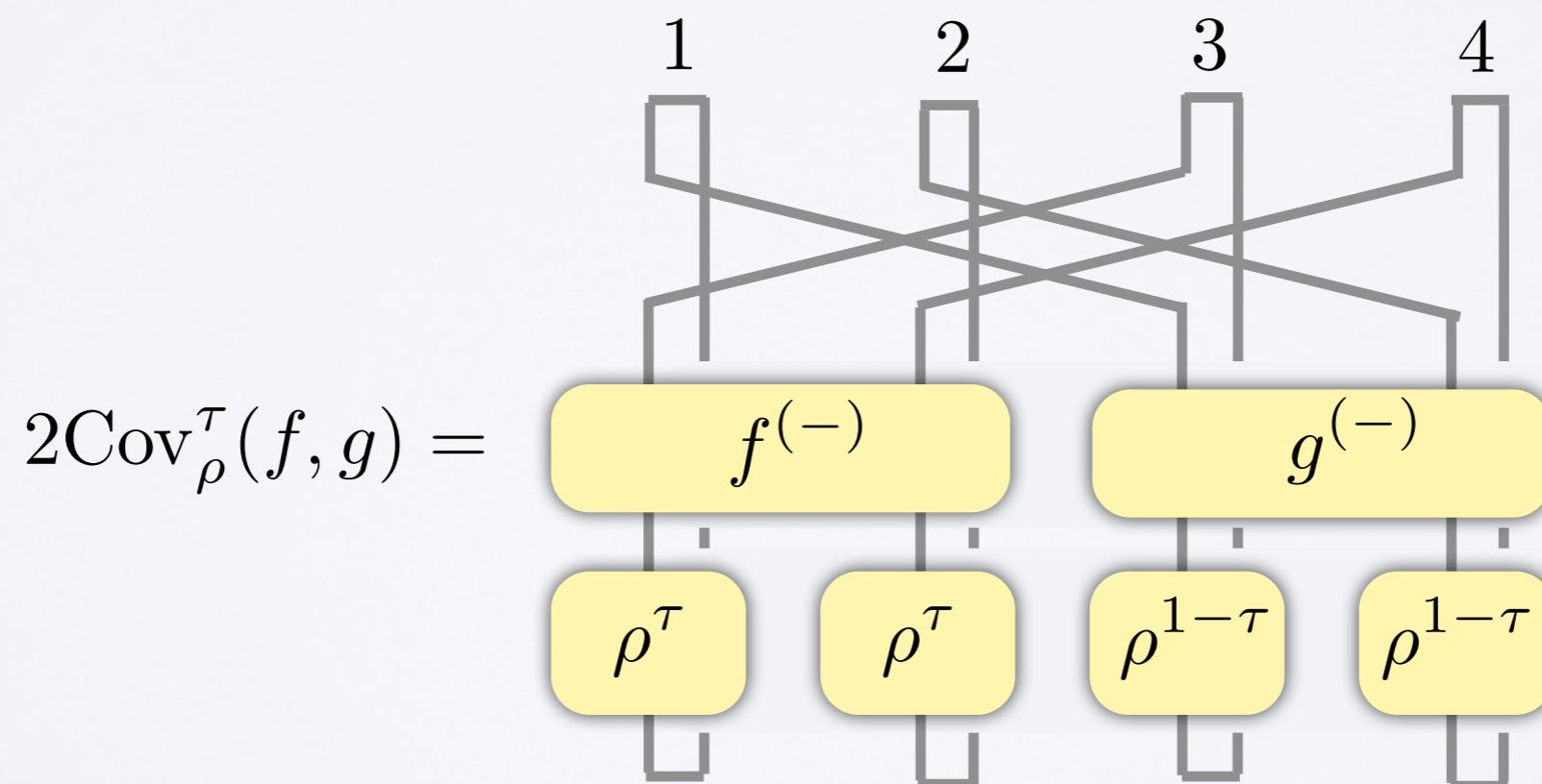
- Connected set of edges: **Lattice animal**
- Number  $a_m$  of lattice animals  $F$  of size  $m = |F|$
- Lattice animal constant: Smallest  $\alpha$  such that  $a_m \leq \alpha^m$

# Flavour of (involved) proof

Define **generalized covariance**  $\text{Cov}_\rho^\tau(f, g) = \text{tr}(\rho^\tau f \rho^{1-\tau} g) - \text{tr}(\rho f) \text{tr}(\rho g)$ ,  $\tau \in [0, 1]$

↓  
**Multiple "swap-trick"**

Write  $\text{Cov}_\rho^\tau(f, g) = \frac{1}{2} \text{tr} \left( \mathcal{S}^{(1,3)} \mathcal{S}^{(2,4)} (f^{(-)} \otimes g^{(-)}) (\rho^\tau \otimes \rho^\tau \otimes \rho^{1-\tau} \otimes \rho^{1-\tau}) \right)$ ,  
on four copies, where  $f^{(-)} = f \otimes 1 - 1 \otimes f$



# Flavour of (involved) proof

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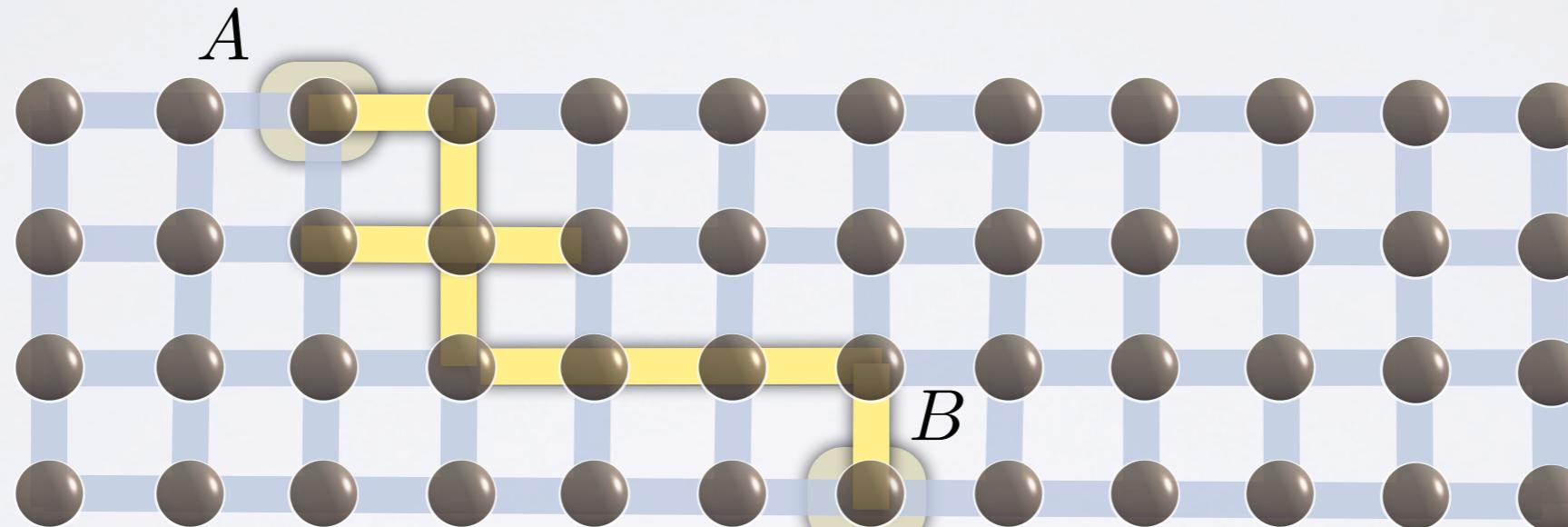
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**Multiple "swap-trick"**

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on four copies, where  $f^{(-)} = f \otimes 1 - 1 \otimes f$

↓  
**Cluster expansion of new Hamiltonian  $\tilde{H}$**

$$\frac{e^{-\beta \tilde{H}}}{\text{tr}(e^{-\beta \tilde{H}})} = \rho^\tau \otimes \rho^\tau \otimes \rho^{1-\tau} \otimes \rho^{1-\tau}$$

# Flavour of (involved) proof



**Symmetry:** Only clusters connecting  $A$  and  $B$  contribute

**Cluster expansion of new Hamiltonian  $\tilde{H}$**

$$\frac{e^{-\beta \tilde{H}}}{\text{tr}(e^{-\beta \tilde{H}})} = \rho^\tau \otimes \rho^\tau \otimes \rho^{1-\tau} \otimes \rho^{1-\tau}$$

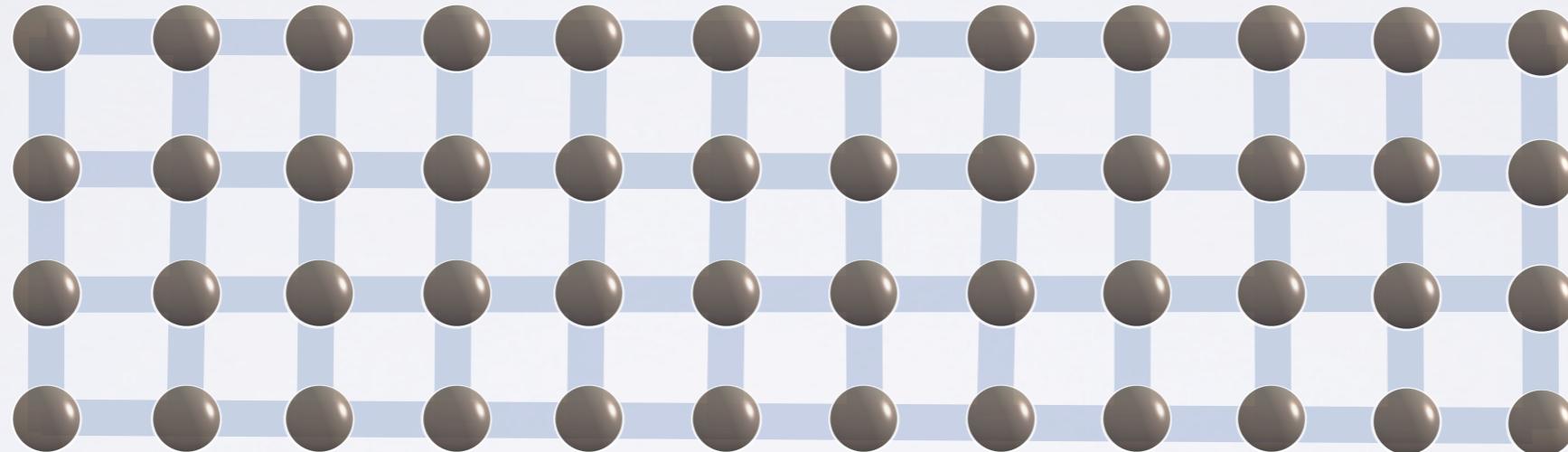
**Truncated cluster expansion**

$$\left\| \sum_{w \in C_{\geq L}(F)} \frac{(-\beta)^{|w|}}{|w!|} h(w) \right\|_1 \leq Z(\beta) \left( e^{|F| \frac{b(\beta)L}{1-b(\beta)}} - 1 \right)$$

**Combinatorics**

# Physical implications: Bounds to Curie temperatures

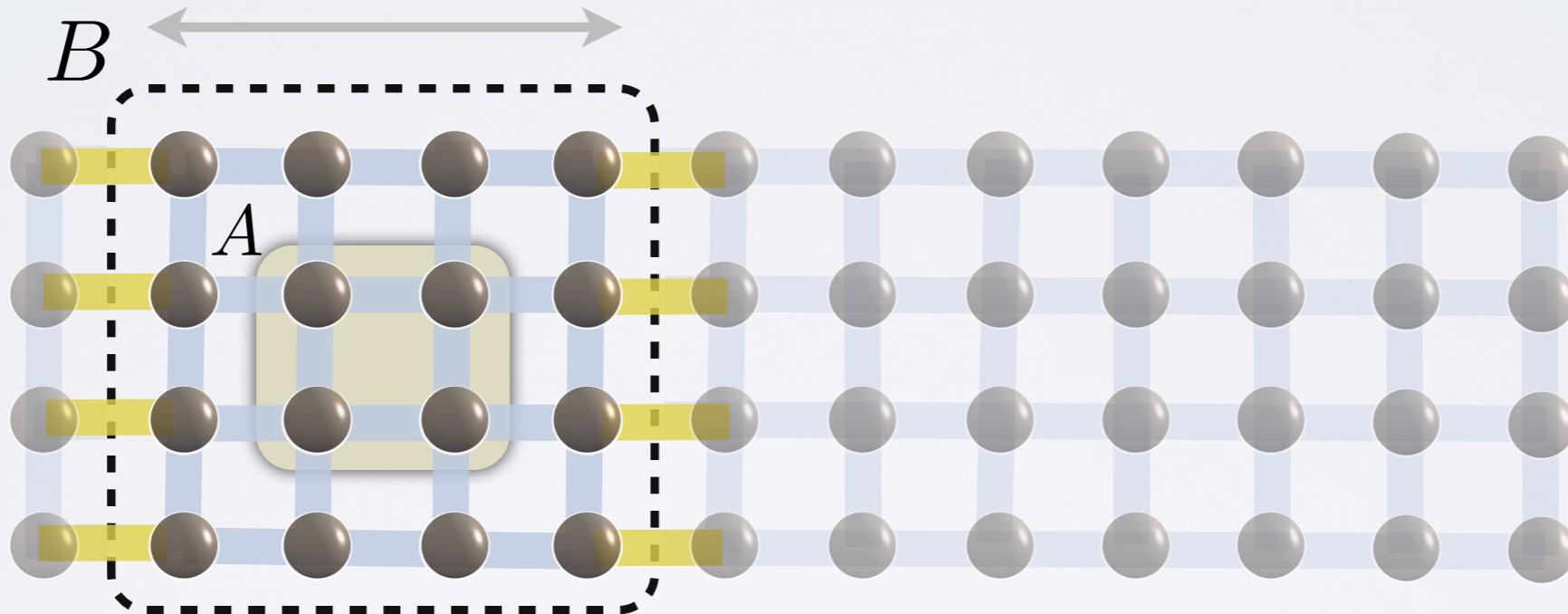
1. "Critical temperature" is **universal upper bound** to phase transition points



- E.g., ferromagnetic **2d isotropic Ising model** without external field,  
 $1/(J\beta^*) = 24.58$ , while phase transition known to happen at 2.27

# Locality of temperature

## 2. Length scale of temperature

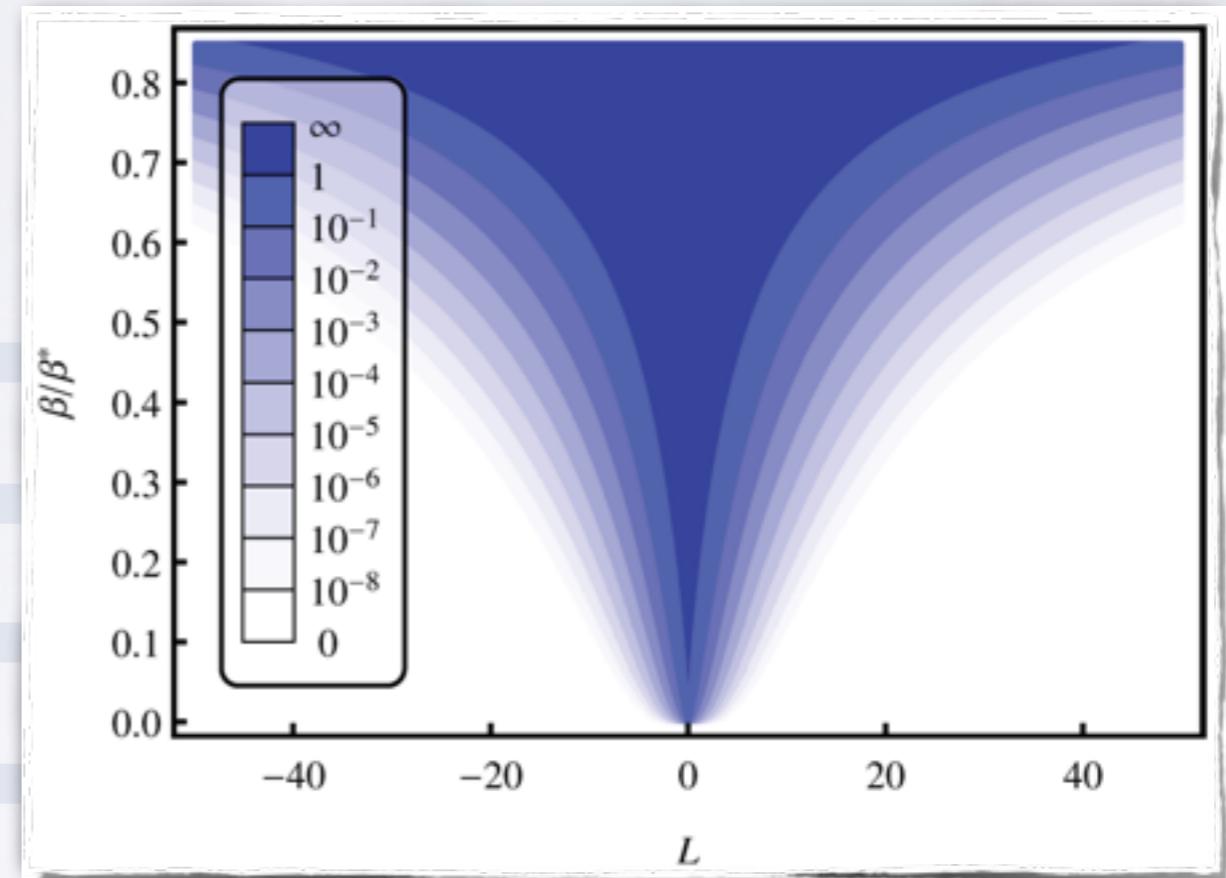
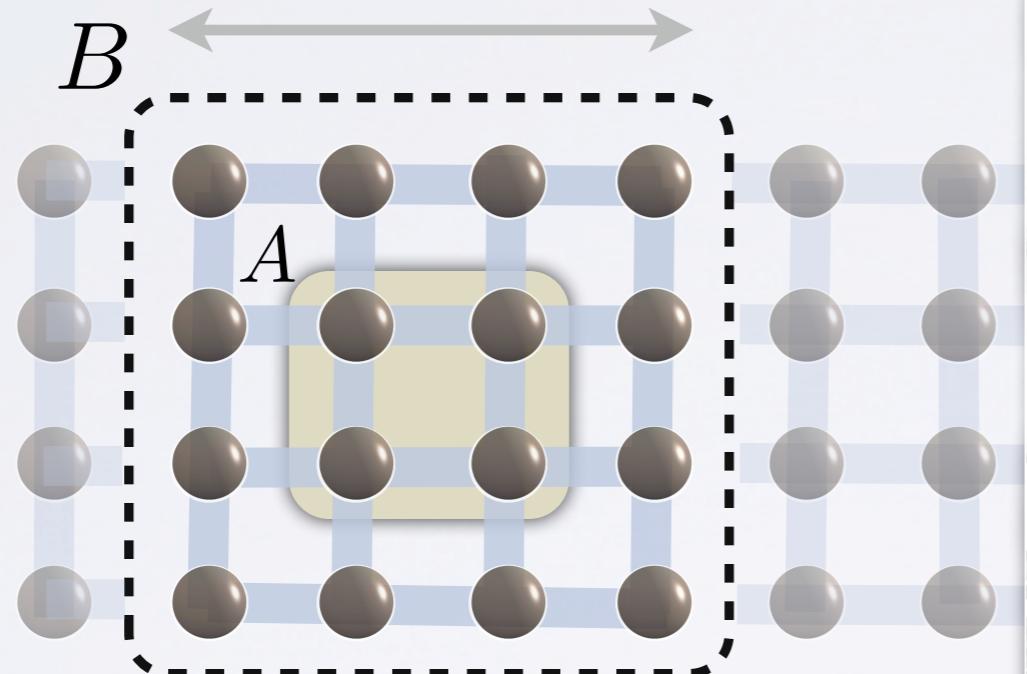


$$\text{tr}(Ag[H(0)]) - \text{tr}(Ag[H]) = \beta \int_0^1 d\tau \int_0^1 ds \text{Cov}_{g[H(s)]}^\tau(A, H_I)$$

$$H(s) = H - (1 - s)H_I$$

# Locality of temperature

## 2. Length scale of temperature



$$\|g_A[H] - g_A[H(0)]\|_1 \leq \frac{v|\beta|J}{1 - e^{-1/\xi(\beta)}} e^{-d(A, \partial A)/\xi(\beta)}$$

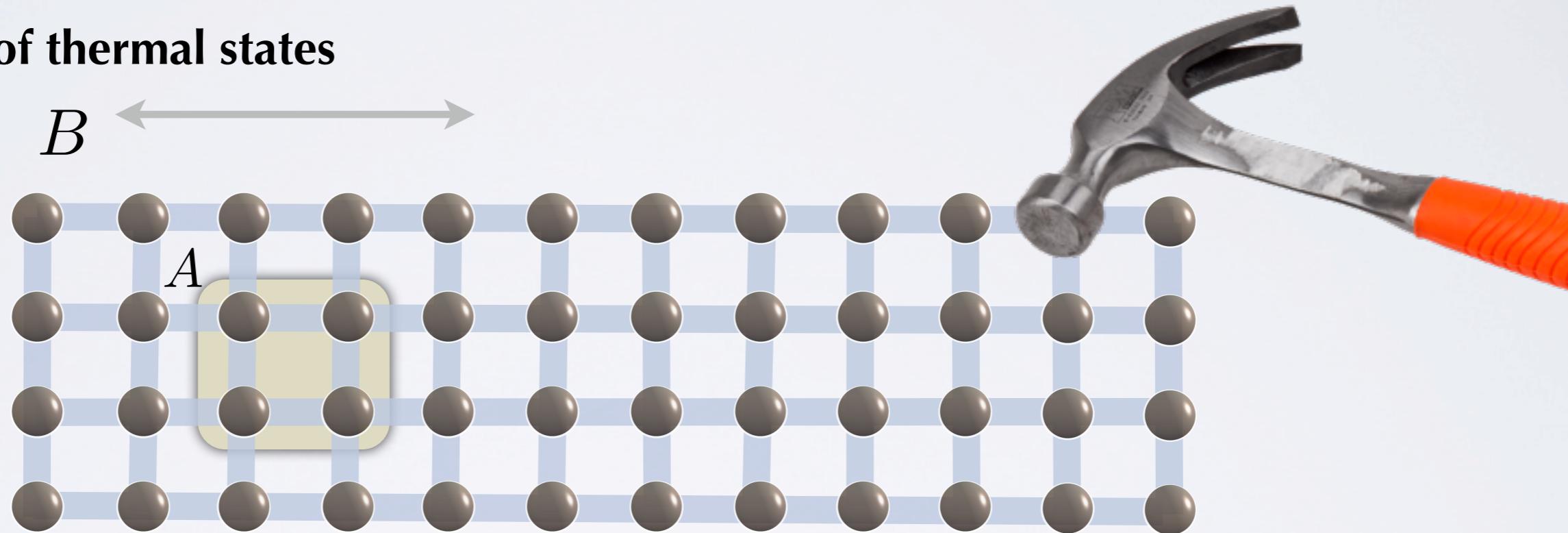
Kliesch, Gogolin, Kastoryano, Riera, Eisert, arXiv:1309.0816

- Tool in rigorous approaches in **quantum thermodynamics** and **canonical typicality**

Popescu, Short, Winter, Nature Phys 2, 754 (2006)  
Goldstein, Lebowitz, Tumulka, Zanghi, Phys Rev Lett 96, 050403 (2006)  
Reimann, Phys Rev Lett 101, 190403 (2008)  
Linden, Popescu, Short, Winter, Phys Rev E 79, 061103 (2009)  
Riera, Gogolin, Eisert, Phys Rev Lett 108, 080402 (2012)

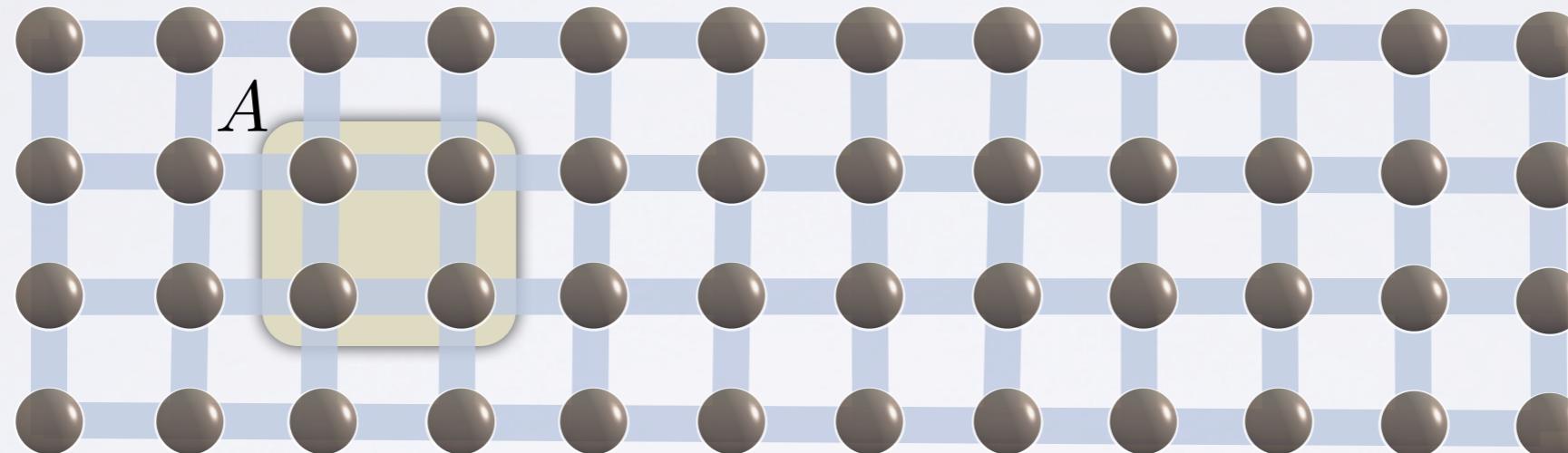
# Stability of high temperature thermal states

## 3. Stability of thermal states



# Efficient computation of expectation values

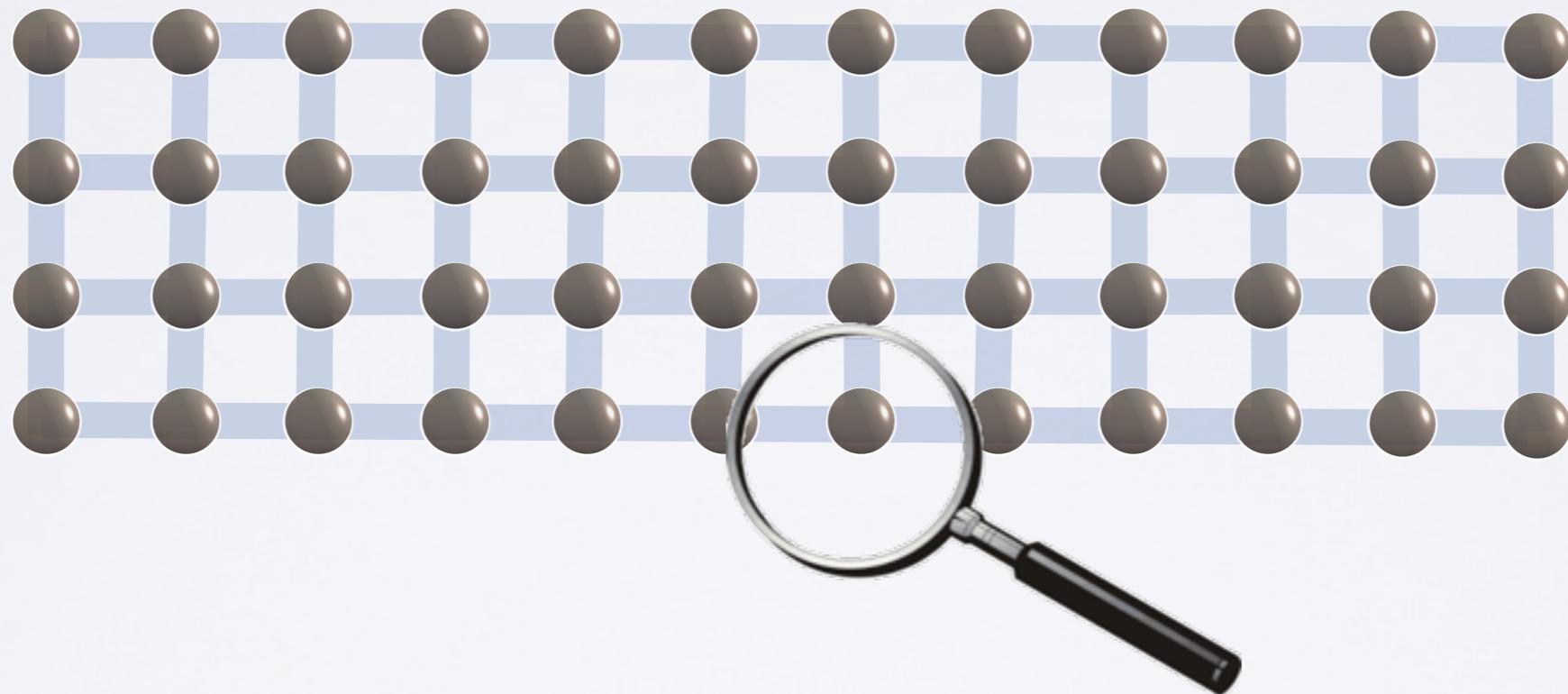
## 4. Computing of **local expectation values** is in P



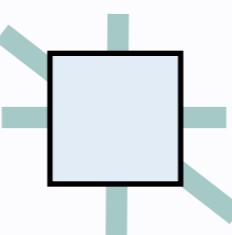
- Guideline for quantum Monte Carlo etc

# Matrix-product operators

## 5. Matrix-product operator approximation

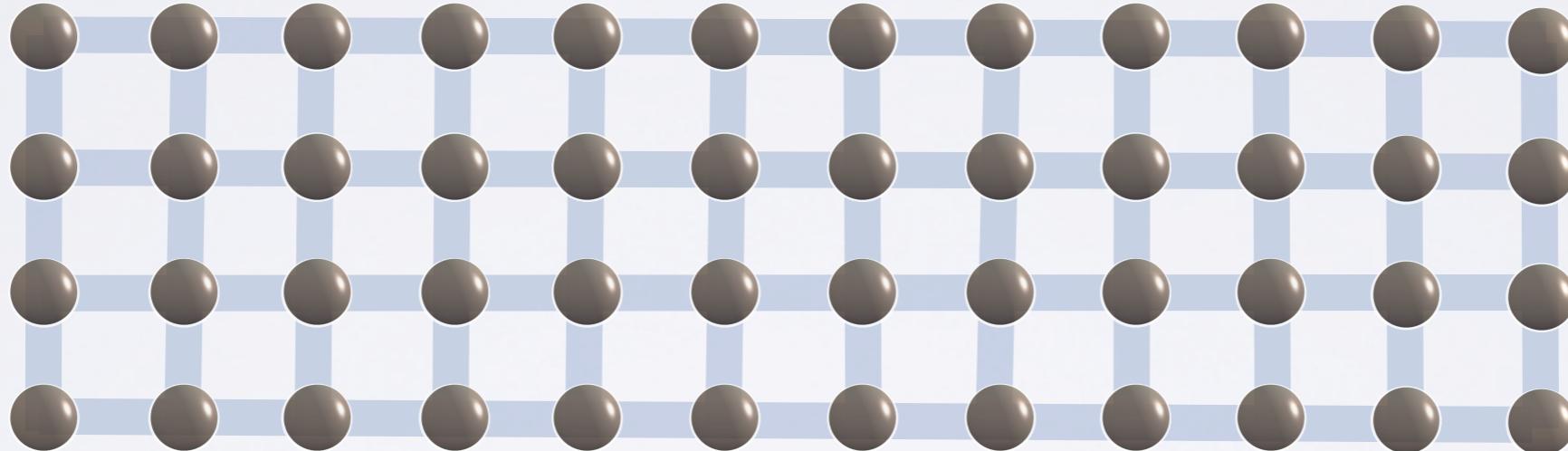


MPO approximation of sub-exponential dimension, efficient in 1D



# Interacting fermions

## 6. All also true for interacting fermions



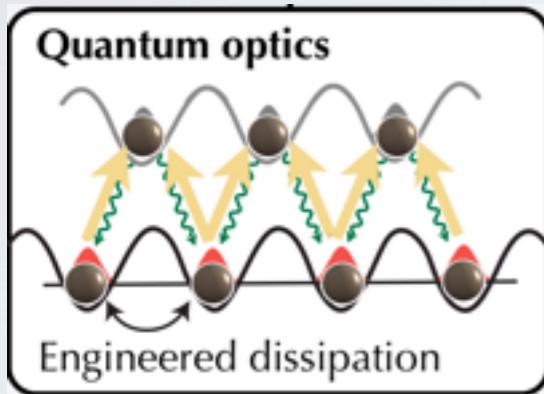
$$\{f_j, f_k^\dagger\} = \delta_{j,k}$$

- Generalizing earlier results on fermionic covariance matrices

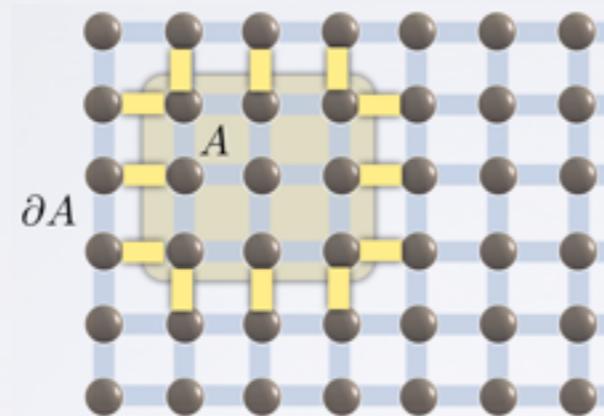
# Thermal many-body systems

- **Lessons:**
- **Length scale** at which one can speak of temperature!
- High temperature thermal states have **clustering correlations**
- Can **efficiently compute local expectation values**

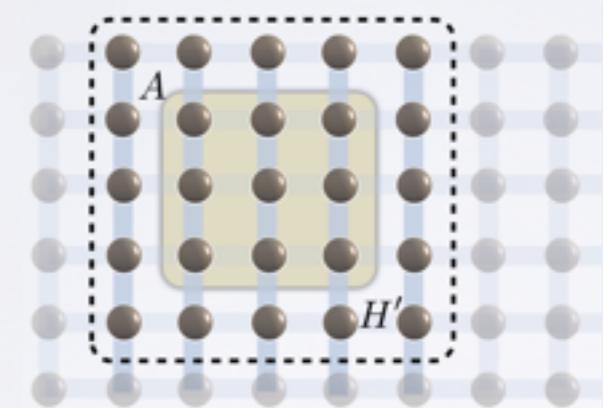
# Summary: Liouvillians and thermal states



"Liouvillian complexity":  
Interesting arena to study  
many-body problems



Clustering of correlations, area  
laws, topological order



Intensity of temperature  
and correlations in thermal  
many-body systems

- Does the exponential complexity of general quantum systems persist at high temperature?

No

## Thanks for your attention!